JOVIAN JETS AND VORTICES AS STATISTICAL EQUILIBRIA

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Self Organization of Large Scales of Geophysical Flows



itut non linéaire de **N**ice

Jupiter : The Great Red Spot



Jupiter : A Brown Barge

- Geophysical flows have the property to self-organize at large scale.
 - This is a general property (all planet atmospheres, oceans ...)
- These flows also have a turbulent nature. More quantitatively : the Reynolds' number or the number of exited degrees of freedom $R_e = \frac{UL}{\nu} \simeq 10^{12}$ and $N \simeq 10^{20}$
- Stability : the Jupiter's Great Red Spot exists from more than three century.
- Is this paradoxical ? An explanation ?

A Statistical Explanation

Thermodynamical Analogy

• A variational problem for the curve formed by the jet

 $\min\left\{F_{R}\left[\phi_{R}\right]=2Re_{c}L-2Ru\int_{A_{1}}d\mathbf{r}\,h_{0}(y)+o\left(R\right)\right\}$

• Laplace equation: link between the curvature radius *r* and the free energy difference

 $\frac{e_c}{r} = -u\left(\alpha_1 - h_0\left(y\right)\right)$

• Prediction: the structure is located on extrema of the equivalent topography.



Right: The actual equivalent topography for the GRS and White Ovals (computed using Dowling and Ingersoll (1989) analysis of the ob-

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Typical vortex shape served velocity fields)

Phase Diagram : Jets and Elongated Vortices



E is the energy and *B* measures the asymmetry of the initial PV distribution

- Like for a usual gas example, the statistical effects leads to deterministic behaviors
- The statistical mechanics
 - Phase space : uniform density
 - A huge particle number : entropy
 - Dynamical mixing
- Statistical mechanics for geophysical flows ?

The Quasi-Geostrophic Model

 $\frac{\partial q}{\partial t} + \mathbf{u} \cdot \nabla q = 0$ with $q = -\Delta \psi + \frac{\psi}{R^2} - h(y)$ and $\mathbf{u} = -\mathbf{e}_z \wedge \nabla \psi$

where *q* is the potential vorticity (PV), u the velocity field, ψ the stream function, and *h* the equivalent topography induced by the deep zonal flow.

Conservation laws :

Energy:
$$E = \frac{1}{2} \int_D d\mathbf{r} \left(\mathbf{u}^2 + \frac{\psi^2}{R^2} \right)$$
 Casimirs: $C_f(q) = \int_D d\mathbf{r} f(q)$

Statistical Mechanics of the Quasi-Geostrophic Model



Left : Typical potential vorticity (PV) field, from a numerical simulations with initially only two PV values (red and blue) Let $p(\mathbf{r})$ be the local probability to have one of the two initial PV values (red or blue)

Entropy (proposed by Robert and Sommeria (1991), Miller (1991)):

$S = -\int_{D} [p(\mathbf{r}) \ln p(\mathbf{r}) + (1 - p(\mathbf{r})) \ln(1 - p(\mathbf{r}))] d\mathbf{r}$

The entropy counts the number of states corresponding to a given *p*: the entropy maximum is the most probable state after complete PV mixing

Statistical Equilibria : $\max \{S \mid \text{with } E = E_0 \text{ and } A = A_0\}$

where *A* is the area occupied by one of the potential vorticity levels, *E* is the energy and *S* the entropy (most probable state for a given energy and PV distribution)

• Critical points: a stationnary state given by

$$q = -\Delta \psi + \frac{\psi}{R^2} - h(y) = f_{\alpha,\beta}(\psi)$$



The Shape of QG Equilibria and of Jovian Vortices







Oval BC



A Brown Barge

Statistical Model of the Great Red Spot's Velocity Field





Observation data

Observations : from Voyager data analyses (Dowling and Ingersoll 1994).

• A good quantitative agreement + Rossby deformation radius determination.

White Ovals from Random PV Distributions



Left : Dynamical evolution from a random initial distribution Right : This evolution con-

• Analytical results in the limit of small Rossby deformation radius: $R \rightarrow 0$.



verges to a statistical equilibrium, similar to a White Oval

The Great Red Spot: Coexistence of 2 Thermodynamical

Phases Separated by an Interface (Strong Jet)

Entropy maximization is equivalent to the variational problem:

 $\min \{F_R[\phi] \mid \text{with } A[\phi] \text{ given} \}$ $\text{with } F_R[\phi] = \int_D d\mathbf{r} \left[\frac{R^2(\nabla \phi)^2}{2} + f(\phi) - \frac{R\phi h_0(y)}{2} \right] \text{ and } A[\phi] = \int_D d\mathbf{r} \phi$

This describes a first order phase transition (analogous to a gaz buble in a liquid)



f has two minima

 ϕ thus takes the two values where *f* reaches its minima, in two subdomains (phase separaration) separated by an interface (see right figure)



- With an asymptotic expansion $(R \rightarrow 0)$, we describe the jet
- By analogy with usual thermodynamics, this interface should minimize its length, for a fixed area. The topography will slightly change this picture.

The Brown Barges' Velocity Field



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- F. BOUCHET and J. SOMMERIA, 2002, J. Fluid. Mech., 464 465-207
- F. BOUCHET and T. DUMONT Sub. to Journal of Atmospherical Sciences
- F. BOUCHET, P.H. CHAVANIS and J. SOMMERIA (Shallow Water model)