

Consequences of Bose-Einstein condensation

For weekly extracting base gas

So far : working assumptions that
weak interactions do not destroy BEC.

nearby
100% condensate : spontaneous - symmetry breaking (SSB)

$$\hat{\psi}(\vec{r}) = \hat{\psi}_0(\vec{r}) + \delta\hat{\psi}(\vec{r})$$

$$\text{nearly } 100\% \quad \langle \hat{\Sigma}^+ \hat{\Sigma}^- \rangle \ll \overline{\hat{\Sigma}}_0(\vec{r})$$

↑
Density of
volcanic depletion
↓
cinder cone
density

$$\int d^4r |\tilde{\Sigma}_0|^2 = N_0 \sim (N)$$

$$J_0 = \gamma_0, \quad B \in C \quad \langle \delta \gamma^+, \delta \gamma^- \rangle = 0$$

It has hit all

Gross-Pitaevskii equation (GPE)

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V_{ext}(\vec{r}) + \int d^3r' V(\vec{r}-\vec{r}') |\psi_0(\vec{r}')|^2 \right] \psi_0(\vec{r}) = \mu \psi_0(\vec{r})$$

\Rightarrow solution minimizes the energy in the absence of condensate depletion

time-dependent GPE

$$i\hbar \frac{\partial}{\partial t} \psi_0^{(r,t)} = \left[-\frac{\hbar^2}{2m} \nabla^2 + V_{ext} + \frac{1}{2} \int d^3r' V(\vec{r}-\vec{r}') |\psi_0(\vec{r}',t)|^2 \right] \psi_0(\vec{r},t)$$

how to obtain $\hat{\psi}$

$$\begin{aligned} i\hbar \frac{\partial}{\partial t} \hat{\psi}(\vec{r},t) &= [\hat{\psi}^\dagger, \hat{\psi}] = \hat{\psi} \rightarrow \psi_0 \\ &= \left[-\frac{\hbar^2}{2m} \nabla^2 + V_{ext} + \frac{1}{2} \int d^3r' V(\vec{r}-\vec{r}') \hat{\psi}^\dagger(\vec{r}',t) \hat{\psi}(\vec{r}',t) \right] \hat{\psi}(\vec{r},t) \\ \hat{\psi} &= \hat{H}_{\text{one-body}} + \frac{1}{2} \int d^3r' d^3r'' V(\vec{r}-\vec{r}'') \hat{\psi}^\dagger(\vec{r}') \hat{\psi}^\dagger(\vec{r}'') \\ &\quad \hat{\psi}(\vec{r}') \hat{\psi}(\vec{r}'') \end{aligned}$$

$$i\hbar \frac{\partial \hat{\psi}^\dagger}{\partial t} = [\hat{\psi}^\dagger, \hat{H}] = - \left[\hat{\psi}^\dagger \right]^+$$

$$\left(\frac{\partial \hat{\psi}^\dagger}{\partial t} \right) \hat{\psi} + \hat{\psi}^\dagger \frac{\partial \hat{\psi}}{\partial t} = - \frac{\partial}{\partial t} (\hat{\psi}^\dagger \hat{\psi})$$

$$= \int_{\text{St}} \left[\dots \right]$$

$$= - \vec{\nabla} \cdot \vec{j}$$

$$\frac{\partial}{\partial t} (\hat{\psi}^+ \hat{\psi}) + \vec{\nabla} \cdot \vec{j} = 0$$

$$\langle \vec{j}(\vec{r}, t) \rangle = \frac{i}{2\omega} \left[\hat{\psi}^+ \vec{\sigma} \hat{\psi} - (\vec{\sigma} \hat{\psi}^+) \hat{\psi} \right]$$

$$\langle \psi^+(\vec{r}) \cdot \vec{\sigma} \psi(\vec{r}) \rangle = \lim_{\vec{r}' \rightarrow \vec{r}} \vec{\nabla}_{\vec{r}'} \langle \psi^+(\vec{r}) \psi(\vec{r}') \rangle$$

$$S^{(1)}(\vec{r}, \vec{r}')$$

$$\langle \vec{j}(\vec{r}, t) \rangle = \frac{i}{2\omega} \lim_{\vec{r}' \rightarrow \vec{r}} \vec{\nabla}_{\vec{r}'} \left[S^{(1)}(\vec{r}, \vec{r}'; t) - S^{(1)}(\vec{r}', \vec{r}; t) \right]$$

