

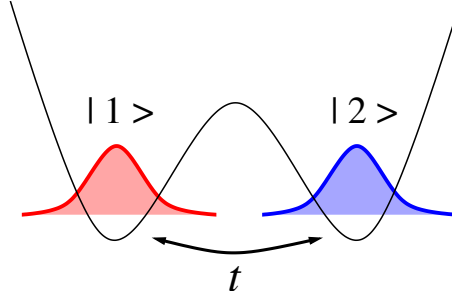
### TD3: Bosons in a double-well potential

In this TD we will investigate the physics of  $N$  identical bosons trapped in a double-well potential (bosonic Josephson junction). In the limit of a very deep potential, we can consider only two orthonormal single-particle states,  $|1\rangle$  and  $|2\rangle$ , localized on the two sides of the double-well potential.

The many-body Hamiltonian on this reduced basis takes the form:

$$\hat{\mathcal{H}} = -t \left( \hat{a}_1^\dagger \hat{a}_2 + \hat{a}_2^\dagger \hat{a}_1 \right) + \frac{U}{2} [\hat{n}_1 (\hat{n}_1 - 1) + \hat{n}_2 (\hat{n}_2 - 1)] \quad (1)$$

where  $\hat{a}_1, \hat{a}_1^\dagger$  and  $\hat{a}_2, \hat{a}_2^\dagger$  are bosonic creation/destruction operators.



## 1 Many-body states: Fock states, Schrödinger's cat states and BEC states.

In this section we consider trial wave-functions to describe the ground state of the model.

### 1.1

Consider the Fock state

$$|N_1, N_2\rangle = \frac{(\hat{a}_1^\dagger)^{N_1}}{\sqrt{N_1!}} \frac{(\hat{a}_2^\dagger)^{N_2}}{\sqrt{N_2!}} |0\rangle \quad (2)$$

Write down its one-body density matrix  $g_{ij}^{(1)} = \langle \hat{a}_i^\dagger \hat{a}_j \rangle$ . For what values of  $N_1$  and  $N_2$  is the system showing condensation/fragmentation?

### 1.2

Consider then the *Schrödinger's cat* state

$$|N_1, N_2; N_2, N_1\rangle = \frac{1}{\sqrt{2}} (|N_1, N_2\rangle + |N_2, N_1\rangle) \quad (3)$$

Calculate its one-body density matrix  $g_{ij}^{(1)}$  for  $|N_1 - N_2| > 1$ ; can this state show condensation?

### 1.3

We then take the BEC state

$$|\alpha, \phi\rangle = \frac{1}{\sqrt{N!}} \left( \frac{\hat{a}_1^\dagger + \alpha e^{i\phi} \hat{a}_2^\dagger}{\sqrt{1 + \alpha^2}} \right)^N |0\rangle \quad (4)$$

Write down its one-body density matrix  $g_{ij}^{(1)}$ . To do so, it can be useful to introduce the orthonormal states

$$|+\rangle = \frac{|1\rangle + \alpha e^{i\phi} |2\rangle}{\sqrt{1 + \alpha^2}} \quad |-\rangle = \frac{\alpha |1\rangle - e^{i\phi} |2\rangle}{\sqrt{1 + \alpha^2}} \quad (5)$$

and the associated creation/destruction operators  $\hat{a}_\pm, \hat{a}_\pm^\dagger$ . Diagonalise  $g_{ij}^{(1)}$  and show that the state corresponds to a perfect condensate (as it was obvious from the first definition....).

### 1.4

Show that the Fock state can be obtained by Fourier transformation of the BEC state with respect to  $\phi$ . Conclude on the existence of a phase-difference/number-difference uncertainty.

## 2 Variational determination of the ground state

### 2.1

Calculate the energy expectation value for the Fock state  $E(N_1, N_2) = \langle N_1, N_2 | \hat{\mathcal{H}} | N_1, N_2 \rangle$ , and for the Schödinger's cat state. Find the combination  $(N_1, N_2)$  minimizing the energy in the case of repulsive ( $U > 0$ ) and attractive ( $U < 0$ ) interactions. Is the Fock state with  $N_1 \neq N_2$  an acceptable equilibrium state of the system?

### 2.2

Calculate the energy expectation value for the BEC state  $E(\alpha, \phi)$ , and show that the energy is an explicit function of the phase  $\phi$ . Assuming  $\alpha = 1$ , minimize with respect to  $\phi$  for both attractive and repulsive interactions.

### 2.3

Comparing the minimum energies  $E_{\min}(N_1, N_2)$  and  $E_{\min}(\alpha = 1, \phi)$ , determine the transition points between fragmented states and condensate states upon varying the ratio  $U/t$ . Draw the corresponding phase diagram on the  $U/t$  axis.

## 3 Schwinger pseudo-spin representation

### 3.1

Introducing the operators

$$\hat{J}_x = \frac{1}{2} (\hat{a}_1^\dagger \hat{a}_2 + \hat{a}_2^\dagger \hat{a}_1) \quad \hat{J}_y = \frac{1}{2i} (\hat{a}_1^\dagger \hat{a}_2 - \hat{a}_2^\dagger \hat{a}_1) \quad \hat{J}_z = \frac{1}{2} (\hat{a}_1^\dagger \hat{a}_1 - \hat{a}_2^\dagger \hat{a}_2) \quad (6)$$

show that they satisfy the commutation relations of angular momentum,  $[\hat{J}_\alpha, \hat{J}_\beta] = i\epsilon^{\alpha\beta\gamma} \hat{J}_\gamma$ .

