

Equation de Schrödinger indépendante du temps en représentation x

1D=1

①

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \psi_{\epsilon}(x) = \epsilon \psi_{\epsilon}(x)$$

$V(x) = 0$

$$-\frac{d^2}{dx^2} \psi = 2m\epsilon \psi \quad \psi = N e^{\pm i \sqrt{\frac{2m\epsilon}{\hbar^2}} x}$$

$$\psi_{\epsilon}(x) = N e^{ikx}$$

solution générale

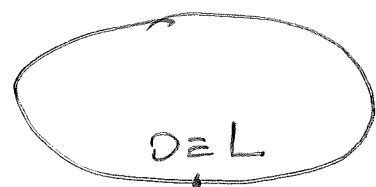
$$\psi(x) = A e^{ikx} + B e^{-ikx}$$

$$k = \pm \sqrt{\frac{2m\epsilon}{\hbar^2}}$$

$$\epsilon = \frac{\hbar^2 k^2}{2m}$$

$$|\psi| \sim \frac{1}{\sqrt{L}}$$

Conditions aux limites périodiques



$$\psi_n(x) = \frac{e^{ik_n x}}{\sqrt{L}}$$

$$k_n = \frac{2\pi}{L} n$$

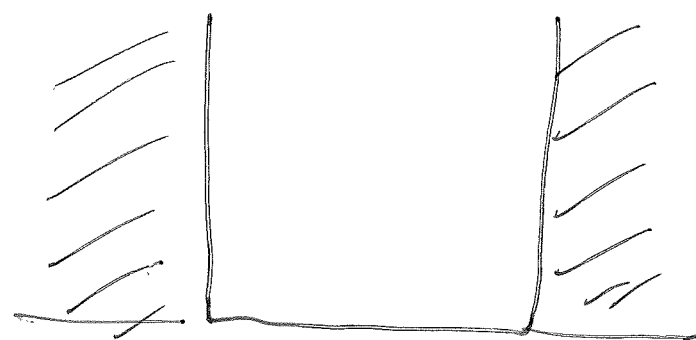
$$E_n = \frac{\hbar^2}{2m} \left(\frac{2\pi}{L} n \right)^2$$

$$n \in \mathbb{Z}$$

$$\begin{cases} \psi_k(x) = \frac{1}{\sqrt{2\pi}} e^{ikx} \\ \phi_p(x) = \frac{1}{\sqrt{2\pi\hbar}} e^{i\frac{p}{\hbar}x} \\ \text{Normalisée à } \delta \end{cases}$$

$$\begin{aligned} & \text{---} & n=2, -2 \\ & \text{---} & n=1, -1 \\ & \text{---} & n=0 \end{aligned}$$

Condition aux limites ouvertes / libres



$$V(x) = \begin{cases} 0 & 0 < x < L \\ \infty & x < 0 \text{ or } x > L \end{cases}$$

$$\psi(0) = \psi(L) = 0$$

$$\psi(0) = 0 \quad A+B=0 \quad A=-B$$

$$\psi_k(x) = 2iA \sin(kx)$$

$$\psi(L) = 0$$

$$\sin(kL) = 0 \rightarrow$$

$$kL = n\pi \quad \left[k_n = \frac{n\pi}{L} \right]$$

$$n \in \mathbb{N}^*$$

$$n=1, 2, \dots$$

$$\psi_n(x) = \underbrace{2iA}_N \sin\left(\frac{n\pi}{L}x\right) = 2iA \frac{e^{i\frac{n\pi}{L}x} - e^{-i\frac{n\pi}{L}x}}{2i}$$

(2)

$$N^{-1} \left[\int_0^L dx \sin^2\left(\frac{n\pi}{L}x\right) \right]^{1/2} = \frac{e^{i\frac{n\pi}{L}x}}{\sqrt{L}} - \frac{e^{-i\frac{n\pi}{L}x}}{\sqrt{L}}$$

