

Particule dans un potentiel central

①

$$\hat{H} = \frac{\hat{p}^2}{2\mu} + \hat{V}(r)$$

représentation $r, \vartheta, \phi \rightarrow -\frac{\hbar^2}{2\mu r} \frac{\partial^2}{\partial r^2} r(\cdot) + V(r) + \frac{\hat{L}^2}{2\mu r^2}$

$\hat{L}^2$ est un opérateur qui dépend seulement de ϑ et ϕ

$$[\hat{H}, \hat{L}^2] = 0$$

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$$[\hat{L}^2, \hat{L}] = 0$$

$$\psi(\vec{r}) = \psi(r, \vartheta, \phi) = \underline{R_{nl}}(r) Y_{lm}(\vartheta, \phi)$$

$$\left[-\frac{\hbar^2}{2\mu r} \frac{\partial^2}{\partial r^2} r + V(r) + \frac{\hat{L}^2}{2\mu r^2} \right] R_{nl}(r) Y_{lm}(\vartheta, \phi) = E_{nl} Y_{lm}(\vartheta, \phi) R_{nl}(r)$$

$$\begin{cases} \hat{L}^2 Y_{lm}(\vartheta, \phi) = \hbar^2 l(l+1) Y_{lm}(\vartheta, \phi) \\ \hat{L}^2 Y_{lm}(\vartheta, \phi) = \hbar^2 m Y_{lm}(\vartheta, \phi) \end{cases}$$

$$\left[-\frac{\hbar^2}{2\mu r} \frac{d^2}{dr^2} r + V(r) + \frac{\hbar^2 l(l+1)}{2\mu r^2} \right] R_{nl}(r) = E_{nl} R_{nl}(r)$$

$$R_{nl}(r) = \frac{u_{nl}(r)}{r}$$

$$-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} u_{nl} + V(r) \cancel{\frac{u_{nl}}{r}} + \frac{\hbar^2 l(l+1)}{2\mu r^2} \cancel{\frac{u_{nl}}{r}} = E_{nl} \cancel{\frac{u_{nl}}{r}}$$

$$\left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + V(r) + \frac{\hbar^2 l(l+1)}{2\mu r^2} \right] u_{nl} = E u_{nl}$$

$\underbrace{\quad}_{V_{eff}(r)} \rightarrow$ barrière centrifuge

Hypothèse sur $V(r)$:

$$\begin{cases} V(r) \text{ ne diverge pas quand } r \rightarrow 0 \\ \text{ou diverge moins vite que } \frac{1}{r^2} \\ V(r) < 0 \end{cases}$$

$$\boxed{l \geq 0}$$

$$L = I\omega$$

$$E = \frac{1}{2} I \omega^2$$

$$\underline{r \rightarrow 0} \quad \left(-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + \frac{\hbar^2 l(l+1)}{2\mu r^2} - E \right) u_{nl} \approx 0$$

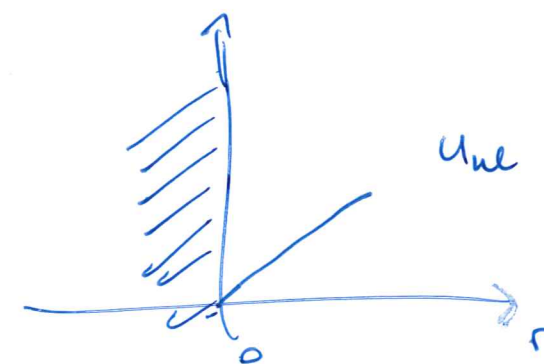
$$u_{nl} \sim r^\alpha$$

$$-\frac{\hbar^2}{2\mu} \left[\alpha(\alpha-1) r^{\alpha-2} - l(l+1) r^{\alpha-2} \right] = 0$$

$$\alpha(\alpha-1) = l(l+1) \quad \begin{cases} \alpha = l+1 \\ \alpha = -l \end{cases}$$