$$S^{(1)}(\vec{r}, \vec{r}'; t) = \sum_{\alpha} N_{\alpha} \vec{x}_{\alpha}(\vec{r}; t) \vec{x}_{\alpha}(\vec{r}'; t)$$

$$= \sum_{\alpha} \vec{j}_{\alpha}(\vec{r}, t)$$

$$\downarrow \quad \downarrow$$

$$\downarrow \quad \downarrow$$

$$N_{\alpha} \frac{i}{2\omega} \left(\vec{x}_{\alpha}^* \vec{\sigma} \vec{x}_{\alpha} - (\vec{\sigma} \vec{x}_{\alpha}^*) \vec{x}_{\alpha} \right)$$

$$= \boxed{\vec{j}_0(\vec{r}, t)} + \sum_{\alpha \neq 0} \vec{j}_{\alpha}(\vec{r}, t)$$

dominant term

$$\chi_o(\vec{r}, t) = |\chi_o(\vec{r}, t)| e^{i \phi_o(\vec{r}, t)}$$

$$\vec{j}_o(\vec{r}, t) = \frac{1}{m} N_o |\chi_o| (\vec{s} \vec{H}^2) \vec{\nabla} \phi_o(\vec{r}, t)$$

$\rho(\vec{r}, t)$
particle # density

$$= \rho(\vec{r}, t) \vec{v}_s(\vec{r}, t)$$

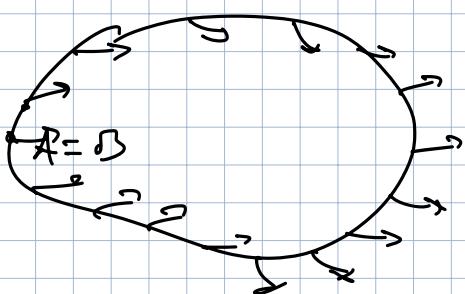
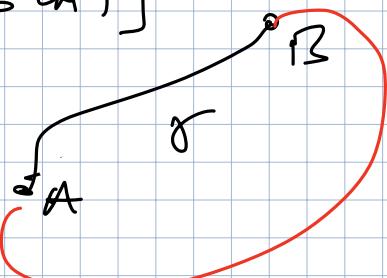
↑
"equilibrium velocity"

$$\vec{v}_s(\vec{r}, t) = \frac{h}{m} \vec{\nabla} \phi_o(\vec{r}, t)$$

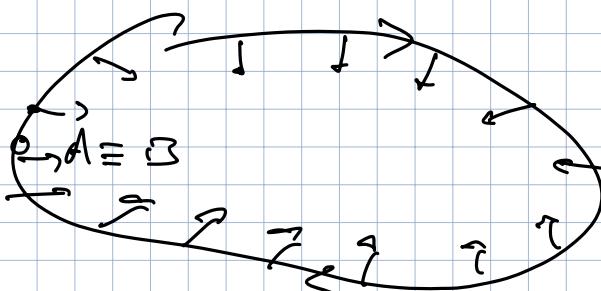
$$K = \int_A^B \vec{dl} \cdot \vec{v}_s(\vec{r}, t) = \frac{h}{m} [\phi_o(B) - \phi_o(A)]$$

