

# Lieb-Liniger variational ansatz and one-dimensional dipolar gases

Roberta Citro<sup>1</sup>, Marilu Chiofalo<sup>2</sup>, Stefania De Palo<sup>3</sup> and Edmond Orignac<sup>4</sup>

<sup>1</sup> Dipartimento di Fisica "E. R. Caianiello", Università di Salerno, Italy

<sup>2</sup> Dipartimento di Fisica, Università di Pisa, Italy

<sup>3</sup> CNR-IOM Democritos and Università di Trieste, Italy

<sup>4</sup> Laboratoire de Physique de l'ENS de Lyon, CNRS UMR 5672, France

November 3rd, 2020

## Abstract

We introduce a variational ansatz using the ground state wavefunction of the integrable Lieb-Liniger model as trial wavefunction. The expectation value of the potential energy is obtained from an approximation of the static structure factor of the Lieb-Liniger model introduced by Cherny and Brand. We apply the variational ansatz to the calculation of the ground state energy of a quasi-one dimensional gas of dipolar bosons as a function of density and tilt angle of the dipoles. We discuss the breathing modes of the dipolar gas.

## Dipolar bosons trapped in one dimension

$$H_{Q1D} = -\frac{\hbar^2}{2m} \sum_i \frac{\partial^2}{\partial x_i^2} + \sum_{i<j} V_{Q1D}(x_i - x_j) + g_{VdW} \sum_{i<j} \delta(x_i - x_j)$$

## Dipolar interactions in the single mode approximation

$$V_{Q1D}(x) = V(\theta) \left[ V_{DDI}^{1D} \left( \frac{x}{l_{\perp}} \right) - \frac{8}{3} \delta \left( \frac{x}{l_{\perp}} \right) \right],$$

$$V(\theta) = \frac{\mu_0 \mu_D^2}{4\pi} \frac{1 - 3 \cos^2 \theta}{4l_{\perp}^3},$$

$$V_{DDI}^{1D} \left( \frac{x}{l_{\perp}} \right) = -2 \left| \frac{x}{l_{\perp}} \right| + \sqrt{2\pi} \left[ 1 + \left( \frac{x}{l_{\perp}} \right)^2 \right] \times e^{\frac{1}{2} \left( \frac{x}{l_{\perp}} \right)^2} \operatorname{erfc} \left[ \left| \frac{x}{\sqrt{2} l_{\perp}} \right| \right],$$

F. Deuretzbacher, J. C. Cremon, and S. M. Reimann, Phys. Rev. A **81**, 063616 (2010); **87**, 039903(E) (2013).

## Variational method

$$H_{1D} = \mathcal{K} + \mathcal{V},$$

$$H_{\text{var}} = \mathcal{K} + g\mathcal{U},$$

$$H_{\text{var}} |\psi_0(g)\rangle = E_0(g) |\psi_0(g)\rangle$$

$$E_{\text{var}}(g) = E_0(g) - g \frac{\partial E_0(g)}{\partial g} + \langle \psi_0(g) | \mathcal{V} | \psi_0(g) \rangle.$$

The best approximation is obtained by minimizing with respect to  $g$ .

## Choice of variational Hamiltonian

$$\mathcal{U} = \sum_{1 \leq i < j \leq N} \delta(x_i - x_j),$$

$H_{\text{var}}$  is the Hamiltonian of the Lieb-Liniger model, integrable by the Bethe Ansatz.  $E_0(g)$  is obtained from the solution of an integral equation. E. H. Lieb and W. Liniger, Phys. Rev. **130**, 1605 (1963)

## Variational energy as a function of the structure factor

$$\frac{\langle \psi_0(g) | \mathcal{V} | \psi_0(g) \rangle}{L} = \frac{n}{2} \int_0^{+\infty} \frac{dk}{\pi} \hat{v}(k) [S(k) - 1] + \frac{n^2}{2} \hat{v}(k=0),$$