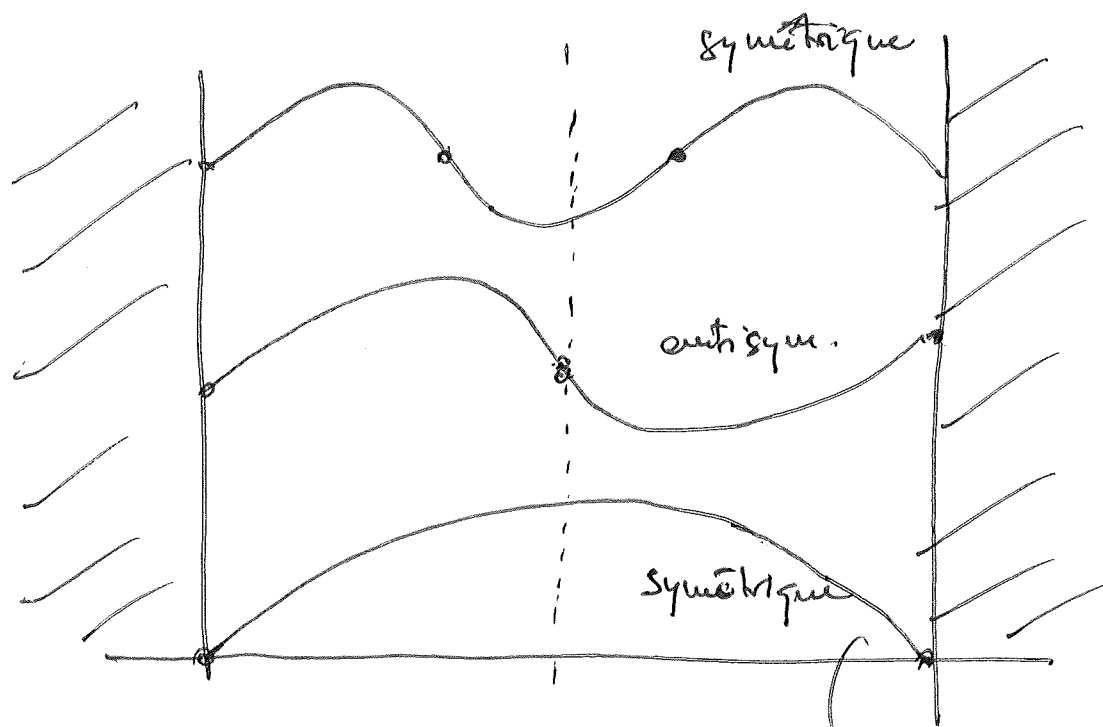
$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi}{L}x\right)$$

$$E_n = \frac{\hbar^2}{2m} \left(\frac{n\pi}{L}\right)^2$$

$$= \frac{1}{\sqrt{2}} \left[\frac{e^{i\frac{n\pi}{L}x}}{\sqrt{L}} - \frac{e^{-i\frac{n\pi}{L}x}}{\sqrt{L}} \right]$$

$$= 2i \frac{1}{\sqrt{2L}} \sin\left(\frac{n\pi}{L}x\right) = \cancel{2i} \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi}{L}x\right)$$

$$2A = \sqrt{\frac{2}{L}}$$



de nœuds = $n-1$

$$n=3$$

$$n=2 \quad \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi}{L}x\right)$$

