

Quantum Monte Carlo for quantum systems on a lattice

(finite-T calculations)

So far:

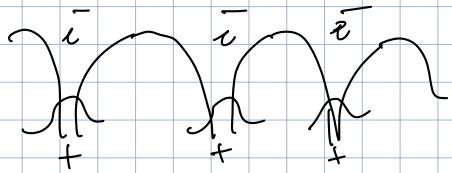
- exact diagonalization + diffusion MC
- variational MC : ground states / time evolutions
- tensor-network techniques: \sim / \sim / \sim / thermodynamics

→ PIMC : thermodynamics of systems in continuous space

⇒ PIMC / QMC for lattice models $\leftarrow \neq \downarrow \xrightarrow{P, X}$

→ quantum spin models : localized electrons / quantum magnetism / quantum simulations based on cold atoms / SC circuit (ensembles of qubits)

→ itinerant identical particles : electrons in a solid on a lattice



- cold atoms in optical lattices
- fermions
- bosons



limitation: sign problem

fermionic systems ($d > 1$)

frustrated quantum magnets

1) PIMC \leftarrow

2) Stochastic series expansion \leftarrow

(transverse field)

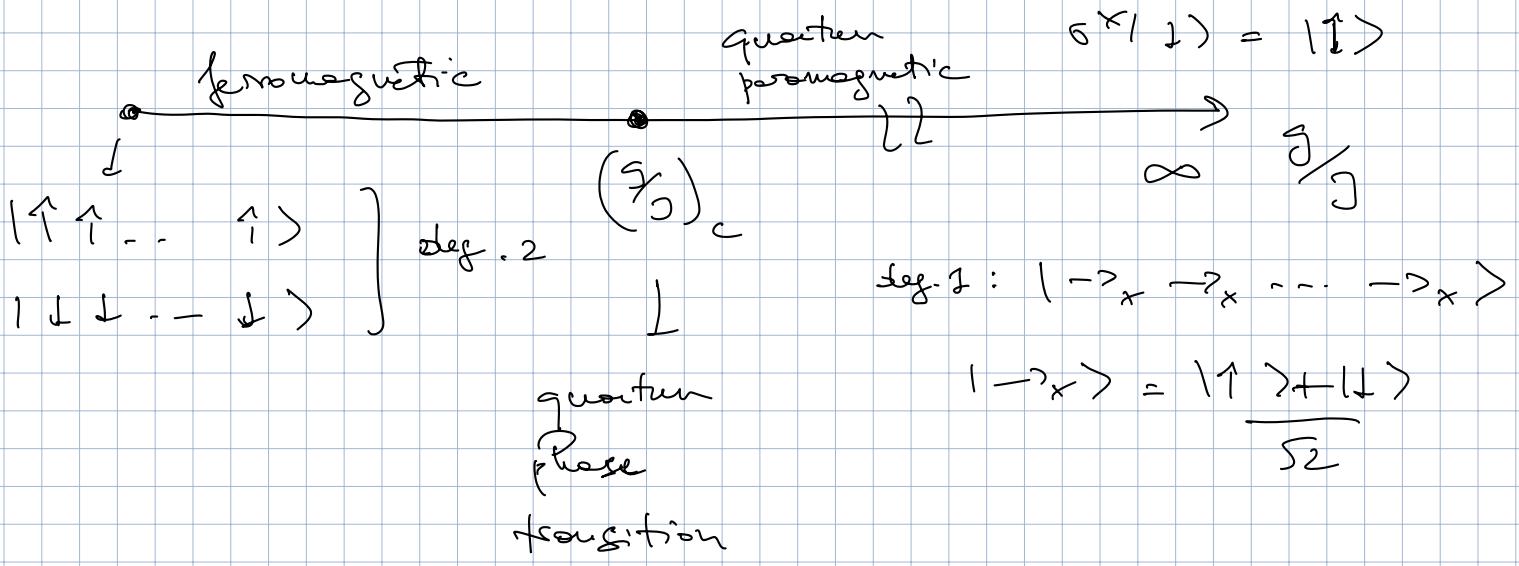
PIMC for quantum Ising model

$$H = - J \left(\sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z + \frac{g}{J} \sum_i \sigma_i^x \right) \quad \text{for interactions}$$