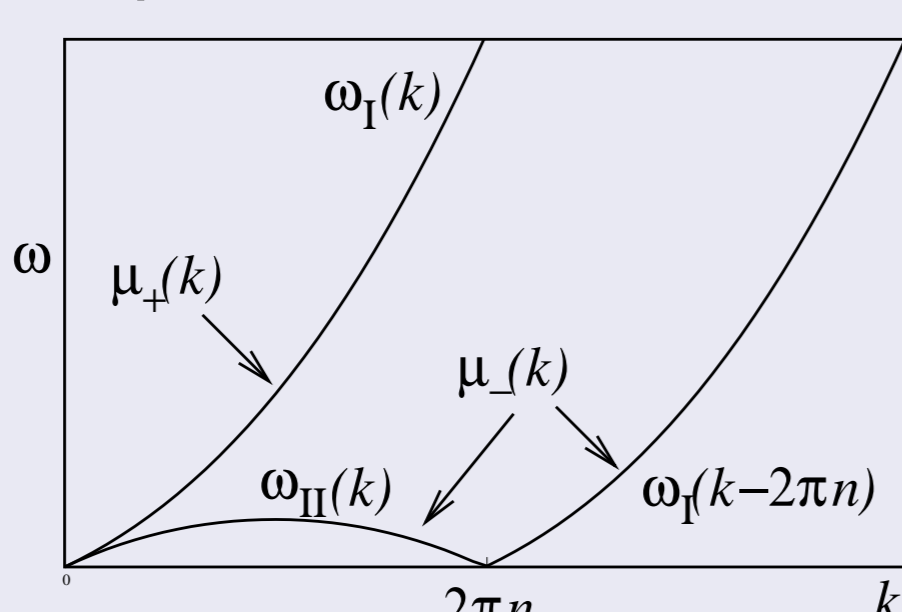
where  $S(k)$  is the static structure factor of the Lieb-Liniger model.

## Cherny Brand Ansatz for the structure factor of the Lieb-Liniger gas

$$S(k) = \frac{k^2}{2m\omega_{II}(k)}$$

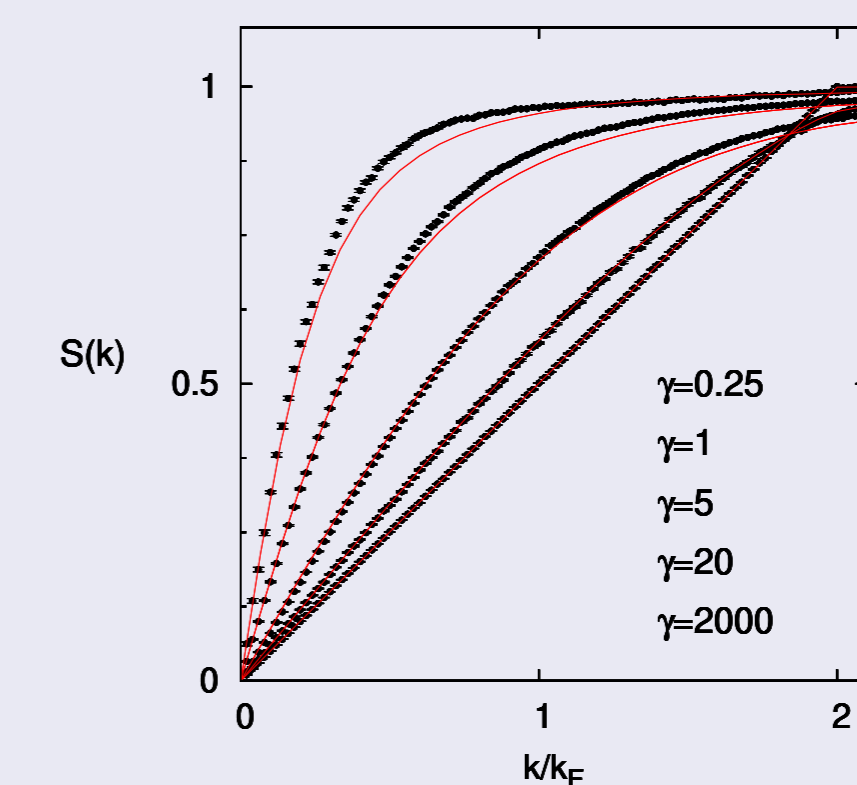
$$\times \frac{2F_1 \left( 1 + \frac{\sqrt{K}}{1+\sqrt{K}} + \mu_-(k) + \mu_+(k), 1 + \mu_-(k); 2 + \mu_-(k) - \mu_+(k), 1 - \frac{\omega_{II}^2(k)}{\omega_I^2(k)} \right)}{2F_1 \left( 1 + \frac{2\sqrt{K}}{1+\sqrt{K}} + \mu_-(k) + \mu_+(k), 1 + \mu_-(k); 2 + \mu_-(k) - \mu_+(k), 1 - \frac{\omega_{II}^2(k)}{\omega_I^2(k)} \right)}$$

