# Hilbert metric, singular elliptic PDEs and gas dynamics

Denis SERRE UMPA, UMR 5669 CNRS École Normale Supérieure de Lyon, France

> Workshop SISSA / ENS Lyon 5–6 december, 2017

# Team PDEs and Modelling

# Faculty (5)

Emmanuel Grenier (Prof.), Laure Saint-Raymond (Prof., team coordinator), Denis Serre (Prof.), Sergio Simonella (CNRS researcher), Paul Vigneaux (Ass. Prof., HDR).

# PhD, post-doctorate (5)

Arthur Marly, Édouard Ollier, Florian Patout, Valentine Roos, Monika Twarogowska.

## Main themes

**Project INRIA - NUMED.** Complex math. modelling in Medecine: brain stroke, cancer. Multiscale analysis. Wave propagation in complex media. Statistical parameter estimates for PDEs. Numerical simulations.

Collaborations with medical teams.

**Fluid dynamics.** Hydrodynamics. Boundary layers. Statistical Physics. Boltzmann & kinetic equations. Mathematical modelling. Relative entropy method.

# The equation ; 1st formulation

Unknown  $\phi:\Omega\to\mathbb{R}$  ; planar bounded domain.

$$\operatorname{div} \frac{\nabla \phi}{\sqrt{2\phi + |\nabla \phi|^2}} + \frac{2}{\sqrt{2\phi + |\nabla \phi|^2}} = 0. \tag{1}$$

In quasilinear form

$$(2\phi + |\nabla\phi|^2)\Delta\phi - \mathbf{D}^2\phi(\nabla\phi,\nabla\phi) + 4\phi + |\nabla\phi|^2 = 0.$$

The equation ; 2nd formulation

Auxiliary unknown  $w = \sqrt{2\phi}$  (when  $\phi > 0$ )

$$\operatorname{div} \frac{\nabla w}{\sqrt{1 + |\nabla w|^2}} + \frac{2}{w\sqrt{1 + |\nabla w|^2}} = 0.$$
 (2)

or

$$(1 + |\nabla w|^2)\Delta w - D^2 w (\nabla w, \nabla w) + 2 \frac{1 + |\nabla w|^2}{w} = 0.$$

#### Motivation (I)

The PDE is the Euler-Lagrange equation of either

$$L[\phi] = \int \sqrt{2\phi + |\nabla\phi|^2} \frac{dx}{\phi^{3/2}} \quad \text{or} \quad \mathcal{L}[w] = \int \sqrt{1 + |\nabla w|^2} \frac{dx}{w^2}.$$

Interpretation :

The graph  $x_3 = w(x_1, x_2)$  is a non-parametric minimal surface in the Poincaré's half-space

$$\mathbb{H}^{3} = \{ x \in \mathbb{R}^{3} | x_{3} > 0 \}, \qquad ds^{2} = \frac{1}{x_{3}^{2}} (dx_{1}^{2} + dx_{2}^{2} + dx_{3}^{2}),$$

of constant *negative* curvature.

# Motivation (II)

The PDE can be recast in terms of principal curvatures of a rotationally symmetric hypersurface in  $\mathbb{R}^4$  as

$$\kappa_3 = \frac{1}{2}(\kappa_1 + \kappa_2).$$

Again, a parametric equation

$$\sqrt{x_4^2 + x_3^2} = w(x_1, x_2).$$

## Motivation (III)

Euler equations for isentropic gas dynamics

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0,$$
  
$$\partial_t (\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) + \nabla p(\rho) = 0.$$

For self-similar flows (*Riemann Problem*)

$$\rho = \rho\left(\frac{x}{t}\right), \qquad \mathbf{u} = \mathbf{u}\left(\frac{x}{t}\right),$$

the system becomes

div
$$(\rho \mathbf{v}) + 2\rho = 0$$
,  $(\mathbf{v} \cdot \nabla)\mathbf{v} + \mathbf{v} + \frac{1}{\rho}\nabla p = 0$ ,

where  $\mathbf{v}(y) := \mathbf{u}(y) - y$  is the *pseudo-velocity*.