$$\boxed{u_{nl} \sim r^{l+1} \text{ pour } r \rightarrow 0}$$

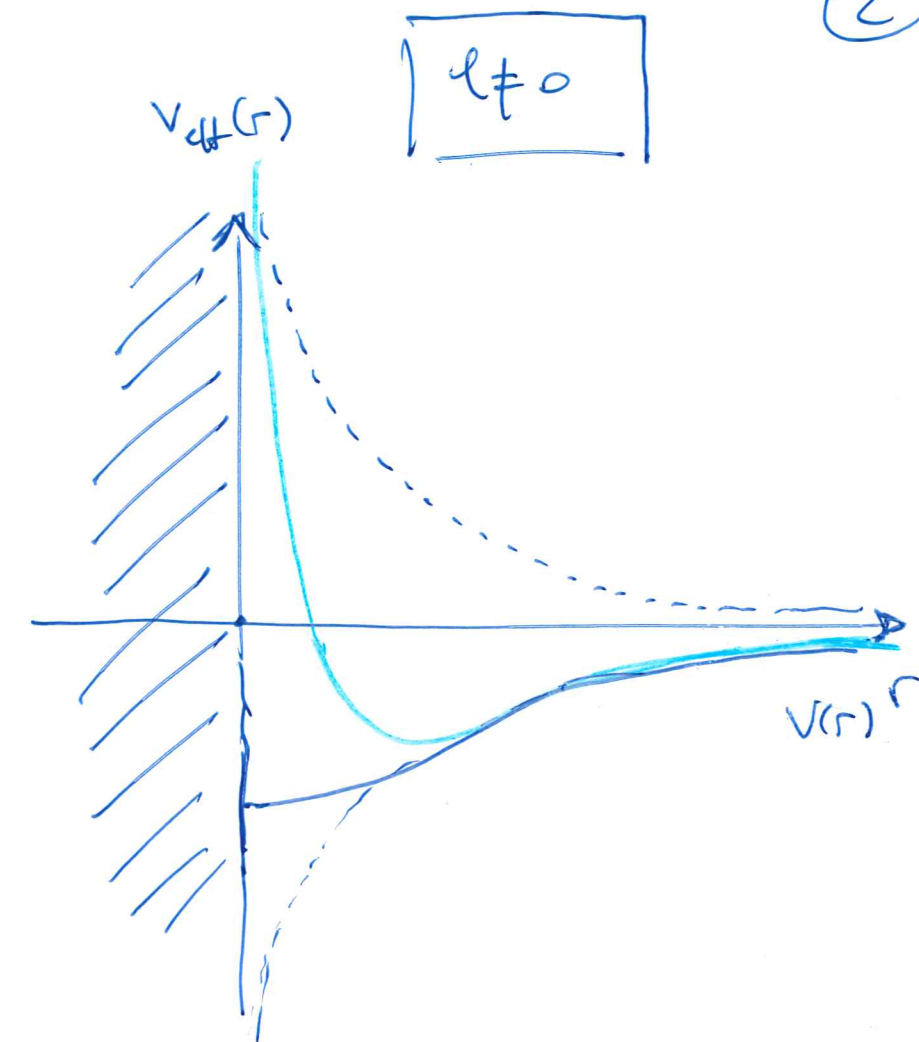
valable aussi pour $\underline{l=0}$



$$u_{nl} \sim r^{-l}$$

$$l > 0$$

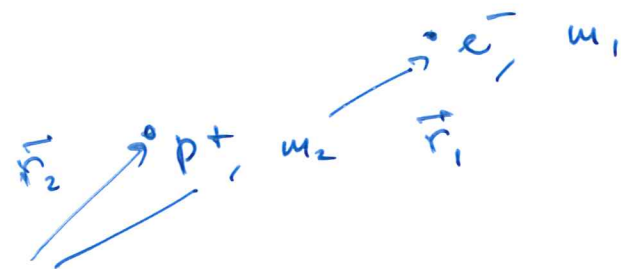
$$R_{nl} \sim r^l$$



(2)

Atome d'hydrogène

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$$E = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 + V(\vec{r}_1 - \vec{r}_2)$$

centre de masse $\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}$

variable relative

$$\vec{r} = \vec{r}_1 - \vec{r}_2$$

$$= \underbrace{\frac{1}{2} \mu \dot{\vec{r}}^2 + \frac{1}{2} M \dot{\vec{R}}^2}_{\text{CM}} + V(\vec{r})$$

relative

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

masse réduite

$$M = m_1 + m_2$$

référentiel de repos du centre de masse $\dot{\vec{R}} = 0$

$$\vec{p} = \mu \dot{\vec{r}}$$

impulsion relative

$$\hat{H} = \frac{\hat{\vec{p}}^2}{2\mu} + V(\vec{r})$$

$$V(\vec{r}) = V(r)$$

$$V(r) = - \frac{e^2}{4\pi\epsilon_0 r}$$

$$\begin{cases} e = 1.6 \times 10^{-19} \text{ C} \\ \epsilon_0 = 8.85 \times 10^{-12} \text{ F/m} \end{cases}$$

$$m_1 = 10^{-30}$$

$$m_2 = 10^{-27}$$

$$\vec{R} = \frac{1}{2} \vec{r}_2$$

$$\vec{r} = \vec{r}_1$$

$$\vec{r}_2 = 0$$

$$\left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + \frac{\hbar^2 l(l+1)}{2\mu r^2} - \frac{e^2}{4\pi\epsilon_0 r} \right] u_{nl} = E_{nl} u_{nl}$$

$\downarrow$ r_{ao}

$$\left[-\frac{d^2}{dr^2} + \frac{l(l+1)}{r^2} - 2 \frac{\mu e^2}{\hbar^2} \frac{1}{r} \right] u_{nl} = \frac{2\mu}{\hbar^2} E_{nl} u_{nl}$$

$$a_0^2 \frac{d^2}{dr^2} = \frac{d^2}{d\rho^2}$$

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{\mu e^2} = \text{rayon de Bohr} \hat{=} 0.5 \text{ \AA} = 5 \times 10^{-11} \text{ m}$$

$$\rho = r/a_0$$

$$v_{nl}(\rho) = u_{nl}(r)$$

$$\left[-\frac{d^2}{d\rho^2} + \frac{l(l+1)}{\rho^2} - \frac{2}{\rho} \right] v_{nl}(\rho) = \underbrace{\left(\frac{2\mu a_0^2}{\hbar^2} E_{nl} \right)}_{E_{nl}} v_{nl}$$