- $\omega_I(k)$ ,  $\mu_+(k)$  dispersion and threshold exponent of Type-I excitations.
- $\omega_{II}(k)$ ,  $\mu_-(k)$  dispersion and threshold exponent of Type-II excitations.
- $K$  Tomonaga-Luttinger exponent.

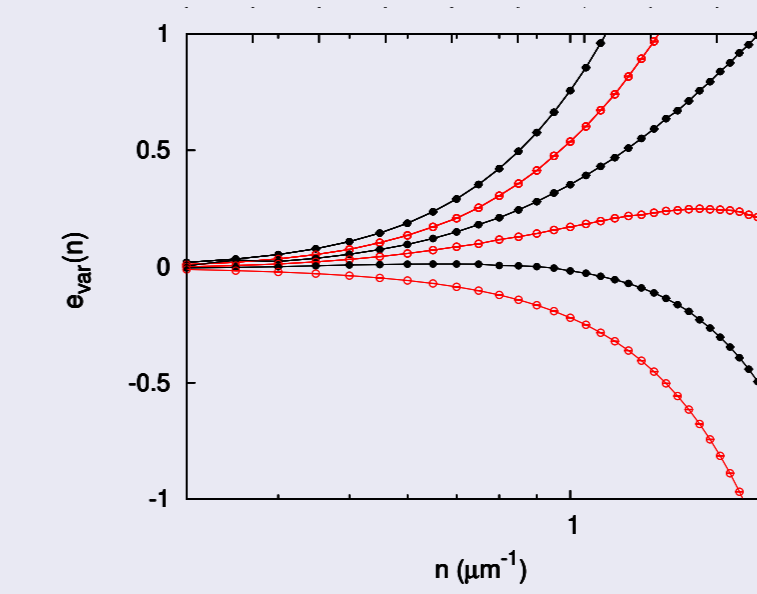


A. Y. Cherny and J. Brand, Phys. Rev. A **79**, 043607 (2009)

## Comparison of the Ansatz for $S(k)$ with Quantum Monte Carlo



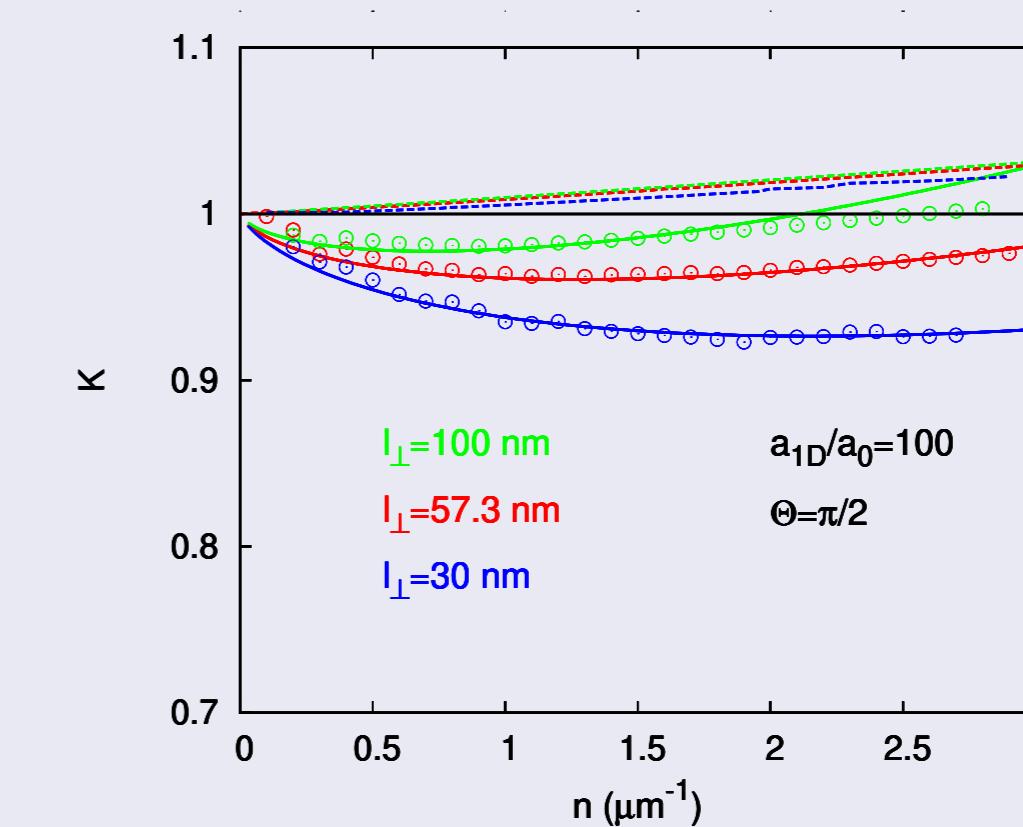
## Variational energy per unit length for $\theta = 0$



Energy per unit length as a function of atom density for selected scattering lengths. Top to bottom  $a_{1D}/a_0 = 5000, 6000, 7000, 8000, 9000$  and  $10000$

## Tomonaga-Luttinger exponent $K$

$$\frac{u}{K} = \frac{1}{\pi \hbar} \frac{\partial^2 e_{GS}}{\partial n^2}, \quad uK = \frac{\hbar \pi n}{m},$$

