"Drawing sounds, listening to images" The art of time-frequency analysis

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a general framework

« physics »

(laws of Nature, real world applications)

« mathematics »

(models, proofs)

« computer science »

(algorithms)

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the Fourier example



Fourier (1811)

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2 quotations and 2 references

2 quotations from Fourier

- "The profound study of nature is the most fertile source of mathematical discoveries."
- "The [proposed] method does not leave anything vague and indefinite in its solutions; it drives them to their ultimate numerical applications, necessary condition for any research, and without which we would only end up with useless transformations."
- 2 references on Fourier (in French)
 - Jean Dhombres et Jean-Bernard Robert, *Fourier, créateur de la physique mathématique*, Belin 1999.
 - Jean-Pierre Kahane, "Fourier, un mathématicien inspiré par la physique", *Images des Mathématiques*, CNRS 2009.

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analysis/synthesis

Fourier decomposition based on $e_f(t) := \exp\{i2\pi ft\}$

$$x(t) \rightarrow X(f) = \langle x, e_f \rangle, \ s.t. \ x(t) = \int \langle x, e_f \rangle \ e_f(t) \ df$$

- mathematics: "any" waveform is made of the superimposition of a (possibly infinite) number of harmonic modes which are *everlasting*, *undamped* and with a *fixed frequency*
- **physics**: keyrole played by the concept of *frequency* in relation with vibrations and waves
- computer science: further development of efficient algorithms (FFT = 1965) which favoured its practical use

cycles

- physics "of Nature", from macrophysics (celestial mechanics, tides, ...) to microphysics (Quantum Mechanics)
- physics "of engineers" (rotating machines, modal analysis, surveillance of vibrating structures, ...)

W. Thomson (Lord Kelvin), 1876-1878





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lenses

- o diffracted field
- Fourier image in the focal plane
- spatial filtering













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magnitude and phase



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magnitude and phase



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tones

- eigenmodes of cavities
- Helmholtz resonators
- inner ear (cochlea)



Appareil de Kœnig pour l'analyse du timbre des sons. Document Laboratoire de Mécanique et d'Acoustique, CNRS, Marseille.

beyond Fourier

back to the definition

$$egin{aligned} & x(t)
ightarrow X(f) = \langle x, e_f
angle, \ s.t. \ x(t) = \int \langle x, e_f
angle \ e_f(t) \ df, \ e_f(t) &:= \exp\{i 2\pi f t\} \end{aligned}$$

Fourier

- spectrum without any time dependence
- Iocalization on fixed frequencies
- ③ harmonic modes

beyond Fourier

back to the definition

$$egin{aligned} x(t) &
ightarrow X(f) = \langle x, e_f
angle, \ s.t. \ x(t) = \int \langle x, e_f
angle \ e_f(t) \ df, \ e_f(t) &:= \exp\{i 2\pi f t\} \end{aligned}$$

Fourier "+"

- spectrum with time dependence
- Iocalization on varying frequencies
- Inon-harmonic modes

music notation



from sounds to images...



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... and back



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mathematical notes

Issue

"localized modes" \Rightarrow switch to a 2-parameter transformation group that includes time

$$\mathbf{x}(t) \to \mathcal{T}(t,\lambda) = \langle \mathbf{x}, \mathbf{h}_{t,\lambda} \rangle, \ \mathbf{s}.t. \ \mathbf{x}(t) = \iint \langle \mathbf{x}, \mathbf{h}_{s,\lambda} \rangle \ \mathbf{h}_{s,\lambda}(t) \ \mathbf{d}\mu(s,\lambda)$$

1) time-frequency:
$$\lambda = f$$
 and $h_{s,f}(t) = h(t - s) e_f(t)$

 \rightarrow short-time Fourier transform

2) time-scale:
$$\lambda = a$$
 and $h_{s,a}(t) = |a|^{-1/2} h((s-t)/a)$

 \rightarrow wavelet transform

the wavelet connection (\sim 1980-90)

« physics »

vibroseismics for oil exploration

(Morlet)

« mathematics »

CWT, MRA, bases, etc.