Rydberg

$$Ry = \frac{\hbar^2}{2\mu a_0^2} = 13.6 \text{ eV}$$

cherche solution pour des états liés

$$\underline{E_{nl} < 0}$$

$$\underline{E_{nl} = -|E_{nl}|}$$

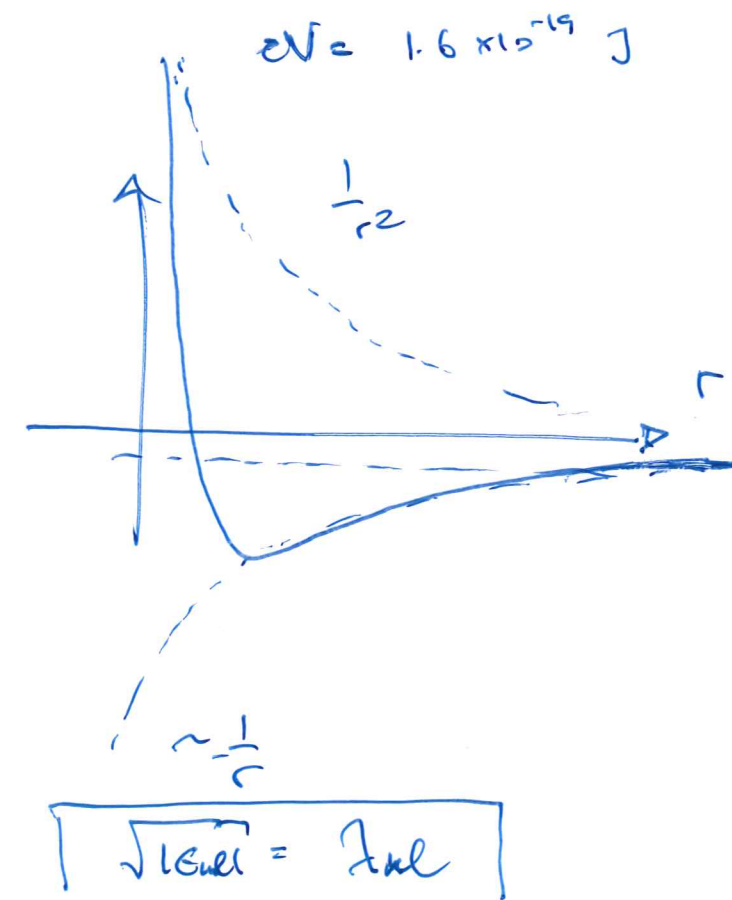
$$\underline{E_{nl} = \frac{E_{nl}}{Ry}}$$

$$v_{nl}(\rho) \underset{\rho \rightarrow 0}{\sim} \rho^{l+1}$$

$\rho \rightarrow \infty$?

$$\frac{d^2}{d\rho^2} v_{nl} = (E_{nl}) v_{nl}$$

$$v_{nl} \sim \exp(\pm \sqrt{|E_{nl}|} \rho) = \exp(-\sqrt{|E_{nl}|} \rho)$$



$$v_{nl}(\rho) = y_{nl}(\rho) e^{-\lambda_{nl}\rho}$$

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$$\left[-\frac{d^2}{d\rho^2} + 2\lambda_{nl} \frac{d}{d\rho} + \frac{l(l+1)}{\rho^2} - \frac{2}{\rho} \right] y_{nl}(\rho) = 0 \quad (**)$$

$$y_{nl}(\rho) \sim \rho^{l+1} \quad \rho \rightarrow 0$$

$$y_{nl}(\rho) = \rho^{l+1} \sum_{q=0}^{\infty} c_q \rho^q \quad (*)$$

$$c_0 \neq 0$$

en injectant (*) en (**) $\rightarrow$

$$\sum_{q=0}^{\infty} A_q(\rho^q) \rho^q = 0$$

$$\underline{A_q = 0} \quad \forall q$$

$$c_q = \frac{2[(q+1)\lambda_{nl} - 1]}{q(q+l+1)} c_{q-1}$$

$$q \rightarrow \infty$$

$$c_q \approx \frac{2\lambda_{nl}}{q} c_{q-1}$$

$$\frac{c_q}{c_{q-1}} \underset{q \rightarrow \infty}{\sim} \frac{2\lambda_{nl}}{q}$$

$$\boxed{q \rightarrow \infty}$$

$$\frac{c_q}{c_{q-1}} = \frac{\frac{(2\lambda_{nl})^q}{q!}}{\frac{(2\lambda_{nl})^{q-1}}{(q-1)!}} = \frac{2\lambda_{nl}}{q}$$

$$\exp(2\lambda_{nl}\rho) = \sum_{q=0}^{\infty} \frac{(2\lambda_{nl})^q}{q!} \rho^q$$

$$y_{nl}(\rho) \xrightarrow{\rho \rightarrow \infty} \exp(2\lambda_{nl}\rho)$$

$$v_{nl}(r) \sim e^{-\lambda_{nl} r} \quad e^{2\lambda_{nl} r} = e^{\lambda_{nl} r}$$