$\sigma_i = \pm 1$

$\sigma^z | \uparrow \rangle = |\uparrow \rangle$
 $\sigma^z | \downarrow \rangle = -|\downarrow \rangle$
 $\sigma^x | \uparrow \rangle = |\downarrow \rangle$
 $\sigma^x | \downarrow \rangle = |\uparrow \rangle$

ground-state phase diagram



$d=1$: exactly solvable \leftrightarrow $d=2$ classical Ising transition

d quantum Ising model \leftrightarrow ($d+1$) classical Ising model

Trotter-Suzuki mapping \rightarrow PI MC

$$Z = \text{Tr} \left(e^{-\beta H} \right)$$

$$H = - J \sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z$$

$$H_g = - g \sum_i \sigma_i^x$$

$$= \text{Tr} \left(e^{-\beta (H_S + H_S)} \right)$$

$$= \lim_{M \rightarrow \infty} \text{Tr} \left(\left(e^{-\beta_{1n} H_S} \wedge e^{-\beta_{1n} H_S} \right)^M \right)$$

$\underbrace{\text{Tr} = \sum_{\vec{\sigma}} |\vec{\sigma}\rangle \langle \vec{\sigma}|}_{\text{decomposition}}$

$$= \lim_{M \rightarrow \infty} \sum_{\vec{\sigma}_1} \langle \vec{\sigma}_1 | \left(\quad \right)^M | \vec{\sigma}_1 \rangle$$

$$|\vec{\sigma}\rangle = (\sigma_1, \sigma_2, \dots, \sigma_N)$$

↓

$$\sigma_i^\dagger |\sigma_i\rangle = \sigma_i |\sigma_i\rangle$$

$$|\vec{\sigma}_k\rangle = (\sigma_{1,k}, \sigma_{2,k}, \dots, \sigma_{N,k})$$

$$\begin{aligned} & \frac{-\beta_{1n} H_S}{e} = \sum_{\vec{\sigma}_n} e^{-\beta_{1n} H_S} (\vec{\sigma}_n) \langle \vec{\sigma}_n | e^{-\beta_{1n} H_S} \\ & = \sum_{\vec{\sigma}_n} \underbrace{e^{-\beta_{1n} \sum_{i,j} \sigma_{i,n} \sigma_{j,n}}} \underbrace{|\vec{\sigma}_n\rangle \langle \vec{\sigma}_n|}_{|\vec{\sigma}_n\rangle \langle \vec{\sigma}_n|} e^{-\beta_{1n} H_S} \end{aligned}$$

$$= \left(\lim_{M \rightarrow \infty} \sum_{\vec{\sigma}_1, \vec{\sigma}_2, \dots, \vec{\sigma}_M} \langle \vec{\sigma}_1 | \right) \underbrace{e^{-\beta_{1n} H_S} (\vec{\sigma}_2) \langle \vec{\sigma}_2 |}_{\text{redacted}} \underbrace{e^{-\beta_{1n} H_S} (\vec{\sigma}_3) \langle \vec{\sigma}_3 |}_{\text{redacted}} \dots \langle \vec{\sigma}_M | \underbrace{e^{-\beta_{1n} H_S} (\vec{\sigma}_1)}_{\text{redacted}}$$

$$= \frac{\beta}{n} \sum_{i=1}^n \sum_{j=1}^n \langle \sigma_{i,n} \sigma_{j,n} \rangle$$