$$n=1 \quad \sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x\right)$$

valeurs propres
non dégénérées

$$E_{n=1} = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 \neq 0$$

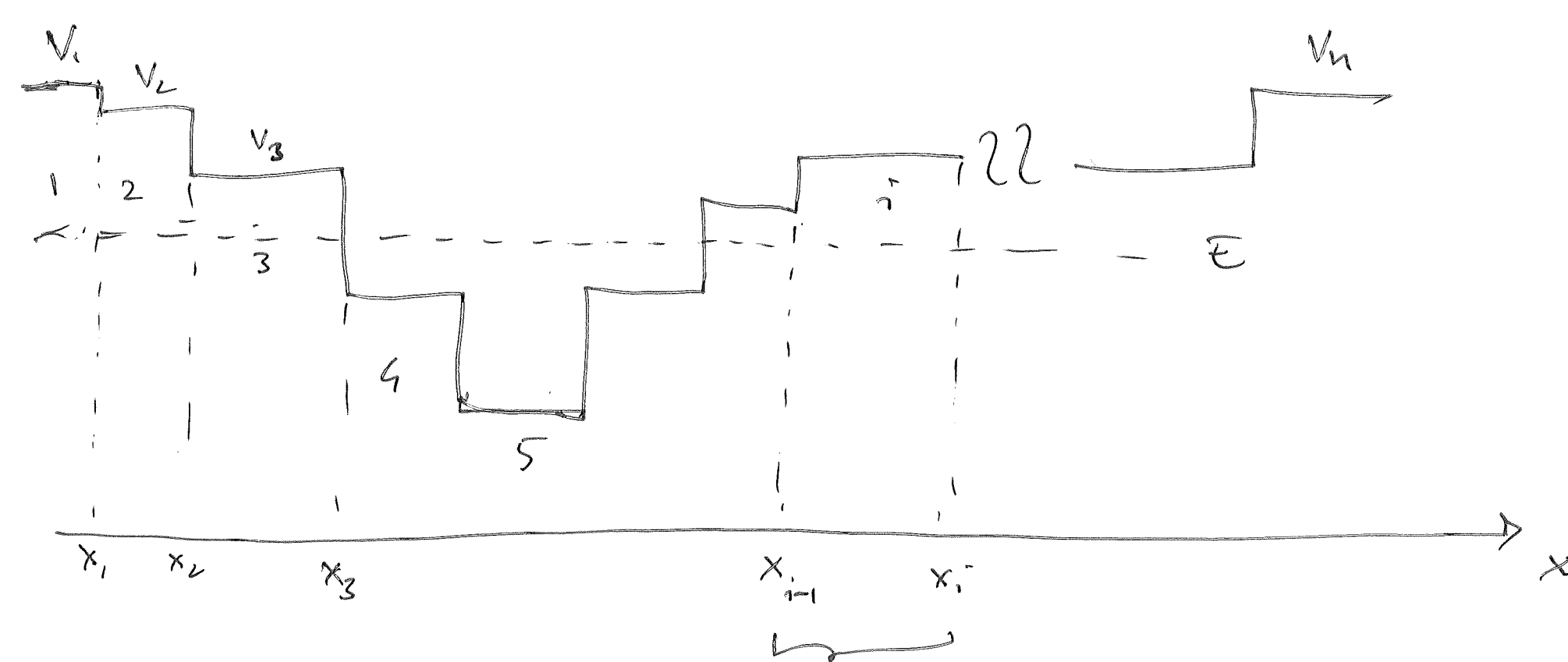
énergie de point zéro

$$\Delta x = L$$

$$\Delta p \sim \frac{\hbar}{L}$$

$$E \approx \frac{p^2}{2m} \approx \frac{(\Delta p)^2}{2m} = \frac{\hbar^2}{2m} \left(\frac{1}{L}\right)^2$$

Solution générale de l'équation de Schrödinger pour un ^{puits de} V potentiel constant par morceaux (3)



$$V(x) = \begin{cases} V_1 & x \leq x_1 \\ V_2 & x_1 < x \leq x_2 \\ \vdots & \vdots \end{cases}$$

n morceaux

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi = (E - V_i) \psi$$

n morceaux

2n variables A_i, B_i

$$\psi = A_i e^{i k_i x} + B_i e^{-i k_i x}$$

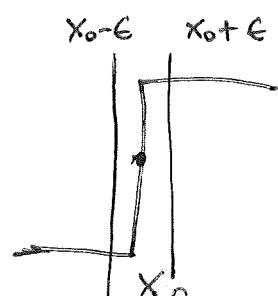
$$= A_i e^{-q_i x} + B_i e^{q_i x}$$

$$k_i = \begin{cases} \sqrt{\frac{2m(E - V_i)}{\hbar^2}} & E > V_i \\ i \sqrt{\frac{2m(E - V_i)}{\hbar^2}} = i q_i & E < V_i \end{cases}$$

Conditions de passage

n-1 points de raccord

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi = [E - V(x)] \psi(x)$$



$$\int_{x_0-E}^{x_0+E} dx$$

$$\frac{d^2 \psi}{dx^2} = \psi'(x_0+E) - \psi'(x_0-E)$$

$$= \int_{x_0-E}^{x_0+E} dx \left[-\frac{2m}{\hbar^2} [E - V(x)] \psi(x) \right] \xrightarrow{E \rightarrow 0} 0$$

$\psi(x)$ normalisée

$$\int |\psi(x)|^2 < \infty \quad (4)$$

$$\psi(x) \sim \frac{1}{x^\alpha}$$

deux conditions

$$\begin{cases} \psi'(x_0^-) = \psi'(x_0^+) \\ \psi(x_0^-) = \psi(x_0^+) \end{cases}$$

$$\psi'(x_0^-) = \psi'(x_0^+)$$

derivée première est continue



$$\psi'(x_0) \text{ existe} \Rightarrow$$

$\psi(x)$ est continue en x_0

$$\psi(x_0^-) = \psi(x_0^+)$$

$$\begin{cases} 2(n-1) = 2n-2 \\ 2n \text{ variables} \end{cases} \quad \begin{array}{l} \text{conditions de passage} + 1 \text{ condition de} \\ \text{normalisation} \\ \text{(quand elle est applicable)} \end{array}$$

