

Density of states of a two-band one-dimensional conductor

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Hubbard model in magnetic field

$$\begin{aligned}
 H &= -t \sum_{j,a} (c_{j,a}^\dagger c_{j+1,a} + \text{H.c.}) + U \sum_j n_{j,\uparrow} n_{j,\downarrow} \\
 &\quad - \sum_{j,a} (\mu + g\mu_B B a) n_{j,a}
 \end{aligned}$$

Retardated Green's function

$$\begin{aligned}
 G_a(j, t) &= -i\theta(t) \langle \{c_{j,a}(t), c_{0,a}^\dagger(0)\} \rangle \\
 A(q, \omega) &= \sum_j \int dt G_a(j, t) e^{-i(qj - \omega t)}
 \end{aligned}$$

- map to two bands of interacting spinless fermions.
- bosonization method

Definitions for bosonization

$$\begin{aligned}
 x &= j\alpha \\
 [\phi_a(x), \partial_x \theta_b(y)] &= i\pi \delta_{ab} \delta(x-y) \\
 c_{j,a} &\sim e^{ik_{F,a}x + \theta_a - \phi_a} + e^{-ik_{F,a}x + \theta_a + \phi_a}
 \end{aligned}$$

Bosonized Hamiltonian

$$H = \int \frac{dx}{2\pi} \left[{}^t(\partial_x \theta) M \partial_x \theta + {}^t(\partial_x \phi) N \partial_x \phi \right]$$

perturbative expressions

$$\begin{aligned}
 M_{ab} &= v_a \delta_{ab} \\
 N_{ab} &= v_a \delta_{ab} + (1 - \delta_{ab}) U \alpha / \pi \\
 v_a &= 2t\alpha \sin(k_{F,a}\alpha)
 \end{aligned}$$

Diagonalization of bosonized Hamiltonian

$$\Delta_1 = {}^t R_1 M R_1$$

$$\Delta_2 = {}^t R_2 \Delta_1^{1/2} {}^t R_1 N R_1 \Delta_1^{1/2} R_2$$

$$\theta = R_1 \Delta_1^{-1/2} R_2 \Delta_2^{1/4} \vartheta = Q \vartheta$$

$$\phi = R_1 \Delta_1^{1/2} R_2 \Delta_2^{-1/4} \varphi = P \varphi$$

$$H = \int \frac{dx}{2\pi} \left[{}^t (\partial_x \vartheta) \Delta_2^{1/2} \partial_x \vartheta + {}^t (\partial_x \varphi) \Delta_2^{1/2} \partial_x \varphi \right]$$

$\Delta_2^{1/2}$ = diagonal matrix. Its eigenvalues are the velocities of the modes.

Correlation functions

$$\begin{aligned} \langle e^{i(\theta_a - \phi_a)(x,t)} e^{-i(\theta_a - \phi_a)(0,0)} \rangle &= \\ &= e^{\langle (\theta_a(x,t) - \phi_a(x,t))(\theta_a(0,0) - \phi_a(0,0)) \rangle_{conn.}} \end{aligned}$$

$$\langle (\theta_a(x,t) - \phi_a(x,t))(\theta_a(0,0) - \phi_a(0,0)) \rangle_{conn.}$$

$$= \sum_b \nu_b \ln \left(\frac{\alpha}{\alpha + i(\nu_b t - x)} \right) + \nu'_b \ln \left(\frac{\alpha}{\alpha + i(\nu_b t + x)} \right)$$

$$\nu_b = (P_{ab} + Q_{ab})^2 / 4$$

$$\nu'_b = (P_{ab} - Q_{ab})^2 / 4$$

$T = 0$ spectral functions

$$A(k_F + q, \omega) = I(q, \omega) + I(-q, -\omega)$$

$$I(q, \omega) = \frac{C \alpha^{\nu_1 + \nu_2 + \nu'_1 + \nu'_2}}{2\pi} \times \int \frac{e^{i(\omega t - qx)}}{(\alpha + i(\nu_1 t - x))^{\nu_1} (\alpha + i(\nu_2 t - x))^{\nu_2}} dx dt$$

$$\times \frac{1}{(\alpha + i(\nu_1 t + x))^{\nu'_1} (\alpha + i(\nu_2 t + x))^{\nu'_2}}$$