↓

$$\langle \vec{\sigma}_n | \underbrace{e^{\frac{\beta}{n} \sum_i \sigma_i^\dagger}}_{\langle \vec{\sigma}_{n+1} |} \langle \vec{\sigma}_{n+1} |$$

$$= \overline{n}_i \langle \sigma_{i,n} | e^{\frac{\beta}{n} \sigma_i^\dagger} | \vec{\sigma}_{i,n+1} \rangle$$

↙

$$\begin{aligned}
 & \langle \sigma | e^{\frac{\beta g}{M} \sigma^x} |\sigma' \rangle \\
 &= \langle \sigma | \left[\cosh\left(\frac{\beta g}{M}\right) + \sinh\left(\frac{\beta g}{M}\right) \sigma^x \right] |\sigma' \rangle \\
 &= \cosh\left(\frac{\beta g}{M}\right) \left[\delta_{\sigma\sigma'} + \tanh\left(\frac{\beta g}{M}\right) (1 - \delta_{\sigma\sigma'}) \right] \quad \leftarrow \\
 &= \cosh\left(\frac{\beta g}{M}\right) \frac{\partial^L (\sigma\sigma' - 1)}{\partial} \log \left[\tanh\left(\frac{\beta g}{M}\right) \right]
 \end{aligned}$$

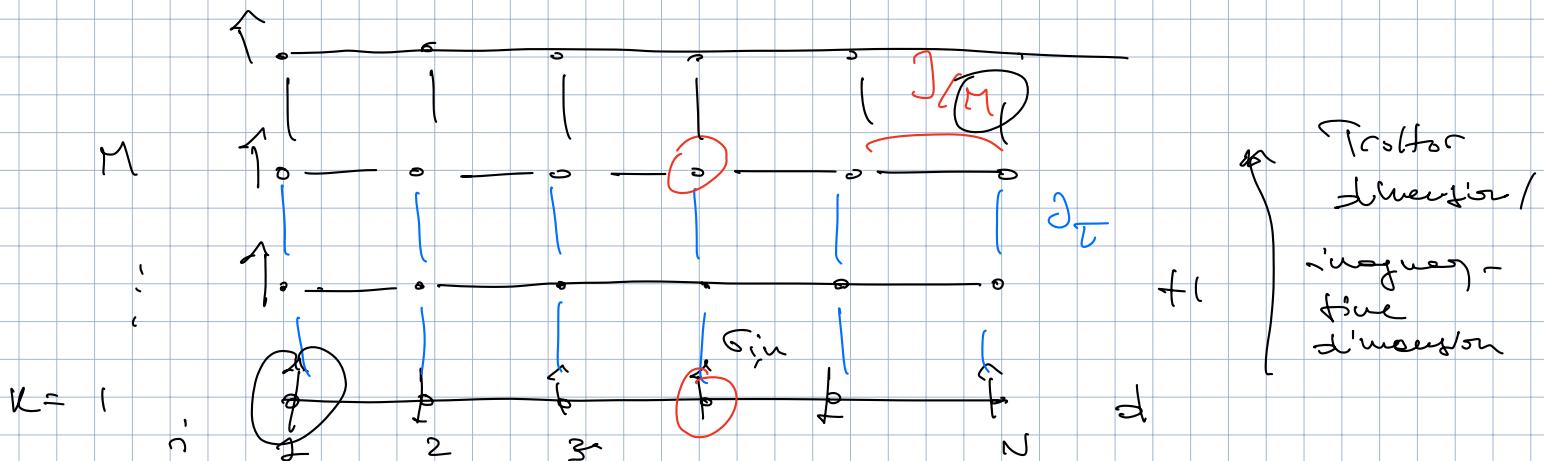
$$\frac{\beta g}{M} \ll 1$$

$$\begin{aligned}
 & \cosh\left(\frac{\beta g}{M}\right) \approx 1 + \frac{1}{2} \left| \log \left[\tanh\left(\frac{\beta g}{M}\right) \right] \right| \quad \text{(Red circle)} \\
 & \approx 1 + \frac{1}{2} \left| \log \left[\tanh\left(\frac{\beta g}{M}\right) \right] \right| (\sigma\sigma' - 1)
 \end{aligned}$$