Trouver la série : $\exists k = q_{\max} : c_{k+1} \neq 0$

$$c_k = \frac{2[(k+l)\lambda_{nl} - 1]}{k(k+2l+1)} \quad c_{k+1} = 0$$

$$\lambda_{nl} = \frac{1}{k+l} = \frac{1}{n}$$

$$k \geq 1$$

$$n = l+k \in \mathbb{N}^*$$

$$l = 0, 1, 2, \dots$$

$$n = 1, 2, \dots$$

si vous fixez n : $l = 0, 1, \dots, n-1$

$$E_{nl} = -R_y (E_{nl}) = -R_y \lambda_{nl}^2 = -\frac{R_y}{n^2} = E_n$$

$$n = \text{nombre quantique principale}$$

$$\psi = R_{nl}(r) Y_{lm}(\theta, \phi) = \psi_{nlm}$$

$$l = 0, 1, \dots, n-1$$

nombre quantique de L^2

$$m = -l, \dots, l$$

" " de L^z

lignes de séquence :

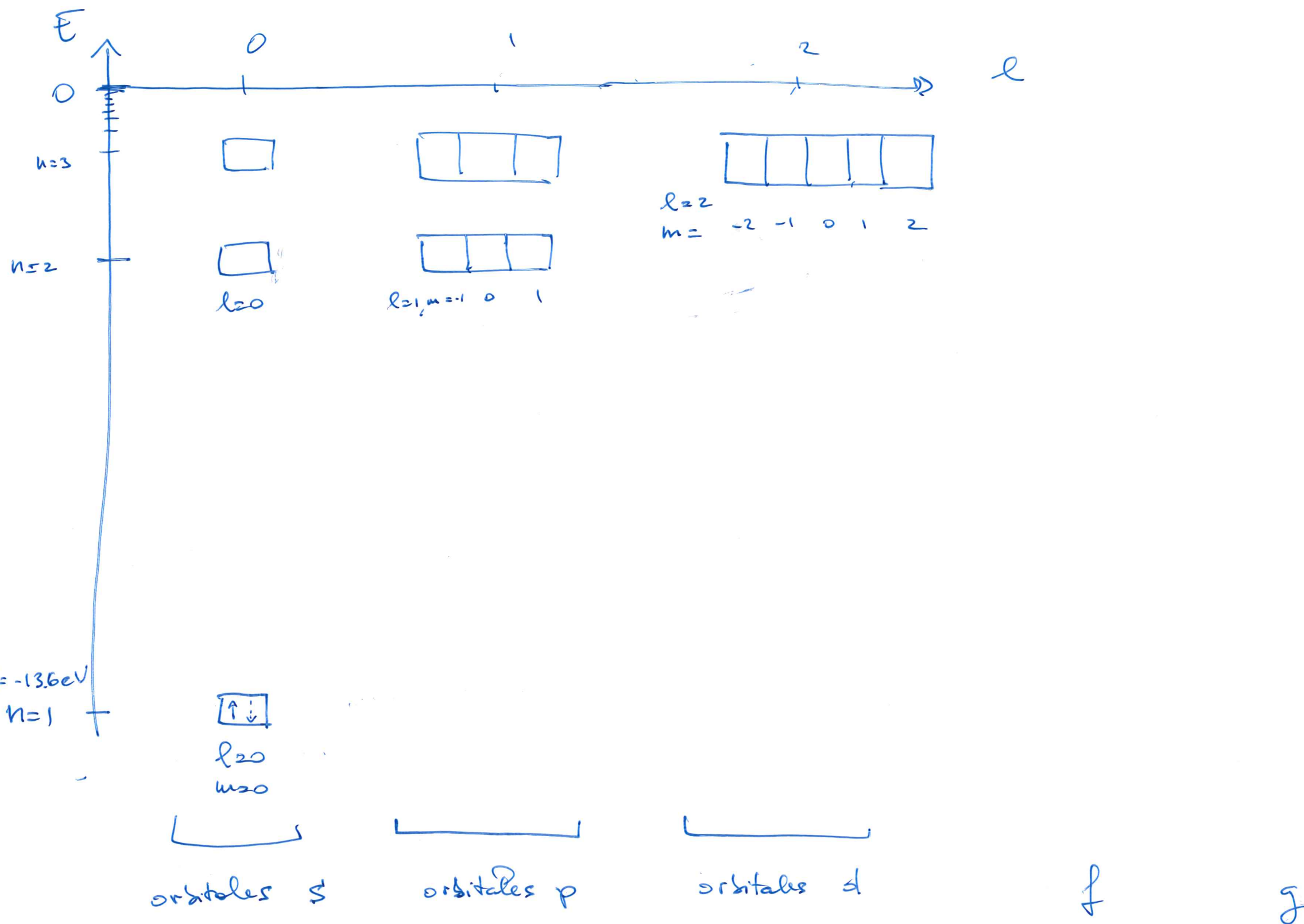
$$J_n = \sum_{l=0}^{n-1} (2l+1) = 2 \frac{n(n-1)}{2} + n = n^2$$