(Grossmann, Meyer, Daubechies)



« computer science »

filter banks, fast algorithms

(Mallat, Cohen, Vetterli)

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exclusion principles

« physics »



« mathematics »

any Fourier pair of variables (Weyl, 1927)



« computer science »

time and frequency

(Gabor, 1946 + ...)

classical formulation

"A fast jig on the lowest register of an organ is in fact not so much bad music but no music at all." (N. Wiener in I am a mathematician)

Localization trade-off

based on a second order (variance-type) measure: $\Delta t_x = (\int t^2 |x(t)|^2 dt)^{1/2}$ and $\Delta f_x = (\int f^2 |X(f)|^2 df)^{1/2} \Rightarrow$

$$\Delta t_x \Delta f_x \geq \frac{\|x\|}{4\pi} \ (>0)$$

- no perfect pointwise localization
- variations: same limitation with other spreading measures, e.g., entropy (Hirschman, 1957)
- o common denominator: minimum achieved with Gaussians

extension

no pointwise localization does not mean no localization

Stronger uncertainty relation (Schrödinger, 1935)

$$\Delta t_x \,\Delta f_x \geq \frac{\|\mathbf{x}\|}{4\pi} \sqrt{1 + 16\pi^2 \left(\int t \,\left(\partial_t \arg x(t)\right) \,|\mathbf{x}(t)|^2 \,dt\right)^2}$$

bound achieved for "squeezed states" $\{\exp(\alpha t^2 + \beta t + \gamma)\},\$ with linear "chirps" as a limit when $\operatorname{Re}\{\beta\} = 0$ and $\operatorname{Re}\{\alpha\} \to 0_-$

local methods and localization

 Back to the short-time FT — One defines the local quantity

$$F_x^{(h)}(t,f) = \int x(s) \,\overline{h(s-t)} \, e^{-i2\pi f s} \, ds,$$

where h(t) is some short-time observation window.

- Measurement The representation results from an interaction between the signal and a measurement device (the window h(t)).
- **Trade-off** A short window favors the "resolution" in time at the expense of the "resolution" in frequency, and vice-versa.

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adaptation

- Chirps Adaptation to pulses if h(t) → δ(t) and to tones if h(t) → 1 ⇒ adapting the analysis to arbitrary chirps suggests to make h(t) (locally) depending on the signal.
- Linear chirp In the linear case f_x(t) = f₀ + αt, the equivalent frequency width δf_S of the spectrogram S_x^(h)(t, f) := |F_x^(h)(t, f)|² behaves as:

$$\delta f_{S} \approx \sqrt{\frac{1}{\delta t_{h}^{2}} + \alpha^{2} \, \delta t_{h}^{2}}$$

for a window h(t) with an equivalent time width $\delta t_h \Rightarrow$ minimum for $\delta t_h \approx 1/\sqrt{\alpha}$ (but α **unknown**...).

Fourier notes localization oscillations distributions "chirps" sparsity

self-adaptation and Wigner-Ville distribution

• Matched filtering — If one takes for the window h(t) the time-reversed signal $x_{-}(t) := x(-t)$, one readily gets that $F_x^{(x_{-})}(t, f) = W_x(t/2, f/2)/2$, where

$$W_x(t,f) := \int x(t+ au/2)\,\overline{x(t- au/2)}\,e^{-i2\pi f au}\,d au$$

is the Wigner-Ville Distribution (Wigner, '32; Ville, '48).