Putting everything together

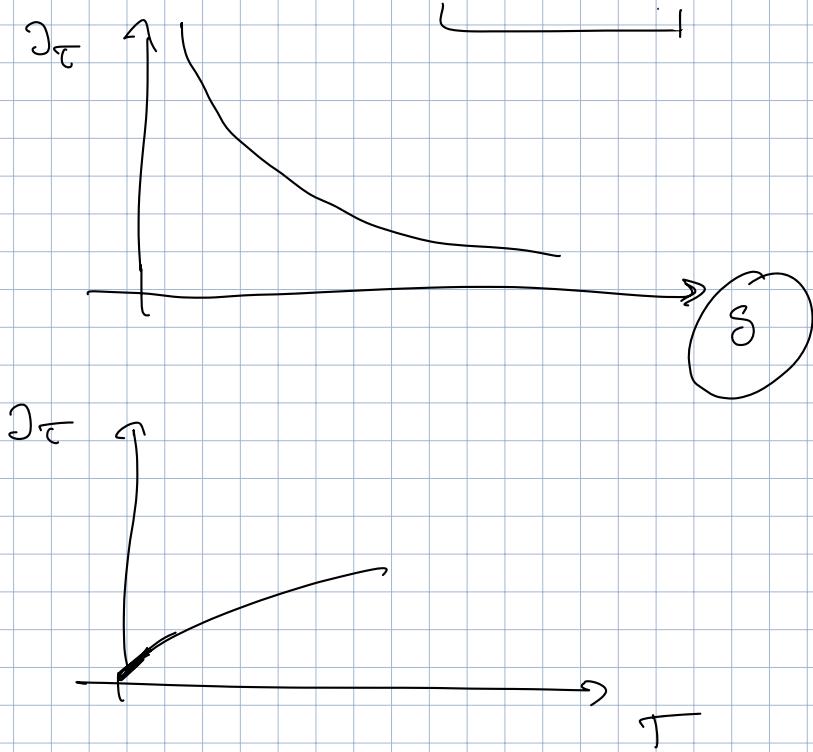
$$Z = \lim_{M \rightarrow \infty} \left[\cosh\left(\frac{\beta g}{M}\right) \right]^M \sum_{\sigma_1, \sigma_2, \dots, \sigma_M} \left(\frac{\beta H_{\text{eff}}(\{\vec{\sigma}_u\})}{M} \right)^M$$

$$H_{\text{eff}}(\{\vec{\sigma}_u\}) = - \left(\frac{J}{M} \right) \sum_{u=1}^M \sum_{j \neq u} (\sigma_{iu} \sigma_{ju} - J_T \sum_i \sum_k (\sigma_{iu} \sigma_{i,u+1} - 1))$$

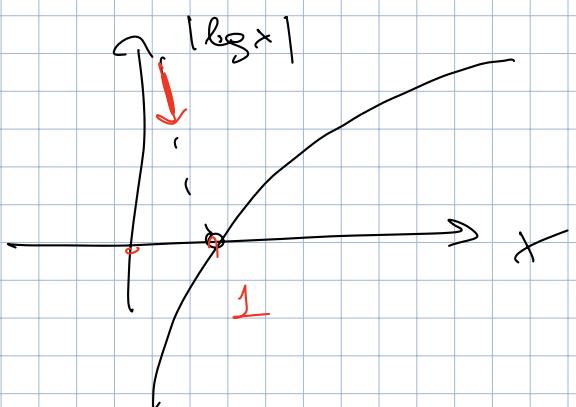


(spatially catastrophic) $(d+1)$ -dimensional Ising model

$$\mathcal{D}_T = \frac{k_B T}{2} \left| \log \left[\tanh \left(\frac{\beta S}{k_B T} \right) \right] \right|$$