$$= \frac{h}{m} 2\pi p$$

Circularity
 $p \in \mathbb{Z}$



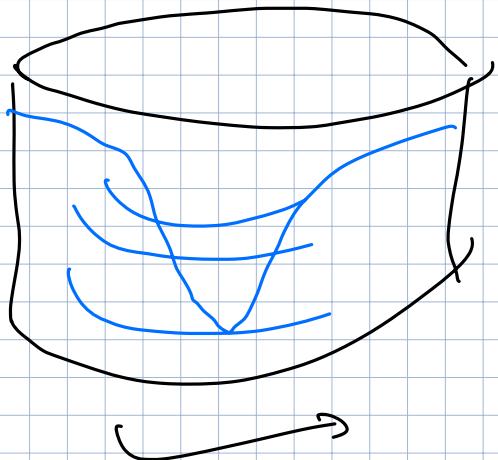
$$p = \infty$$



$$p = -1$$

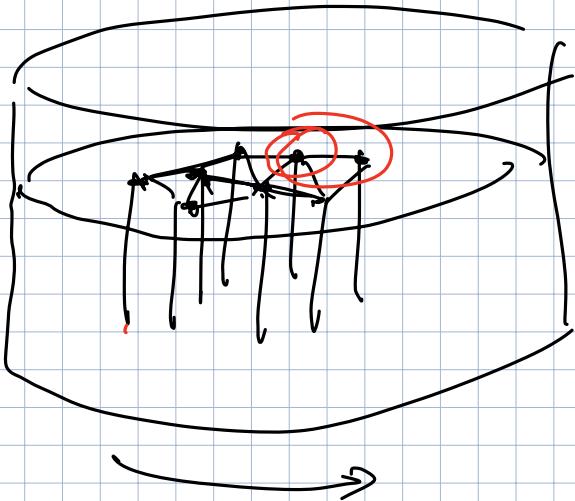
$$V = \frac{h}{m} \cdot P$$

Vertex quantum



Regular fluid

Vortex quantification



BEC

Bogolyubov theory of the condensate depletion

$$\hat{\mathbf{F}}(\vec{r}) = \hat{\mathcal{J}}_0(\vec{r}) + \hat{\mathcal{G}}(\vec{r})$$

$$\hat{H} - \mu \hat{N} = \sum_{\text{G.P.}} [\hat{n}_1, \hat{n}_2] + \hat{H}_1 + \text{quadratic in } \hat{n}_1, \hat{n}_2 + \hat{H}_2 + \hat{H}_3 + \hat{H}_4$$

$$\hat{H}_1 = \int d^d r \sum_{\vec{r}} \psi^+ \left[-\frac{\hbar^2}{2m} \nabla^2 + U_{ext} + \int d^d r' V(\vec{r} - \vec{r}') |\psi_0(\vec{r}')|^2 \right] \psi_0(\vec{r})$$

+ h.c. GPE

= 0

$$\psi^+(\vec{r}) \psi^+(\vec{r}') \psi(\vec{r}') \psi(\vec{r})$$

$$\hat{H}_2 = \int d^d r \sum_{\vec{r}} \psi^+ \left[-\frac{\hbar^2}{2m} \nabla^2 + U_{ext} - \mu \right] \psi$$

+ $\frac{1}{2} \int d^d r \int d^d r' V(\vec{r} - \vec{r}')$

$\left[\underbrace{\delta \psi^+(\vec{r}) \delta \psi^+(\vec{r}')}_{+ 2 \delta \psi^+(\vec{r}) \delta \psi^+(\vec{r}') |h\psi_0(\vec{r}')|^2} \right.$

$\left. + 2 \delta \psi^+(\vec{r}) \delta \psi^+(\vec{r}') \overline{\psi_0(\vec{r}') \psi(\vec{r})} \right]$

+ h.c.

Weakly interacting Bose gas in free space

$$U_{ext} = 0$$

$(\vec{r} - \vec{r}')$

$$V(\vec{r} - \vec{r}') = g \delta(\vec{r} - \vec{r}')$$

$$g = \frac{4\pi \hbar^2 a_s}{m}$$

a_s = s-wave
scattering length

GPE:

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + g |\psi_0|^2 \right] \psi_0 = \mu \psi_0 \quad \leftarrow$$

solution

$$\tilde{H}_0(\vec{r}) = \frac{\sqrt{N_e} e^2}{\sqrt{V}} = \sqrt{n_0} \times \text{const}$$

$$\boxed{\int n_0 = \mu}$$

$$\tilde{H}_0 = \sqrt{n_0}$$

$$\hat{H}_2 = \int d^3r \delta q^\dagger \left(-\frac{\hbar^2}{2m} \nabla^2 + \mu \right) \delta q$$

$$+ \frac{\sqrt{n_0}}{2} \int d^3r \left[\delta q^\dagger \delta q^\dagger + h.c. \right] + 4 \left[\delta q^\dagger \delta q \right]$$

$$\delta q^\dagger = \sum_{\vec{k} \neq 0} \frac{e}{\sqrt{V}} \vec{a}_{\vec{k}}^\dagger$$

$$\hat{H}_2 = \sum_{\vec{k} \neq 0} \left(\frac{\hbar^2 k^2}{2m} - g \mu \right) a_{\vec{k}}^\dagger a_{\vec{k}}$$

$$+ \left(2g \mu \right) \sum_{\vec{k} \neq 0} a_{\vec{k}}^\dagger a_{\vec{k}}$$

$$+ \frac{g \mu}{2} \sum_{\vec{k} \neq 0} \left(a_{\vec{k}}^\dagger a_{-\vec{k}}^\dagger + h.c. \right)$$