• Linear chirps — The WVD perfectly localizes on straight lines of the plane:

$$\mathbf{x}(t) = \exp\{i2\pi(f_0t + \alpha t^2/2)\} \Rightarrow W_{\mathbf{x}}(t, f) = \delta(f - (f_0 + \alpha t)).$$

 Remark — Localization via self-adaptation leads to a quadratic transformation (energy distribution).

interferences

• **Quadratic superposition** — For any pair of signals $\{x(t), y(t)\}$ and coefficients (a, b), one gets

 $W_{ax+by}(t,f) = |a|^2 W_x(t,f) + |b|^2 W_y(t,f) + 2 \operatorname{Re} \{ a \overline{b} W_{x,y}(t,f) \},$ with

$$W_{x,y}(t,f) := \int x(t+ au/2) \,\overline{y(t- au/2)} \, e^{-i2\pi f au} \, d au$$

- Drawback Interferences between disjoint component reduce readability.
- Advantage Inner interferences between coherent components guarantee localization.

interferences

Janssen's formula (Janssen, '81) — It follows from the unitarity of W_x(t, f) that:

$$|W_x(t,f)|^2 = \iint W_x\left(t+\frac{\tau}{2},f+\frac{\xi}{2}\right) W_x\left(t-\frac{\tau}{2},f-\frac{\xi}{2}\right) d\tau d\xi$$

- Geometry (Hlawatsch & F., '85) Contributions located in any two points of the plane plan interfere to create a third contribution
 - 1 midway of the segment joining the two components
 - ② oscillating (positive and negative values) in a direction perpendicular to this segment
 - 3 with a "frequency" proportional to their "time-frequency distance".





















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WV(sum) (N = 3)





WV(sum) (N = 4)







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WV(sum) (N = 6)

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WV(sum) (N = 7)





WV(sum) (N = 8)





WV(sum) (N = 9)

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WV(sum) (N = 10)





WV(sum) (N = 11)





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WV(sum) (N = 12)

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WV(sum) (N = 13)





WV(sum) (N = 14)

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WV(sum) (N = 15)

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WV(sum) (N = 16)

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spectrogram and Wigner-Ville: what else?

Observation

Many quadratic distributions have been proposed in the literature since more than half a century: **none fully extends the notion of spectrum density to the nonstationary case**.

Principle of conditional unicity — **Classes** of quadratic distributions of the form $\rho_x(t, f) = \langle x, \mathbf{K}_{t,f} x \rangle$ can be constructed based on **covariance requirements** :

$$\begin{array}{cccc} \mathbf{x}(t) & \to & \rho_{\mathbf{x}}(t,f) \\ \downarrow & & \downarrow \\ (\mathbf{T}\mathbf{x})(t) & \to & \rho_{\mathbf{T}\mathbf{x}}(t,f) = (\mathbf{\tilde{T}}\rho_{\mathbf{x}})(t,f) \end{array}$$

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classes of quadratic distributions

• Cohen's class — Covariance wrt shifts $(\mathbf{T}_{t_0, f_0} x)(t) = x(t - t_0) \exp\{i2\pi f_0 t\}$ leads to Cohen's class (Cohen, '66) :

$$C_x(t,f) := \iint W_x(s,\xi) \Pi(s-t,\xi-f) \, ds \, d\xi,$$

with $\Pi(t, f)$ "arbitrary" (and to be specified via additional constraints).

• Variations — Other choices possibles, e.g., $(\mathbf{T}_{t_0,f_0}x)(t) = (f/f_0)^{1/2}x(f(t-t_0)/f_0) \rightarrow \text{affine class}$ (Rioul & F, '92), etc.

Cohen's class and smoothing

• **Spectrogram** — Given a low-pass window *h*(*t*), one gets the **smoothing** relation:

$$S_{x}^{(h)}(t,f) := |F_{x}^{(h)}(t,f)|^{2} = \iint W_{x}(s,\xi) W_{h}(s-t,\xi-f) \, ds \, d\xi$$

 From Wigner-Ville to spectrograms — A generalization amounts to choose a smoothing function Π(t, f) allowing for a continuous and separable transition between Wigner-Ville and a spectrogram (smoothed pseudo-Wigner-Ville distributions) :

Wigner - Ville
$$\longrightarrow$$
PWVL \longrightarrow spectrogram $\delta(t) \, \delta(f)$ $g(t) \, H(f)$ $W_h(t, f)$

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extension: data-dependent smoothing



global vs. local

- Global approach The Wigner-Ville Distribution localizes perfectly on straight lines of the plane (linear chirps). One can construct other distributions localizing on more general curves (ex.: Bertrand's distributions adapted to hyperbolic chirps).
- Local approach A different possibility consists in revisiting the smoothing relation defining the spectrogram and in considering localization wrt the instantaneous frequency as it can be measured locally, at the scale of the short-time window ⇒ reassignment.