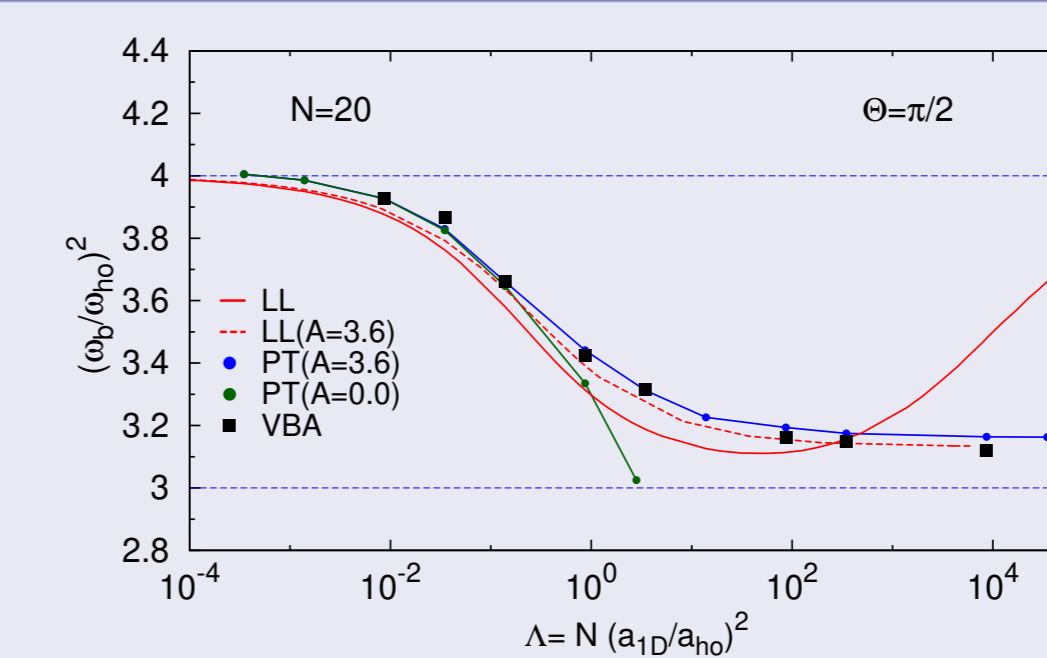


## Breathing modes

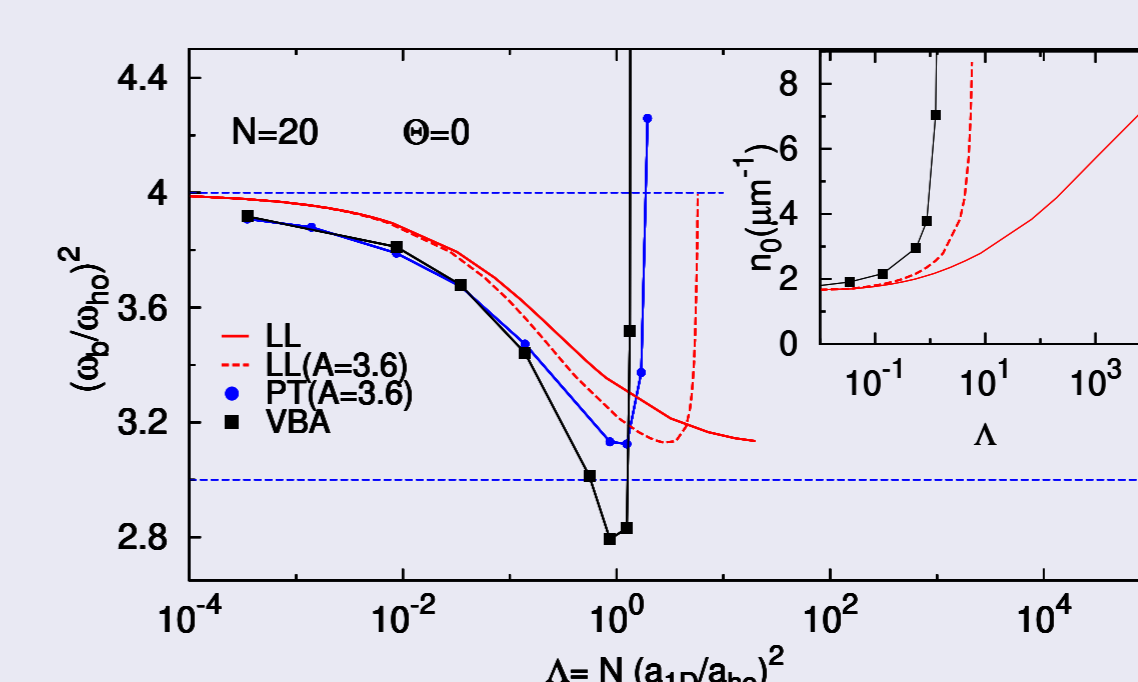
- 1 Add a longitudinal trapping potential  $\frac{1}{2} m \omega_{ho}^2 x r$ .
- 2 find density profile from stationary generalized Gross-Pitaevskii equation
- 3 obtain breathing mode from sum rule approach

$$\omega_B^2 = -2 \left\langle \sum_{i=1}^N x_i^2 \right\rangle \left[ \frac{\partial \langle \sum_{i=1}^N x_i^2 \rangle}{\partial \omega_{ho}^2} \right]^{-1}$$

## Calculated breathing mode in the repulsive case



## Calculated breathing mode in the attractive case



## Perspectives

- 1 Other non-contact interactions: shoulder potential, etc. . .
- 2 Droplets
- 3 Super-Tonks type metastable states
- 4 Generalized time-dependent Gross-Pitaevskii equation