Cas spécifiques

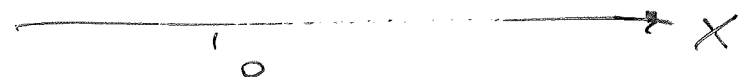
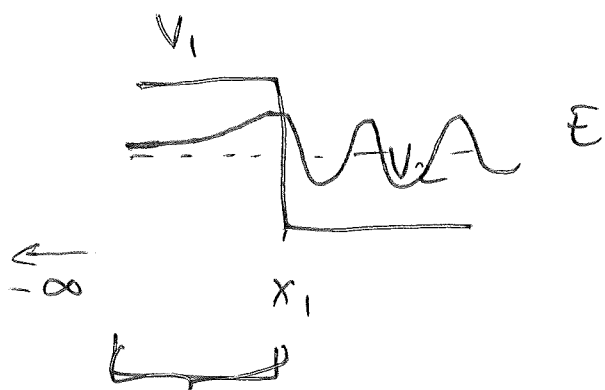
$$1) \text{ ETAT LIÉ : } \begin{array}{l} E < V_1 \\ E < V_n \end{array}$$

$$\begin{cases} 2n-2 \text{ variables} \\ \text{pour } 2n-1 \text{ conditions} \end{cases}$$

condition sur l'énergie E

$\Rightarrow$ spectre discret

spectre non dégénéré

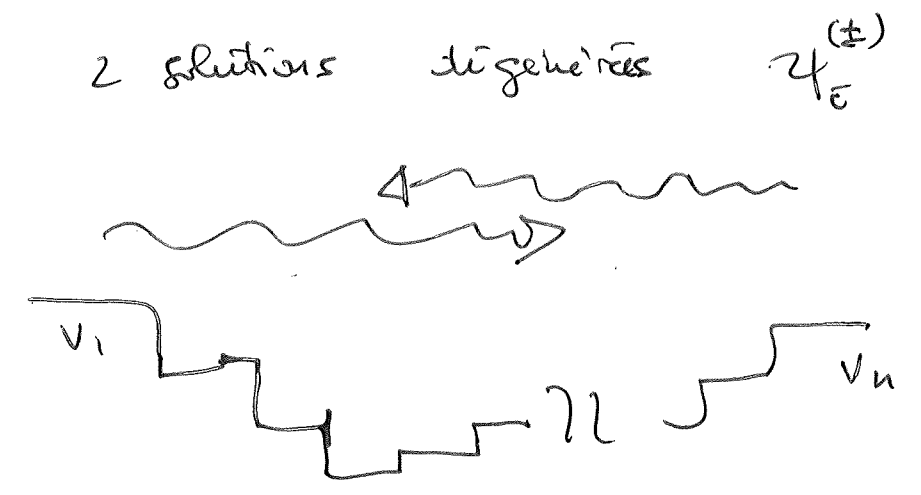


$$A_1 e^{-\kappa_1 x} + B_1 e^{\kappa_1 x}$$

2) $\bar{E}$ est libre $E > V_1, E > V_n$ (5)