$$\frac{\beta S}{k_B T} \ll 1$$



- $\beta \rightarrow 0$
- $T \rightarrow \infty$

$$\mathcal{D}_T = \infty$$

classical limit of
a $d+1$ -dimensional
system

$(d+1)$ -dimensional classical Ising model

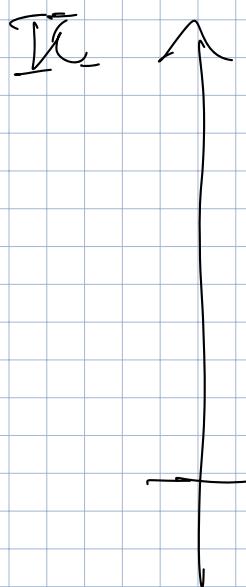
$$\beta J/m$$

$$\Delta H_{ext} = - \overline{J} \sum_{\langle i,j \rangle} \sigma_i \sigma_j - \overline{J} \sum_u \sum_{i,j} (\sigma_i \sigma_j - 1)$$

$$\beta J$$

$$\underline{H}_c = \underline{H}_{c0}$$

spatially isotropic case



FM

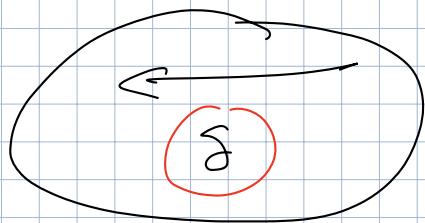
$$\underline{H}_{c0}$$

T

paramagn.

I

$$\frac{\underline{H}_c}{\underline{H}}$$



$$M \rightarrow \infty$$

fermion state physics $\beta \rightarrow \infty$ ($T \rightarrow 0$)

$$\int \delta \tau = \frac{\beta}{M} = \text{const}$$

$\ll 1$

$$\begin{cases} \beta \rightarrow \infty \\ M \rightarrow \infty \end{cases}$$

$$\underline{H}_c$$

increasing β : closed $(d+1)$ -dim Tisir fraction

\equiv quantum PT for the quantum
Tisir model

Gewichtsraster

↓

$$H = - \sum_i h_i \sigma_i^z - \sum_{ij} J_{ij} \sigma_i^z \sigma_j^z$$

$$- \sum_{ijl} Q_{ijl} \sigma_i^z \sigma_j^z \sigma_l^z$$

$$- \sum_{ijklm} P_{ijklm} \sigma_i^z \sigma_j^z \sigma_l^z \sigma_m^z - \sum_i g_i \sigma_i^x$$

lattice gauge theories

+ . . .

↔

PIMC to XXZ model ($S = \frac{1}{2}$)

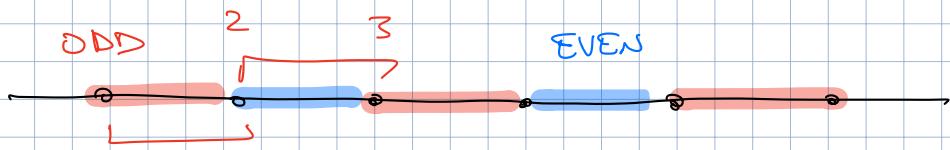
$$H = - \sum_{\langle i,j \rangle} (\xi_i^x \xi_j^x + \xi_i^y \xi_j^y + \Delta \xi_i^z \xi_j^z)$$

$$\begin{array}{c} x \quad x \\ \hline \longrightarrow & \longrightarrow \\ H_{ij}^{xy} & H_{ij}^z \end{array}$$

$$\xi^{\pm} |\sigma\rangle = \frac{\sigma}{2} |\sigma\rangle$$

$$\sigma = \pm 1$$

$$t = \sum_n \langle \sigma \rangle e^{-\beta \sum_{\langle i,j \rangle} (H_{ij}^{xy} + H_{ij}^z)} |\sigma\rangle$$



$$[H_{22}^{xy}, H_{23}^{xy}] \neq 0$$

$$H = \sum_{\langle i,j \rangle \text{ even}} (H_{ij}^{xy} + H_{ij}^z) + \sum_{\langle i,j \rangle \text{ odd}} (H_{ij}^{xy} + H_{ij}^z)$$