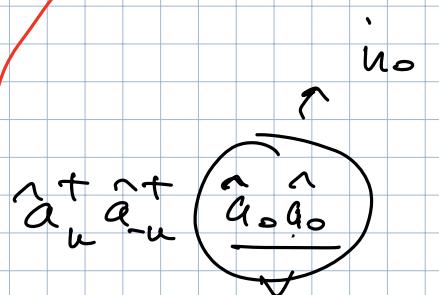
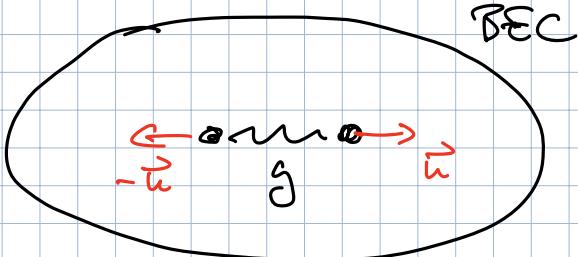
Boyle's
Hamiltonian

$$[\hat{H}_1, \sum_{\vec{k}} \hbar \vec{k} a_{\vec{k}}^\dagger a_{\vec{k}}] = 0$$

$$\hat{H}_2 = \sum_{\vec{u} \neq 0} \left(\frac{\hbar^2 u^2}{2m} + g_{\vec{u}} \right) \hat{a}_{\vec{u}}^\dagger \hat{a}_{\vec{u}}$$

← quadratic in a, a^\dagger

$$+ \frac{g_{\vec{u}}}{2} \sum_{\vec{u} \neq 0} (\hat{a}_{\vec{u}}^\dagger \hat{a}_{\vec{u}} + h.c.)$$



Bogoliubov diagonalization

$$\hat{H}_2 = \sum_{\vec{u} \neq 0} E_{\vec{u}} \hat{b}_{\vec{u}}^\dagger \hat{b}_{\vec{u}} + \text{const.}$$

$$\hat{H}_2 = \frac{1}{2} \sum_{\vec{u} \neq 0} \begin{pmatrix} \hat{a}_{\vec{u}}^\dagger \\ \hat{a}_{\vec{u}} \end{pmatrix}^T \begin{pmatrix} A_{\vec{u}} & B_{\vec{u}} \\ B_{\vec{u}} & A_{\vec{u}} \end{pmatrix} \begin{pmatrix} \hat{a}_{\vec{u}} \\ \hat{a}_{\vec{u}}^\dagger \end{pmatrix} + \text{const.}$$

$$A_{\vec{u}} = \omega_{\vec{u}} + g_{\vec{u}}$$

$$B_{\vec{u}} = g_{\vec{u}}$$

$$\hat{b}_{\vec{u}} = u_{\vec{u}} \hat{a}_{\vec{u}} + v_{\vec{u}} \hat{a}_{-\vec{u}}^\dagger$$

$$\hat{a}_{\vec{u}} = u_{\vec{u}} \hat{b}_{\vec{u}} - v_{\vec{u}} \hat{b}_{-\vec{u}}^\dagger$$

$u_{\vec{u}}, v_{\vec{u}}$ depend only on $|\vec{u}|$

Bogoliubov transformation

$$\begin{pmatrix} \hat{a}_{\vec{u}} \\ \hat{a}_{-\vec{u}}^\dagger \end{pmatrix} = \begin{pmatrix} u & -v \\ -v & u \end{pmatrix} \begin{pmatrix} \hat{b}_{\vec{u}} \\ \hat{b}_{-\vec{u}}^\dagger \end{pmatrix}$$

$$\hat{H}_2 = \frac{1}{2} \sum_{n \neq 0} \left(\begin{pmatrix} \omega_n^+ & \\ & \omega_n^- \end{pmatrix} \begin{pmatrix} u & -v \\ -v & u \end{pmatrix} \begin{pmatrix} A & B \\ B & A \end{pmatrix} \begin{pmatrix} u-v \\ -v-u \end{pmatrix} \right)$$