$2n-2$ conditions de passage
 $2n$ variables

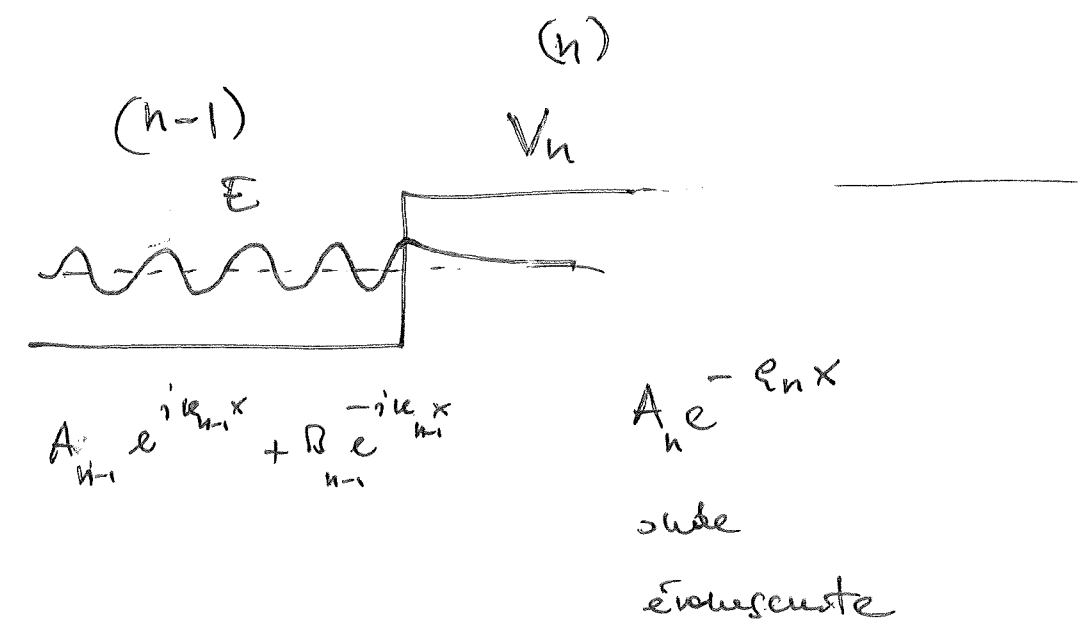
$$\psi = A^{(+)} \psi^{(+)} + A^{(-)} \psi^{(-)}$$



Avec les ψ_E $E > V_1, V_n$
 on construit des paquets d'ondes pour avoir des états physique.