$\hookrightarrow H_{ij}$

$$[H_{ij}, H_{lm}] = 0 \quad (i,j) \text{ and } (l,m) \text{ are 5th even bonds}$$

$$\mathcal{Z} =$$

$$\text{Tr} [e^{-\beta H}] = \lim_{M \rightarrow \infty} \left[\left(e^{-\beta_m H_e} - \beta_m H_{odd} \right)^M \right]$$

$$= \lim_{M \rightarrow \infty} \left[\left(\prod_{\langle i,j \rangle \text{ even}} e^{-\beta_m (H_{ij}^{xy} + H_{ij}^z)} \right) \left(\prod_{\langle i,j \rangle \text{ odd}} e^{-\beta_m (H_{ij}^{xy} + H_{ij}^z)} \right)^M \right]$$

$$= \lim_{M \rightarrow \infty} \text{Tr} \left[\prod_{\langle i,j \rangle \text{ even}} \left(e^{-\beta_m H_{ij}^{xy}} - e^{-\beta_m H_{ij}^z} \right) \prod_{\langle i,j \rangle \text{ odd}} \left(e^{-\beta_m H_{ij}^{xy}} - e^{-\beta_m H_{ij}^z} \right)^M \right]$$

$$\sum_{\vec{\sigma}} (\vec{\sigma}) \langle \vec{\sigma} \rangle$$

$$\sum_{\vec{\sigma}} (\vec{\sigma}) \langle \vec{\sigma} \rangle$$

$$= \lim_{M \rightarrow \infty} \sum_{\{\vec{\sigma}_n\}} \langle \vec{\sigma}_1 \rangle \prod_{\langle i,j \rangle \text{ even}} e^{-\beta_m H_{ij}^{xy}} | \vec{\sigma}_2 \rangle \langle \vec{\sigma}_2 | \prod_{\langle i,j \rangle \text{ odd}} e^{-\beta_m H_{ij}^{xy}} | \vec{\sigma}_3 \rangle \langle \vec{\sigma}_3 | \prod_{\langle i,j \rangle \text{ even}} e^{-\beta_m H_{ij}^{xy}} | \vec{\sigma}_4 \rangle \langle \vec{\sigma}_4 | \prod_{\langle i,j \rangle \text{ odd}} e^{-\beta_m H_{ij}^{xy}}$$