$$[\omega_n, \omega_n^+] = [1 = \omega_n^2 - \omega_n^{-2}] + \text{const}$$

$$\begin{pmatrix} A(u^2 + v^2) - 2Buv \\ B(u^2 + v^2) - 2Avv \end{pmatrix}$$

$$\begin{pmatrix} A(u^2 + v^2) - 2Buv \\ A(u^2 + v^2) - 2Buv \end{pmatrix}$$

$$u^2 = \frac{1}{2} \left(\frac{A}{\epsilon} + 1 \right)$$

$$v^2 = \frac{1}{2} \left(\frac{A}{\epsilon} - 1 \right)$$

$$uv = \frac{B}{2A} (u^2 + v^2)$$

$$\epsilon = \bar{\epsilon}_n = \sqrt{\frac{A^2}{\epsilon_n} - \frac{B^2}{\epsilon_n}}$$

$$d = d_{(\vec{u})} \geq 0$$

$$\epsilon = \bar{\epsilon}_{(\vec{u})}$$

$$u_{\vec{u}} = \sqrt{\frac{1}{2} \left(\frac{A_{\vec{u}}}{\bar{\epsilon}_{\vec{u}}} + 1 \right)}$$

$$v_{\vec{u}} = \text{sign}(B_{\vec{u}}) \sqrt{\frac{1}{2} \left(\frac{A_{\vec{u}}}{\bar{\epsilon}_{\vec{u}}} - 1 \right)}$$

$$\hat{H}_h = \sum_{\vec{u} \neq 0} (\epsilon_{\vec{u}}) \hat{S}_{\vec{u}}^+ \hat{S}_{\vec{u}}^- + \text{const.}$$

$\hat{S}_{\vec{u}}, \hat{S}_{\vec{u}}^+$: destroy / create Bogoliubov (quasi-) particles

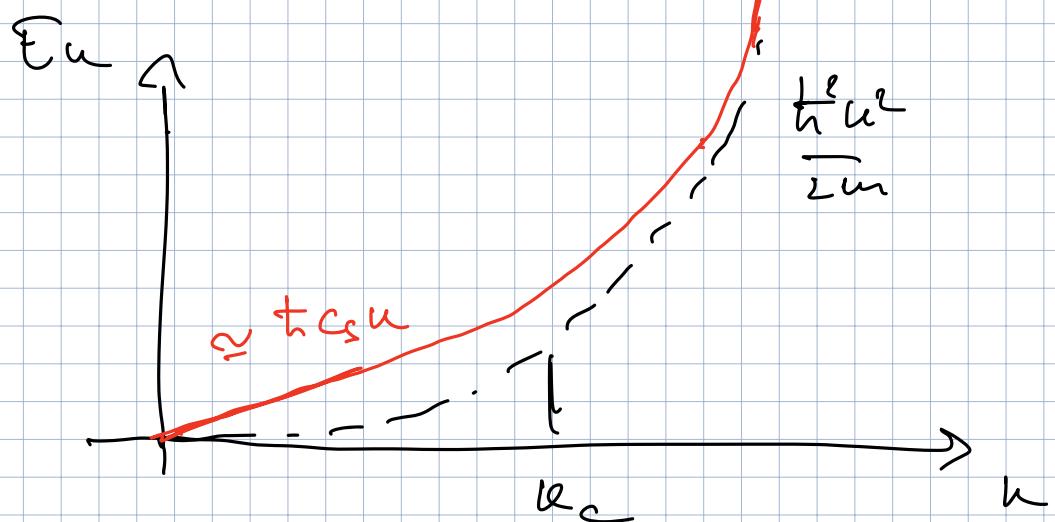
$$\epsilon_{\vec{u}} = \sqrt{\frac{\hbar^2 u^2}{2m} \left(\frac{\hbar^2 u^2}{2m} + g_{us} \right)}$$

$$= \left\{ \frac{\hbar^2 u^2}{2m} \right\}^{1/2} \left(\frac{\hbar^2 u^2}{2m} + g_{us} \right)^{1/2}$$

B. Speed of sound

$$\frac{\hbar^2 u^2}{2m} \ll g_{us}$$

$$\frac{\hbar^2 u^2}{2m} \gg g_{us}$$



$$\frac{\hbar^2 u_c^2}{2m} \approx g_{us}$$

$$\xi = \frac{1}{u_c} = \frac{\hbar}{\sqrt{2m g_{us}}}$$

healing length