Feynman identity

$$\frac{1}{A_1^{\nu_1} A_2^{\nu_2}} = \frac{\Gamma(\nu_1 + \nu_2)}{\Gamma(\nu_1)\Gamma(\nu_2)} \int_0^1 du \frac{u^{\nu_1-1} (1-u)^{\nu_2-1}}{(A_1 u + A_2(1-u))^{\nu_1+\nu_2}}$$

Master formula

$$I(q, \omega) = \frac{\pi^2}{\Gamma(\nu_1)\Gamma(\nu_2)\Gamma(\nu'_1)\Gamma(\nu'_2)} \int_0^1 du_1 u_1^{\nu_1-1} (1-u_1)^{\nu_2-1} \\ \times (\omega - v(u_1)q)^{\nu'_2+\nu'_1-1} \int_0^1 du_2 u_2^{\nu'_1-1} (1-u_2)^{\nu'_2-1} \\ \times |\omega + v(u_2)q|^{\nu_2+\nu_1-1} \frac{\Theta(\omega - v(u_1)q)\Theta(\omega + v(u_2)q)}{(v(u_1) + v(u_2))^{\nu'_2+\nu'_1\nu_2+\nu_1-1}} \\ v(u) = uv_2 + (1-u)v_1$$

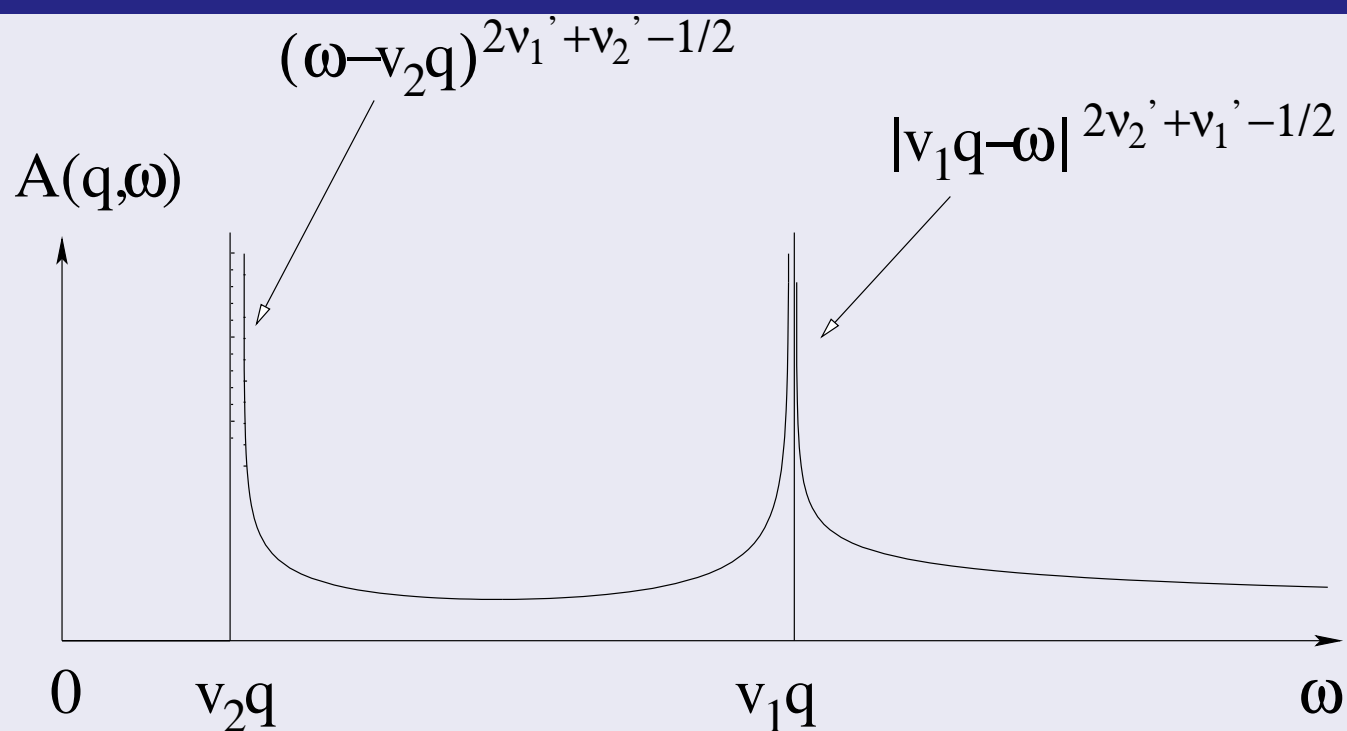
Threshold ($v_2 < v_1$)