Spectre de l'atome d'hydrogène

$$E_n = - \frac{R_H}{n^2}$$

$$n = 1, 2, \dots$$

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$$|n l m; s m_s\rangle$$

↓
 $\pm \hbar/2$

$$\hat{H}, \hat{L}^2, \hat{L}_z, \hat{S}^2, \hat{S}_z$$

forment un ECOC

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Structure des fonctions propres

$$\psi_{nlm}(r, \vartheta, \phi) = R_{nl}(r) Y_{lm}(\vartheta, \phi)$$

$$R_{nl}(r) = \frac{1}{r} v_{nl}\left(\frac{r}{a_0}\right) = \underbrace{\frac{1}{r} Y_{nl}\left(\frac{r}{a_0}\right)}_{\text{polynomes d'ordre } (n-1)} e^{-\frac{r}{na_0}}$$

$$\rho_{nl}(\rho) = \frac{\rho}{na_0}$$

$$\begin{aligned} \rho_{\max} &= k \\ k+l &= n \\ k &= n-l \end{aligned}$$

$$Y_{nl}(\rho) = \rho^l \underbrace{\sum_{q=0}^{n-l-1} c_q \rho^q}_{\text{polynomes d'ordre } n}$$

$$\left\{ \begin{aligned} R_{10} &= \frac{1}{2a_0^{3/2}} e^{-\frac{r}{a_0}} \\ R_{20} &= \frac{2}{(2a_0)^{3/2}} \left(1 - \frac{r}{2a_0}\right) e^{-\frac{r}{2a_0}} \\ R_{21} &= \frac{1}{(2a_0)^{3/2}} \frac{r}{\sqrt{3} a_0} e^{-\frac{r}{2a_0}} \end{aligned} \right.$$

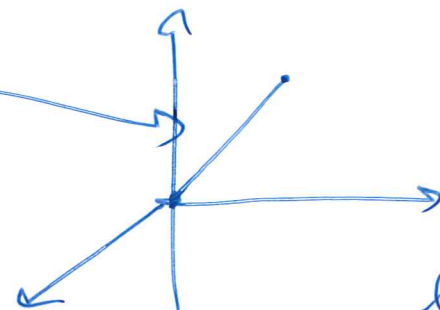
$$\int_0^\infty dr P_{nl}(r) = 1 = \int_0^\infty dr r^2 |R_{nl}|^2$$

$$P_{nl}(r) = r^2 |R_{nl}(r)|^2$$

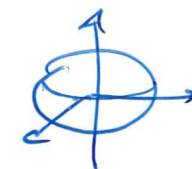
Harmomiques sphériques

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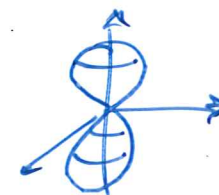
$$|y(\vartheta, \phi)|^2$$



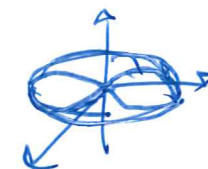
$$l=0 \quad |y_{00}|^2 = \frac{1}{4\pi}$$