Sub-case : the potential flow

 $\mathbf{v} = \nabla \phi$ 

satisfies Bernoulli's Theorem (*i* the *enthalpy*)

$$\frac{1}{2}|\nabla\phi|^2 + \phi + \iota(\rho) = 0, \qquad \iota'(\rho) = \frac{1}{\rho}p'(\rho),$$

that is

$$\rho = h(2\phi + |\nabla\phi|^2).$$

Whence a  $2^{nd}$ -order PDE

div 
$$\left(h(2\phi + |\nabla\phi|^2)\nabla\phi\right) + 2h(2\phi + |\nabla\phi|^2) = 0.$$
 (3)

The Chaplygin / von Kármán equation of state leads to

$$h(s) = s^{-1/2}.$$

Then (3) reduces to our PDE.

## Type of the equation

• Hyperbolic type if  $\phi < 0$ 

(think to the wave equation).

• Elliptic type if  $\phi > 0$ 

(think to the Laplace equation).

• Degeneracy when  $\phi$  vanishes.

We look for positive solution. Hence we may work in terms of  $w = \sqrt{2\phi}$ .

#### The boundary condition

$$\phi|_{\partial\Omega} = 0$$
 (i.e.  $w|_{\partial\Omega} = 0$ ). (4)

#### Interpretations

(I) The minimal surface is *complete*,

M. Anderson (*Inventiones Math.* 1982) proved the existence of parametric complete minimal surface, using geometric measure theory.

He considers also the non-parametric case, but then his "proof" has serious flaws.

(II) The surface of revolution has no boundary,

(III) The domain  $\Omega$  is the subsonic zone of a flow.

The boundary  $\partial \Omega$  is calculated *a priori* from the Riemann data, by solving plnar wave interactions.

Summing up, our BVP is

$$\begin{split} (2\phi + |\nabla\phi|^2)\Delta\phi - D^2\phi(\nabla\phi,\nabla\phi) + 4\phi + |\nabla\phi|^2 &= 0 \quad \text{in }\Omega, \\ \phi &> 0 \quad \text{in }\Omega, \\ \phi &= 0 \quad \text{on }\partial\Omega. \end{split}$$

The equation is elliptic in  $\Omega$ ; it degenerates along the boundary.

At  $x \in \partial \Omega$  the ellipticity is lost. Its failure occurs in the *tangent* direction (unlike Tricomi equation) : a non generic degeneracy (of *Keldysh* type).

But then,

Where is the data ?

The PDE is homogeneous, as well as the boundary condition ...

The data of the problem is the domain  $\Omega$  !

For the Chaplygin gas, only the supersonic part of the flow is not explicit.

## A necessary condition for existence

Non generic degeneracy is associated with an explicit formula of every derivative at the boundary  $\partial \Omega$  !

For instance

$$\kappa \frac{\partial \phi}{\partial v} = -1,$$
 ( $\kappa$  the curvature).

With  $\phi > 0$  in  $\Omega$  and  $\phi = 0$  on  $\partial \Omega$ , this yields:

**Proposition 1** A necessary condition for the existence of a solution, of class  $C^2$  up to the boundary, is that  $\Omega$  be uniformly convex :

 $\inf \kappa > 0.$ 

#### The convexity is sufficient !

**Theorem 1 (DS, 2009, 2015.)** Let  $\Omega$  be a bounded planar convex domain.

Then there exists a unique positive solution to the BVP,

 $\phi \in \mathcal{C}^{\infty}(\Omega) \cap \mathcal{C}(\overline{\Omega}).$ 

**Difficulties** : – non-linearity, – degeneracy at the boundary, – non-uniform ellipticity (unless  $|\nabla w|$  is bounded, unlikely).

Maximum Principle is the only tool ...

 $\bigcirc$ 

#### **The Maximum Principle**

The equation writes  $N(w, \nabla w, D^2 w) = 0$  with

- (ellipticity)  $S \mapsto N(w, p, S)$  strictly increasing in  $\mathbf{Sym}_2(\mathbb{R})$ ,
- $w \mapsto N(w, p, S)$  decreasing in  $(0, +\infty)$ .

**MP 1** If  $w_+$  is a super-solution ( $N(w_+, \nabla w_+, D^2 w_+) \le 0$  in  $\omega$ ), if  $w_-$  is a sub-solution ( $N(w_-, \nabla w_-, D^2 w_-) \ge 0$ ), and if  $w_- \le w_+$  over  $\partial \omega$ , then

 $w_{-} \leq w_{+}$  dans  $\omega$ .