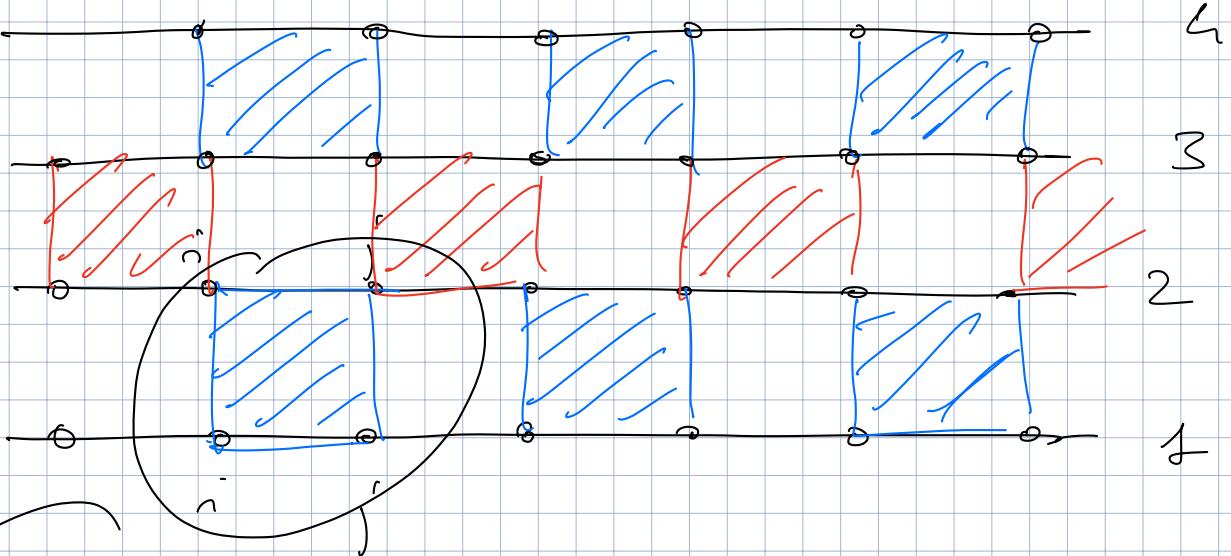
$$\langle \vec{\sigma}_3 | \prod_{\langle i,j \rangle \text{ even}} e^{-\beta_m H_{ij}^{xy}} | \vec{\sigma}_4 \rangle$$

$$\langle \vec{\sigma}_{2n} | \prod_{\langle i,j \rangle \text{ odd}} e^{-\beta_m H_{ij}^{xy}} | \vec{\sigma}_2 \rangle$$

$$e^{-\frac{1}{4m} \sum_{\langle i,j \rangle} \sum_{n=1}^{2m} \vec{\sigma}_{i+n} \cdot \vec{\sigma}_{j+n}}$$

$$\langle \vec{\sigma}_n | e^{-\beta_m \sum_{\langle i,j \rangle \text{ even}} H_{ij}^{xy}} | \vec{\sigma}_{n+1} \rangle$$

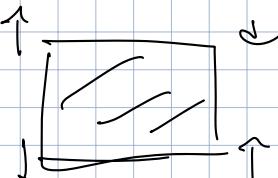
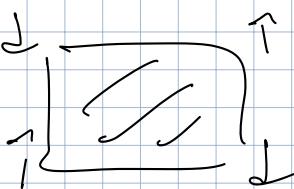
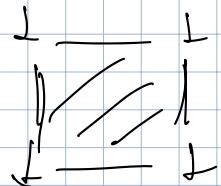
$$= \prod_{\langle i,j \rangle \text{ even}} \langle \sigma_i \sigma_j \rangle e^{-\beta_m H_{ij}^{xy}} | \vec{\sigma}_{n+1} \cdot \vec{\sigma}_{j+n} \rangle$$



$$\langle \sigma_{iu} \sigma_{ju} | e^{-\beta_m H_{ij}^{xy}} | \sigma_{i,u+n} \sigma_{j,u+n} \rangle = w(\sigma_{iu}, \sigma_{ju}; \sigma_{i,u+n}, \sigma_{j,u+n})$$

$$= \begin{pmatrix} \uparrow\uparrow & 1 & 0 & 0 & 0 \\ \uparrow\downarrow & 0 & \cosh\left(\frac{\beta_j}{2m}\right) & \sin\left(\frac{\beta_j}{2m}\right) & 0 \\ \downarrow\uparrow & 0 & \sinh\left(\frac{\beta_j}{2m}\right) & \cosh\left(\frac{\beta_j}{2m}\right) & 0 \\ \downarrow\downarrow & 0 & 0 & 0 & 1 \end{pmatrix}$$

$$w \neq 0$$



6- vertex model

quarter XXZ model
in 2 dimensions

classical

6-vertex model
in $(d+1)$ dimensions

$$d=1$$

FM



-1

dyadic

Correlations

1



Ising - Thirring

QFT