$$l=1 \quad |y_{10}|^2 = \frac{3}{4\pi} \cos^2 \vartheta$$

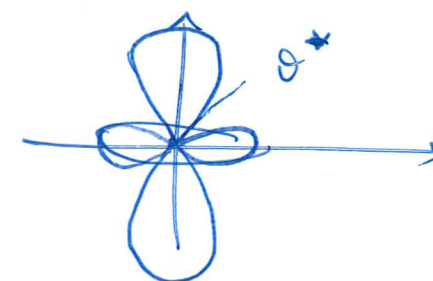


$$|y_{11}|^2 = \frac{3}{8\pi} \sin^2 \vartheta$$



$$l=2 \quad |y_{20}|^2 = \frac{5}{16\pi} (3 \cos^2 \vartheta - 1)^2$$

$$\vartheta^* \approx 54^\circ$$



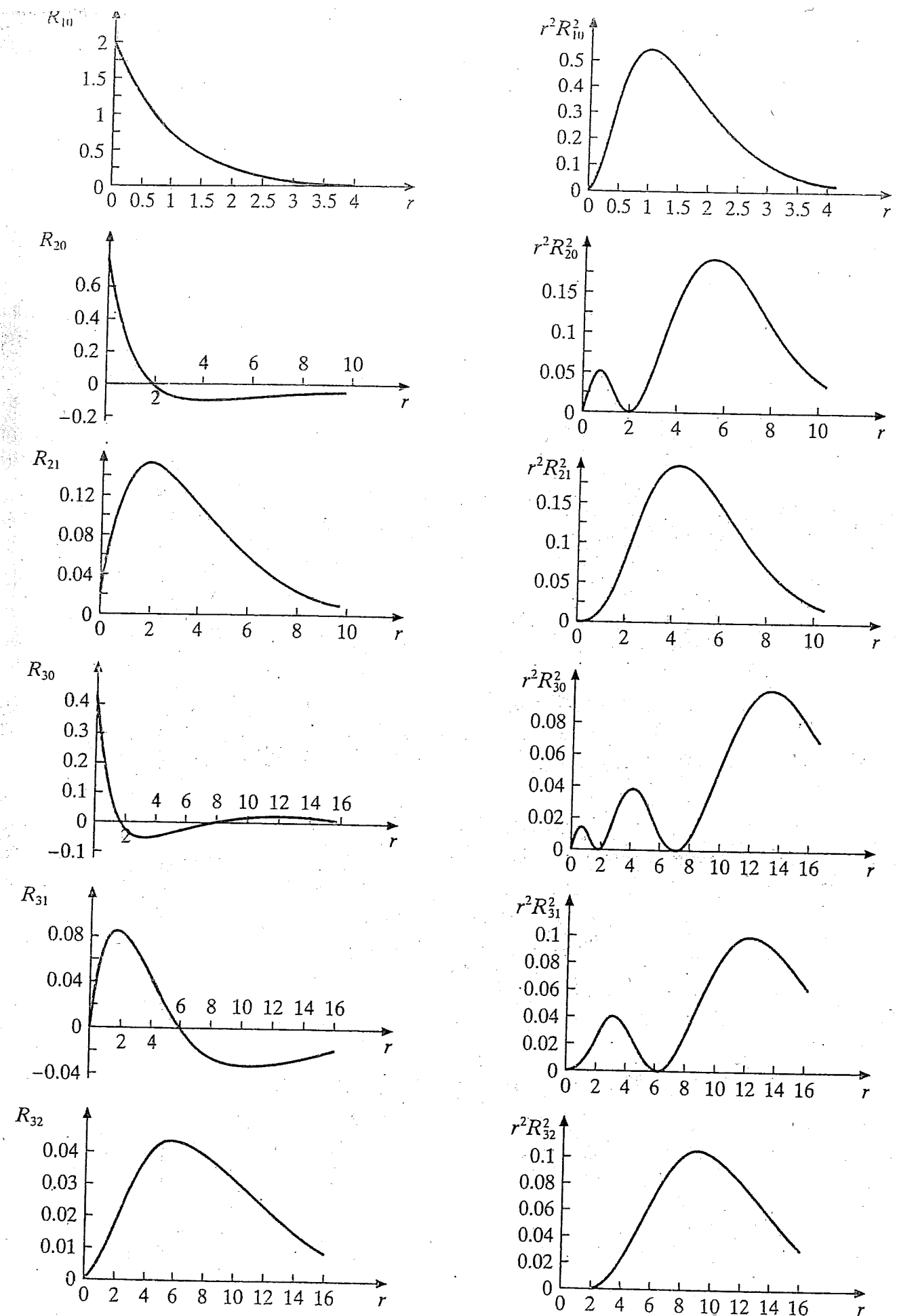
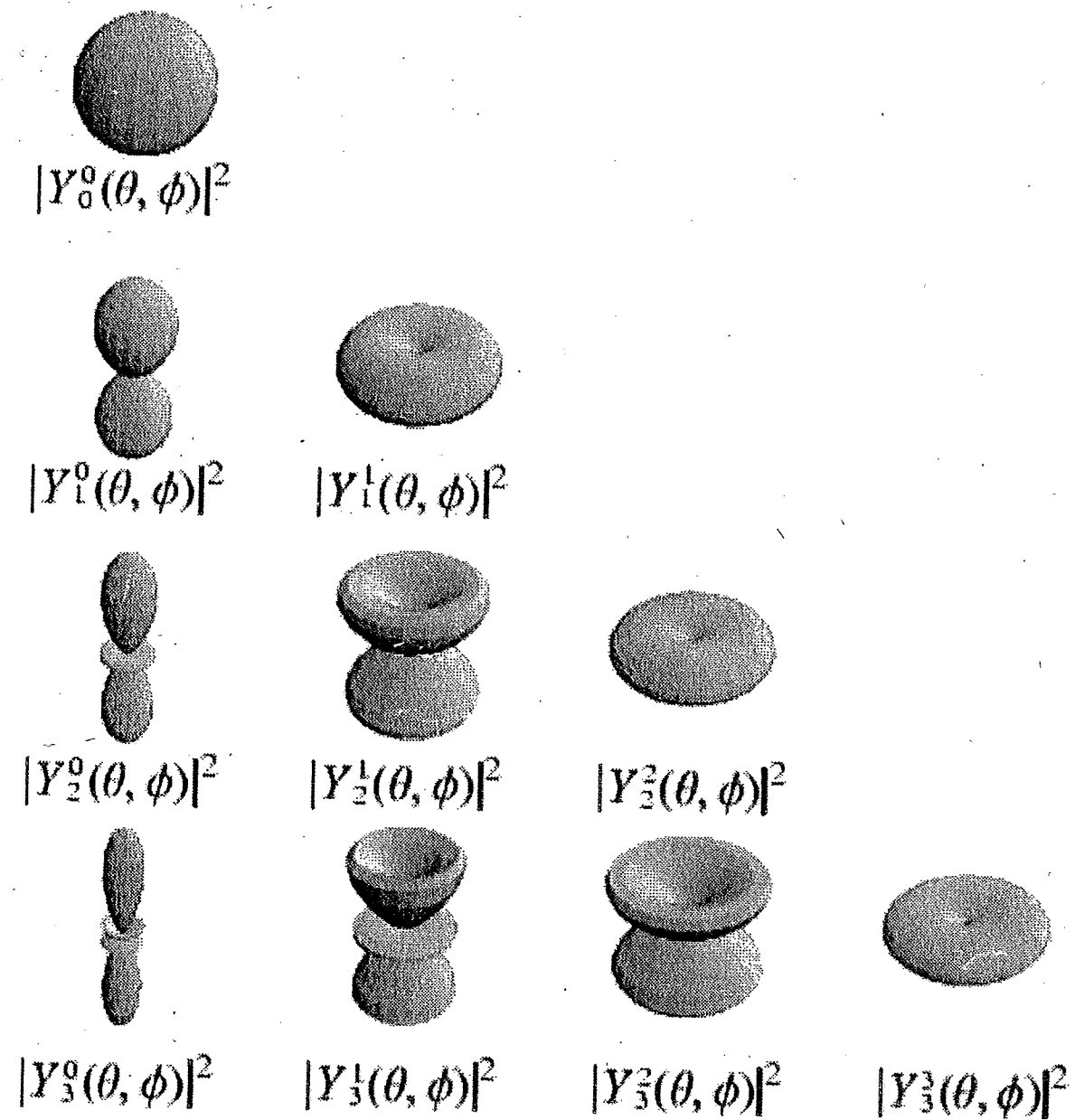
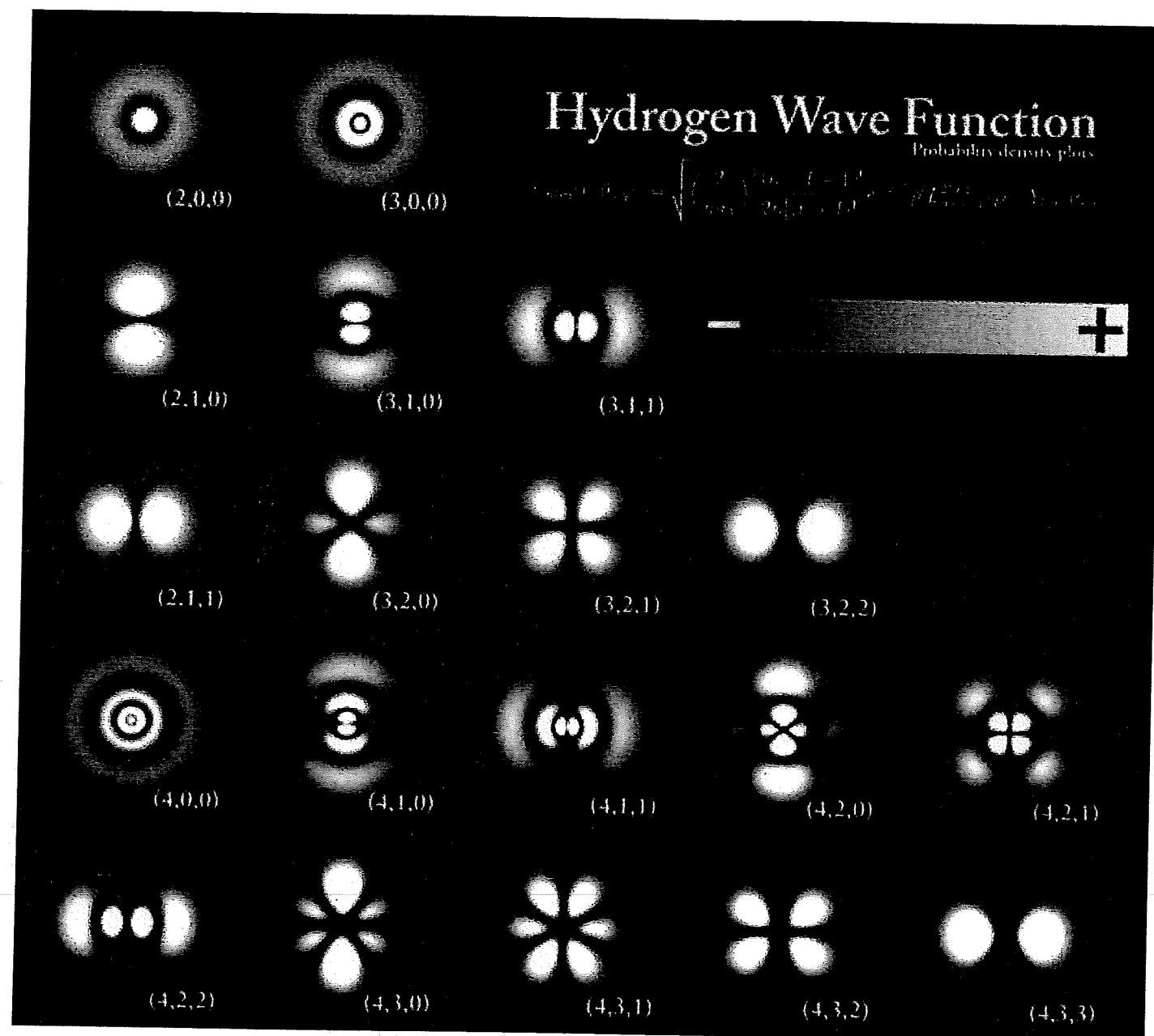
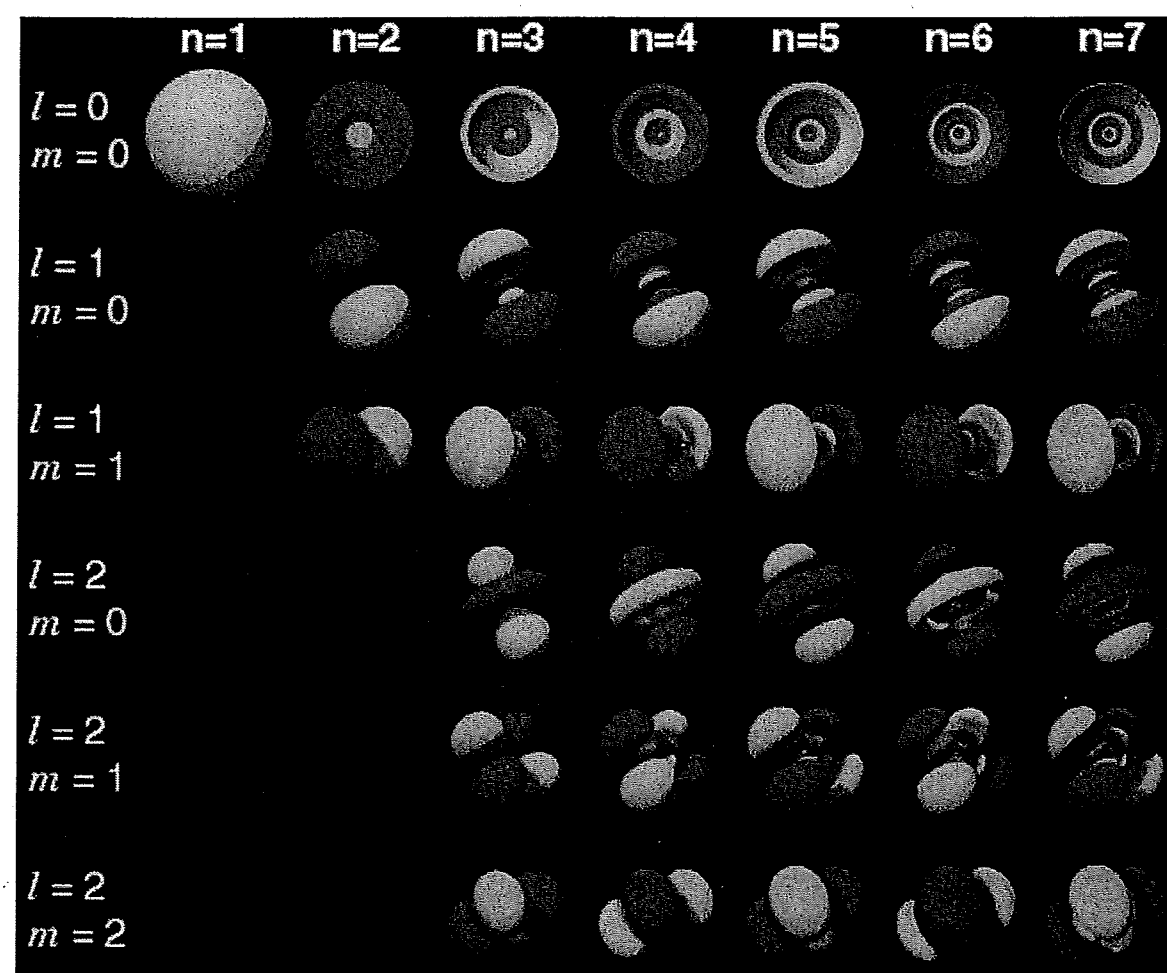


Figure 3.3 Radial functions $R_{nl}(r)$ and radial distribution functions $r^2 R_{nl}^2(r)$ for atomic hydrogen.



http://sevendcolors.org/images/photo/hydrogen_density_plots.jpg



<http://chemistry.beloit.edu/Stars/images/orbitals.jpg>