#### Strategy

- 1. The MP implies uniqueness.
- 2. Comparison with exact solutions  $\phi^{r,m}(x) = \frac{1}{2}(r^2 |x m|^2)$  yield an explicit lower bound  $w_{\min}$  satisfying

 $w_{\min} > 0 \text{ in } \Omega, \qquad w_{\min} = 0 \text{ over } \partial \Omega.$ 

3. Similiar idea, with a little more calculations, yields an explicit upper bound  $w_{\text{max}}$  such that

$$w_{\max} > 0$$
 in  $\Omega$ ,  $w_{\max} = 0$  over  $\partial \Omega$ .

Not the end of the story, because of the lack of uniform ellipticity.

4. Establish a Lipschitz estimate, to ensure the uniform ellipticity.

The most delicate point, since it is known to be only local !

- 5. Apply the regularity theory (*cf* Gilbarg–Trudinger) to get an estimate of  $D^2w$  (or  $D^2\phi$  as well). This ensure the pre-compactness for  $\nabla\phi$ .
- 6. Proceed with a continuation method, with respect to the parameters

div 
$$\frac{\nabla \phi}{\sqrt{2\phi + |\nabla \phi|^2}} + \frac{\mu}{\sqrt{2\phi + |\nabla \phi|^2}} = 0, \qquad \mu \in [0, 2],$$

and

$$\phi|_{\partial\Omega} = \varepsilon > 0, \qquad \varepsilon \to 0^+,$$

(for  $\mu = 0$ , the solution is  $\phi \equiv \epsilon$ ).

#### The Lipschitz estimate

Along  $\partial \Omega$ ,

$$\nabla \phi | = \frac{1}{\kappa}, \qquad |\nabla w| = \infty.$$

**Lemma 1 (2009)** If  $\min_{\partial\Omega} \kappa > 0$  (uniformly convex domain), the Lipschitz estimate of  $\phi$  at the boundary + a (delicate) MP yields a priori

 $\|\nabla\phi\|_{\infty} < C(\min\kappa, \operatorname{diam}\Omega).$ 

 $\diamond$ 

Useless when  $\kappa$  vanishes (flat points) ; one even has  $\|\nabla \phi\|_{\infty} = \infty$ .

Instead, let us use the invariance of the PDE

**under translation :** if *w* is a solution in  $\Omega$ , then  $x \mapsto w(x+a)$  is a solution in  $\Omega - a$ ,

**under rescaling :** likewise,  $x \mapsto \lambda w(\lambda^{-1}x)$  is a solution in  $\lambda \Omega$ .

For  $a, b \in \Omega$ , define

$$p(a|b) = \inf\{\lambda > 0 \mid \Omega - b \subset \lambda(\Omega - a)\} \in [1, +\infty).$$

For  $\lambda \ge p(a|b)$ ,

$$z(x) := \lambda w \left(\frac{x}{\lambda} + a\right)$$

is a solution of the BVP in  $\lambda(\Omega - a)$ . It is thus a super-solution in  $\Omega - b$ . Therefore the MP gives

$$w(x+b) \leq z(x), \quad \forall x \in \Omega - b.$$

Put x = 0:

$$w(b) \le w(a)p(a|b).$$

#### Lipschitz estimate :

 $|\log w(a) - \log w(b)| \le d_{\Omega}(a,b) := \max\{\log p(a|b), \log p(b|a)\},\$ where  $d_{\Omega}$  is the *Thompson* distance over  $\Omega$ .

Equivalence 
$$rac{1}{2}d_H \leq d_\Omega \leq d_H$$
, where  $d_H(a,b):=\log p(a|b) + \log p(b|a).$ 



**Corollary 1** Let *w* be the unique solution *w* of the BVP. Then log *w* is 1-Lipschitz with respect to the Hilbert metric.

 $\diamond$ 

## Summary

- The solution is *a priori* bounded above and below by  $w_{max} > w_{min} > 0$ , both vanishing at the boundary,
- On every compact sub-domain,  $|\nabla w|$  is *a priori* bounded ( $d_H$  is locally equivalent to the Euclidian metric),
- Thus the PDE has uniform ellipticity, away from the boundary,

- Therefore the regularity theory for elliptic equations applies. Whence *a priori* bounds for D<sup>2</sup>*w* on compact sub-domains,
- By Ascoli–Arzela, one has pre-compactness of (approximate) solutions in  $C^1$ . The passage to the limit in the equation is valid.
- The limit of approximate solutions is a solution.
- By the MP, the solution is unique.

Up to technical details, this proves Theorem 1.

Thanks for your attention !