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reassignment

 Principle — The key idea is (1) to replace the geometrical center of the smoothing time-frequency domain by the center of mass of the WVD over this domain, and (2) to reassign the value of the smoothed distribution to this local centroïd:

$$S_x^{(h)}(t,f)\mapsto \iint S_x^{(h)}(s,\xi)\,\delta\left(t-\hat{t}_x(s,\xi),f-\hat{f}_x(s,\xi)
ight)\,ds\,d\xi.$$

 Remark — Reassignment has been first introduced for the only spectrogram (Kodera *et al.*, '76), but its principle has been further generalized to **any** distribution resulting from the smoothing of a localizable mother-distribution (Auger & F., '95).

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spectrogram = smoothed Wigner



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spreading of auto-terms



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cancelling of cross-terms



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Fourier notes localization oscillations

distributions "chirps" sparsity

reassignment (Kodera et al., 1976, Auger & F., 1995)



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reassignment in action

 Spectrogram — Implicit computation of the local centroïds (Auger & F., '95) :

$$\hat{t}_x(t,f) = t + Re\left\{rac{F_x^{(\mathcal{T}h)}}{F_x^{(h)}}
ight\}(t,f)$$

$$\hat{f}_x(t,f) = f - Im \left\{ \frac{F_x^{(\mathcal{D}h)}}{F_x^{(h)}}
ight\} (t,f),$$

with $(\mathcal{T}h)(t) = t h(t)$ and $(\mathcal{D}h)(t) = (dh/dt)(t)/2\pi$.

 Beyond spectrograms — Possible generalizations to other smoothings (smoothed pseudo-Wigner-Ville, scalogram, etc.).

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reassignment in action

 Gaussian spectrogram — Implicit computation of the local centroïds (Auger, Chassande-Mottin & F., '12) :

$$\hat{t}_x(t,f) = t + \frac{\partial}{\partial t} \log M_x^{(h)}(t,f)$$
$$\hat{f}_x(t,f) = f + \frac{\partial}{\partial f} \log M_x^{(h)}(t,f),$$
with $M_x^{(h)}(t,f) = |F_x^{(h)}(t,f)|.$

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independence wrt window size



an example of comparison



music



echolocation

bat echolocation call + echo



time

bats

frequency

"animal sonar"

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gravitational waves



gravitational wave

Riemann's function



a "compressed sensing" approach



Sparsity

minimizing the ℓ_0 quasi-norm not feasible, but almost optimal solution by **minimizing the** ℓ_1 **norm**

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a "compressed sensing" approach"

Idea (F. & Borgnat, 2008-10)

- (1) choose a domain Ω neighbouring the origin of the AF plane
- ② solve the program

$$\min_{\rho} \|\rho\|_{1} \text{ ; } \mathcal{F}\{\rho\} - A_{x} = \mathbf{0}|_{(\xi,\tau)\in\Omega}$$

3 the exact equality over Ω can be relaxed to

$$\min_{\rho} \|\rho\|_{1} ; \|\mathcal{F}\{\rho\} - A_{x}\|_{2} \leq \epsilon|_{(\xi,\tau)\in\Omega}$$

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a toy example



Wigner



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ambiguity



selection



sparse solution



comparison sparsity vs. reassignment



instantaneous frequency

Aim

model a signal
$$x(t) \in \mathbb{R}$$
 as $x(t) = a_x(t) \cos 2\pi \int^t f_x(s) ds$