3) $\bar{E}$ est semi-libre

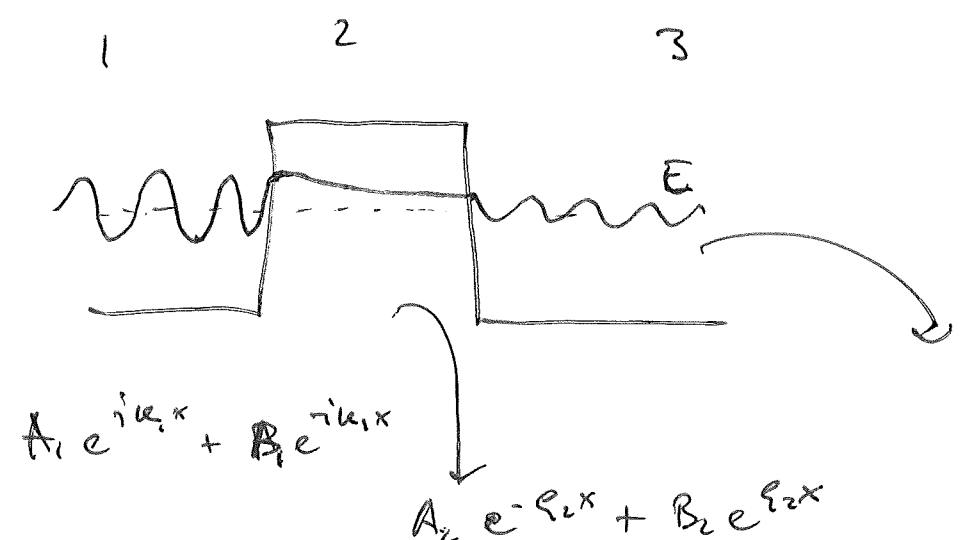
$E > V_1$
 $E < V_n$



$2n-2$ conditions
 $2n-1$ variables

solution non normalisable

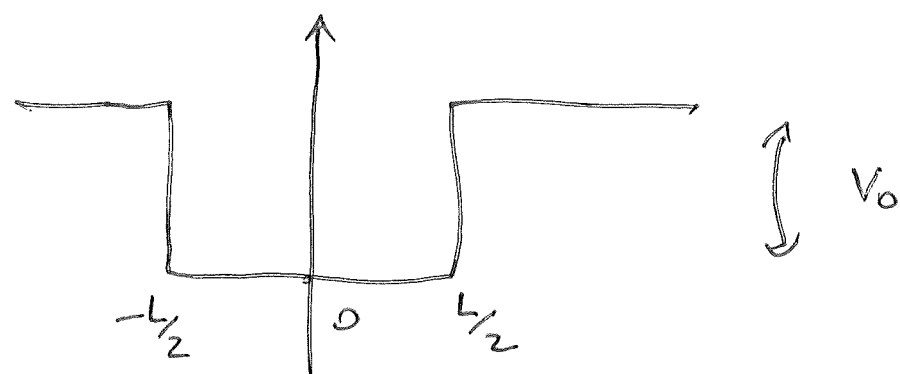
Exemple
 barrière fine



EFFET TUNNEL

$$A_3 e^{i k_3 x} + B_3 e^{-i k_3 x}$$

Puit de potentiel fini



① ② ③

Recherche des états liés

$$E < V_0$$

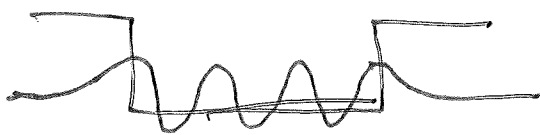
- 1) $A_1 e^{+\xi x}$
- 2) $A_2 e^{i k x} + B_2 e^{-i k x}$
- 3) $A_3 e^{-\xi x}$

$$\xi = \frac{\sqrt{2m |E - V_0|}}{\hbar}$$

$$k = \frac{\sqrt{2m E}}{\hbar}$$

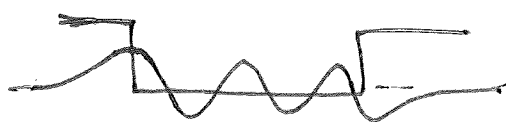
Symétrique

- 1) $A e^{\xi x}$
- 2) $B \cos(kx)$
- 3) $A e^{-\xi x}$



Anti-symétrique

- 1) $A' e^{\xi x}$
- 2) $B' \sin(kx)$
- 3) $-A' e^{-\xi x}$



$$\langle \psi | \hat{P} | \phi \rangle = \int dx \cdot \psi^* \hat{P} \phi = (\langle \phi | \hat{P} | \psi \rangle)^* \quad (6)$$

Symétrique par inversion

Si $\psi(x)$ est solution $\Rightarrow \psi(-x)$ est aussi solution

$$\psi(x) = e^{i\phi} \psi(-x)$$

$$\hat{P} \psi(x) = \psi(-x)$$

opérateur de parité

$\hat{P}$ opérateur hermitique

$$\hat{P} \psi(x) = \psi(-x) = e^{-i\phi} \psi(x)$$

$e^{-i\phi}$ valeur propre de $\hat{P} \rightarrow e^{i\phi} \in \mathbb{R}$
 $\phi = 0, \pi$

$$\psi(x) = \pm \psi(-x)$$

$\psi(x)$ a une parité définie

Passage entre (1) et (2)