$$I(q > 0, \omega < v_2 q) = 0 \Rightarrow A(q > 0, |\omega| < v_2 q) = 0$$

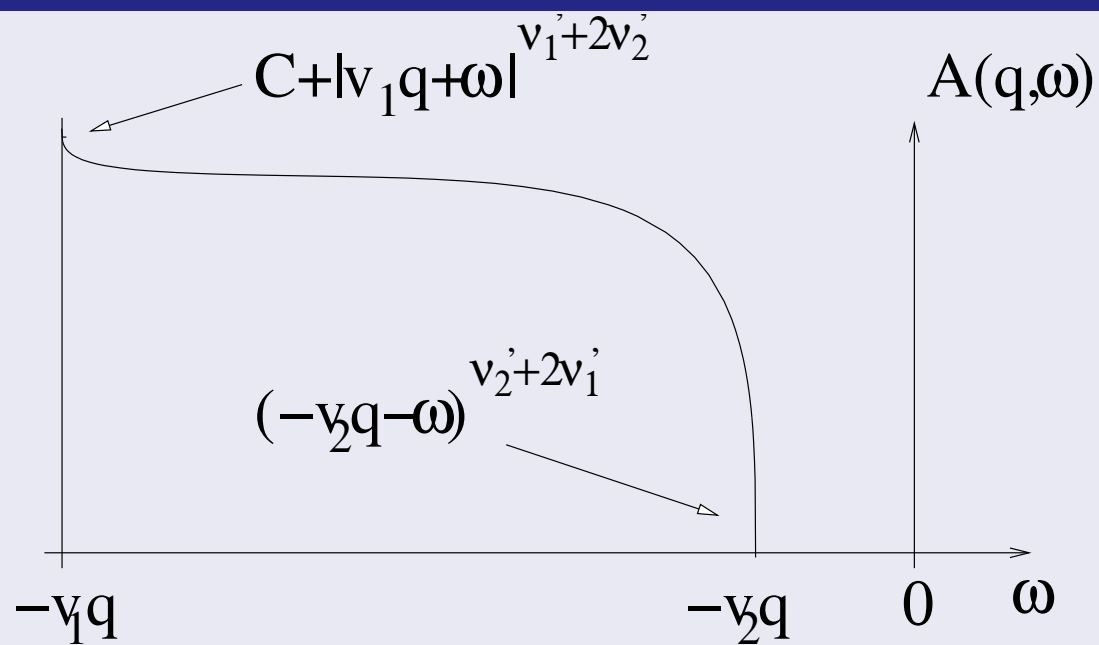
$$0 < v_2 q < \omega < v_1 q$$

$$I(q, \omega) = \frac{\pi^2}{\Gamma(\nu_2)\Gamma(\nu'_1)\Gamma(\nu'_2)\Gamma(\nu_1 + \nu'_1 + \nu'_2)} \times \frac{(\omega - v_2 q)^{\nu'_2 + 2\nu'_1 - 1/2} (v_1 q - \omega)^{2\nu'_2 + \nu'_1 - 1/2}}{(v_1 - v_2)^{2\nu'_1 + 2\nu'_2} (\omega + v_2 q)^{\nu'_2} (\omega + v_2 q)^{\nu'_1}} \times \int_0^1 t^{\nu'_2 - 1} (1 - t)^{\nu'_1 - 1} {}_2F_1\left(2(\nu'_1 + \nu'_2), \nu'_1 + \nu'_2; \nu'_2 + 2\nu'_1 + 1/2; 1 + \frac{\omega - v_1 q}{(v_1 - v_2)q}\right) \times \frac{(\omega + v_2 q)(v_1 + v_2) + (v_1 - v_2)(v_2 q - \omega)t}{(\omega + v_2 q)(\omega + v_1 q)}$$

${}_2F_1$ is Gauss' hypergeometric function.

spectral function for $\omega > 0$: peaks

spectral function for $-v_1q < \omega < 0$: cusps



Conclusion

- 1 Feynman identity used to obtain spectral functions of a 2-band spinless conductor
- 2 expressions in terms of integrals of Gauss hypergeometric functions
- 3 Cusps and peaks easily recovered from general formula.

Perspectives

- 1 More general case with multiple bands (such as 2-leg ladder).
- 2 case of positive temperature