- for a given t, "1 equation and 2 unkowns" ⇒ no unique representation
- multiplicity of solutions under constraints
 - global
 - Iocal
 - o non harmonic

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"global" approach (Gabor, 1946; Ville, 1948)

monochromatic wave = circle in the complex plane + constant speed

2
$$x(t) \rightarrow z_x(t) = x(t) + i \mathcal{H}\{x(t)\}$$
 (analytic signal)

modulated "AM-FM" signal: circle \rightarrow "any" loop around the origin of the complex plane + varying speed

3 amplitude : $a_x(t) = |z_x(t)|$ instantaneous frequency : $f_x(t) = \frac{1}{2\pi} \partial_t \arg z_x(t)$

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variation (Equis, Jacquot & F., 2011)



"local" approach (Teager, 1980 ; Kaiser, 1990)

$$I x(t) = a \cos 2\pi f t \Rightarrow \Psi(x) := (\partial_t x)^2 - x \cdot \partial_t^2 x = 4\pi^2 a^2 f^2$$

 $\Psi(x)$ energy operator taking the form $E(x) = x^2[n] - x[n-1]x[n+1]$ in discrete-time

- ② similar local properties when $a \rightarrow a_x(t)$ and $f \rightarrow f_x(t)$
- 3 instantaneous amplitude : $a_x(t) = \Psi(x)/\sqrt{|\Psi(\partial_t x)|}$ instantaneous frequency : $f_x(t) = \frac{1}{2\pi}\sqrt{|\Psi(\partial_t x)/\Psi(x)|}$

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"non harmonic" approach (Huang et al., 1998)



Idea of Empirical Mode Decomposition (EMD) signal = fast oscillation + slow oscillation [& iteration]

- data-driven "fast vs. slow" disentanglement
- "local" analysis based on neighbouring extrema
- oscillation rather than frequency

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algorithm



$$\begin{aligned} x(t) &= c_1(t) + r_1(t) \\ &= c_1(t) + c_2(t) + r_2(t) \\ &= \dots &= \sum_{k=1}^{K} c_k(t) + r_K(t) \end{aligned}$$

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IMF 1; iteration 0

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в

IMF 1; iteration 1





IMF 1; iteration 1











IMF 1; iteration 1



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simulations

Signal

$$x(t) = \underbrace{a_1 \cos\left(2\pi f_1 t\right)}_{x_1(t)} + \underbrace{a_2 \cos\left(2\pi f_2 t + \varphi\right)}_{x_2(t)}, \quad f_1 > f_2$$

Analysis of its EMD

- only the first IMF is computed: if separation, it should be equal to the highest frequency component x₁(t)
- criterion (= 0 if separation) :

$$c\left(\frac{a_2}{a_1}, \frac{f_2}{f_1}, \varphi\right) = \frac{\|IMF_1(t) - x_1(t)\|_{\ell_2}}{\|x_2(t)\|_{\ell_2}}$$

 sampling effects are neglected : f₁, f₂ ≪ f_s, with f_s the sampling frequency

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- Nonlinear behaviour ⇒ dissymmetry of tones disentanglement w.r.t. amplitude ratio a = a₂/a₁, via the sign of a - 1:
 - smooth variation when a < 1 (HF dominant) & no a-dependence
 - abrupt phase transition when a > 1 (LF dominant) & strong a-dependence



 Data-driven separation ⇒ good match to "beating effect" perception ⇒ connection with hearing?

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concluding remarks

- Fourier: 200 years and still alive!
- basic ideas related to decompositions and frequency still central in "modern" approaches, whatever the variations (localized and/or evolutive tones, nonlinear techniques,...)
- time-frequency as a natural language

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back to music notation

Rainer Wehinger' visual listening score created in the 70's to accompany Gyorgy Ligeti's *Artikulation*



http://www.youtube.com/watch?v=71hNl_skTZQ

(thanks to Laurent Chevillard & Gabriel Rilling)

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