symétrique

$$\psi(-\frac{L}{2}^-) = \psi(-\frac{L}{2}^+)$$

$$A e^{-\kappa \frac{L}{2}} = B \cos(\kappa \frac{L}{2})$$

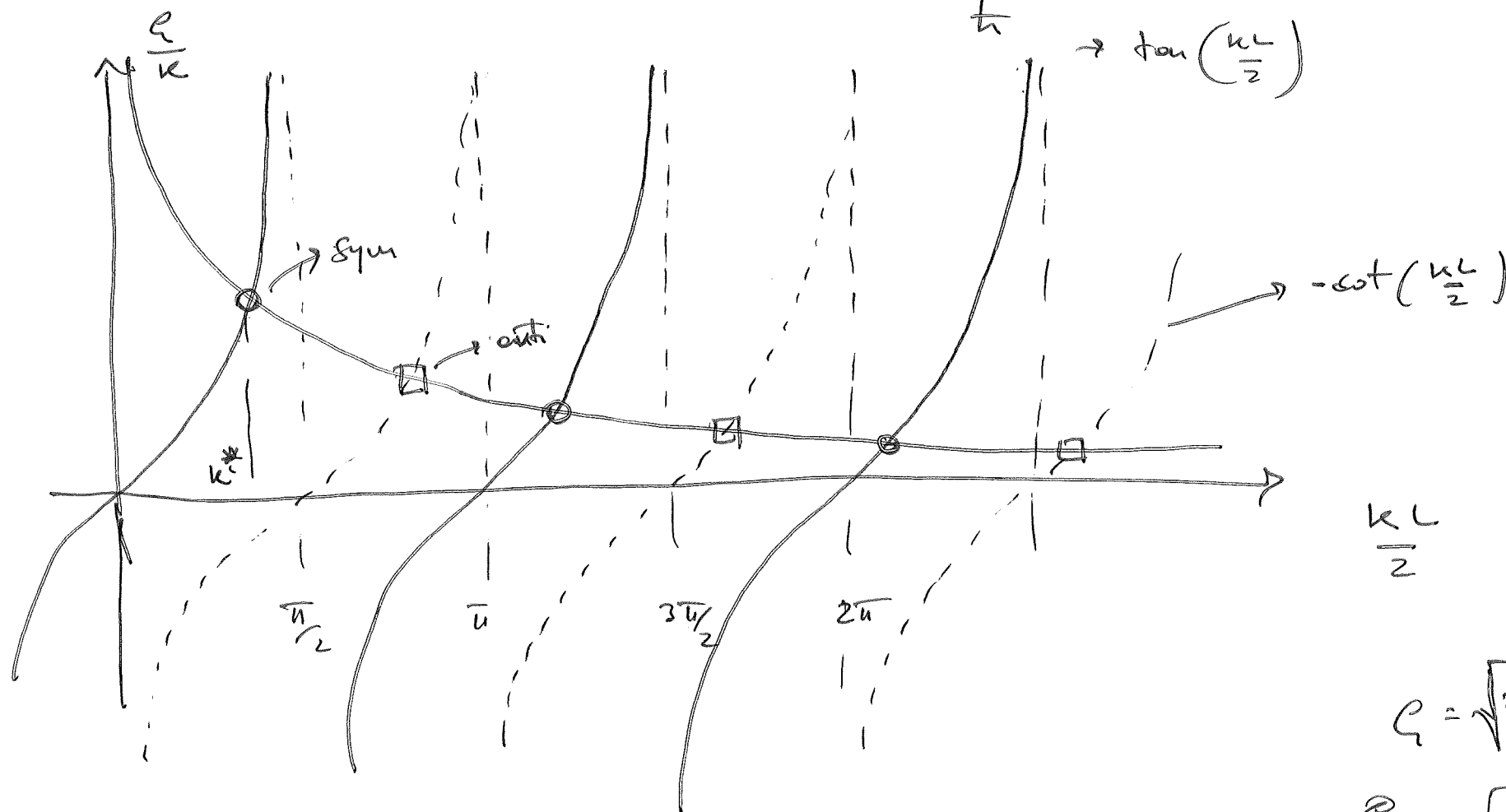
$$\psi'(\quad) = \psi'(\quad)$$

$$A \kappa e^{-\kappa \frac{L}{2}} = -B \kappa \sin(\kappa \frac{L}{2})$$

$$\frac{\kappa}{\kappa} = \tan(\kappa \frac{L}{2})$$

condition sur $\kappa \Rightarrow$

$$\kappa = \frac{\sqrt{2mE}}{\hbar} \rightarrow \tan(\kappa \frac{L}{2})$$



antisymétrique

$$A' e^{-\kappa \frac{L}{2}} = -B' \sin(\kappa \frac{L}{2})$$

$$A' \kappa e^{-\kappa \frac{L}{2}} = -B' \kappa \cos(\kappa \frac{L}{2})$$

$$\frac{\kappa}{\kappa} = -\cot(\kappa \frac{L}{2})$$

$$E < V_0$$

$$\kappa \leq \kappa_{\max} = \frac{\sqrt{2mV_0}}{\hbar}$$

Il en existe un état lié ?

$$\frac{\kappa L}{2}$$

$$\frac{\kappa}{\kappa} = \tan(\kappa \frac{L}{2})$$

$$\kappa \leq \kappa_{\max}$$

$$\kappa = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}} = \sqrt{\kappa_{\max}^2 - \kappa^2}$$

$$\frac{\kappa}{\kappa} = \sqrt{\frac{\kappa_{\max}^2}{\kappa^2} - 1} = \tan(\kappa \frac{L}{2})$$

$$\frac{\kappa_{\max}^2}{\kappa^2} - 1 = \tan^2(\kappa \frac{L}{2})$$

$$\frac{k_{\max}^2}{k^2} = \tan^2(kL) + 1$$

(8)

$$k_{\max}^2 = (\tan^2(kL) + 1) k^2 \geq k^2 \quad k_{\max} \geq k$$

∃ un état lié symétrique pour toute valeur
de V_0 .