### 3.2

Calculating  $J^2 = |\hat{\mathbf{J}}|^2$ , determine the effective spin length. Show that the Hamiltonian of the system in the pseudospin variables takes the form

$$\hat{\mathcal{H}} = -2t\hat{J}_x + U \left( \hat{J}_z^2 + J^2 - N \right) . \quad (7)$$

Treating the  $J_\alpha$  operators as classical variables, what are the values of the spin components which minimize the energy in the opposite limits  $|U|/t \ll 1$ ,  $U/t \gg 1$  and  $-U/t \gg 1$ ?

### 3.3

Calculate the vector  $\langle \hat{\mathbf{J}} \rangle = (\langle \hat{J}_x \rangle, \langle \hat{J}_y \rangle, \langle \hat{J}_z \rangle)$  for the Fock state  $|N_1, N_2\rangle$  and the BEC state  $|\alpha, \phi\rangle$ . Introducing the angle variable  $\theta$  such that

$$\frac{1}{\sqrt{1+\alpha^2}} = \cos(\theta/2) \quad \frac{\alpha}{\sqrt{1+\alpha^2}} = \sin(\theta/2) \quad (8)$$

show that the state vector  $\langle \hat{\mathbf{J}} \rangle$  for the BEC state is a vector of length  $N$  with polar/azimuthal angles  $\phi, \theta$ . Represent the ground state of the system in the various phases, as determined in the previous section, via the length and orientation of the state vector.

## 4 Momentum distribution in the BEC state

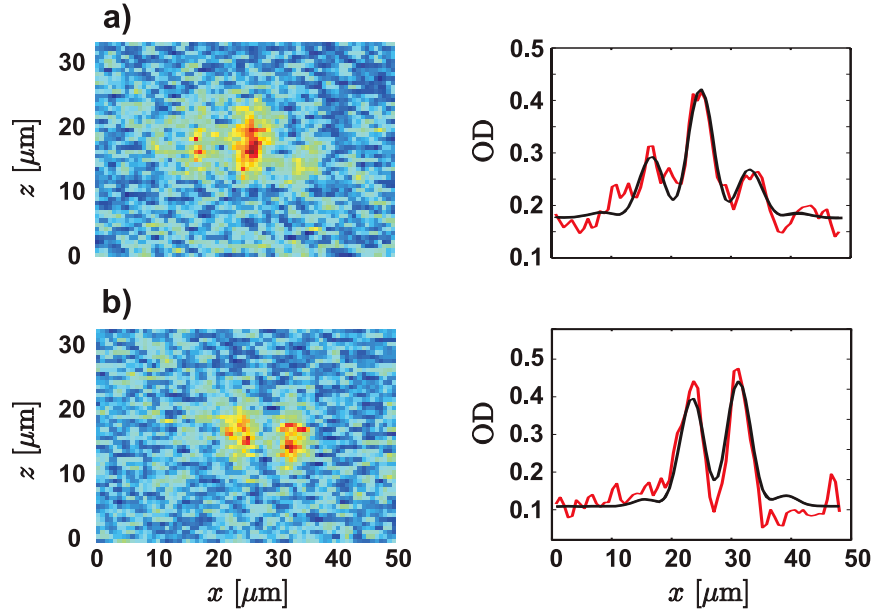


Figure 1: Absorption images of bosons released from a double-well trap. From M. Albiez's PhD thesis, University of Heidelberg (2005).

Be  $\Phi_1(x)$  and  $\Phi_2(x)$  the spatial wavefunction associated with the two states  $|1\rangle$  and  $|2\rangle$ , with the property  $\Phi_2(x) = \Phi_1(x - d)$  ( $d$  = separation between the two wells). The states  $\Phi_1$  and  $\Phi_2$  are part of a basis  $\Phi_i$  of orthonormal wavefunctions.

### 4.1

Calculate  $\langle \hat{\psi}^\dagger(x) \hat{\psi}(x') \rangle$  for the generic BEC state  $|\alpha, \phi\rangle$ .

(Suggestion: expand the field operators on the  $\Phi_i$  basis).

## 4.2

Calculate the momentum distribution  $n(k) = \langle \hat{a}_k^\dagger \hat{a}_k \rangle$ , where

$$a_k = \int dx \frac{e^{-ikx}}{\sqrt{V}} \psi(x) . \quad (9)$$

You should find an expression containing  $N$ ,  $\alpha$ ,  $\phi$ ,  $k$  and  $\tilde{\Phi}_1(k)$  (Fourier transform of the  $\Phi_1$  wavefunction). In the case of Gaussian wavefunctions,  $\Phi_1(x) \sim \exp[-x^2/(2\sigma^2)]$ , sketch the form of  $n(k)$ .

## 4.3

Fig. 1 shows the measurement of the momentum distribution of bosons trapped in double well potential for two different measurement shots. Can you tell the phase difference  $\phi$  between the two wells in the two cases?