

TD8: Anderson's pseudo-spin model and BCS variational wavefunction

In this TD we will explore a very insightful approach to the BCS Hamiltonian provided by P. W. Anderson (1958). Recasting the BCS Hamiltonian in terms of pseudo-spin variables, we will be able to write down the ground state of BCS theory in a very transparent and suggestive way.

NOTE: the angle $\vartheta_{\mathbf{k}}$ that appears in the following is not the same as the angle $\theta_{\mathbf{k}}$ discussed in the lectures on BCS theory. We shall establish the relationship between the two later.

1 Pair operators and spin operators

Consider the pair operators

$$\hat{b}_{\mathbf{k}} = \hat{c}_{-\mathbf{k}\downarrow} \hat{c}_{\mathbf{k}\uparrow} \quad \hat{b}_{\mathbf{k}}^{\dagger} = \hat{c}_{\mathbf{k}\uparrow}^{\dagger} \hat{c}_{-\mathbf{k}\downarrow}^{\dagger} \quad (1)$$

1.1

Show that

$$[\hat{b}_{\mathbf{k}}, \hat{b}_{\mathbf{q}}] = [\hat{b}_{\mathbf{k}}^{\dagger}, \hat{b}_{\mathbf{q}}^{\dagger}] = 0 \quad (2)$$

$$[\hat{b}_{\mathbf{k}}, \hat{b}_{\mathbf{q}}^{\dagger}] = (1 - n_{\mathbf{k},\uparrow} - n_{-\mathbf{k},\downarrow}) \delta_{\mathbf{q},\mathbf{k}} \quad (3)$$

Moreover, justify that $(\hat{b}_{\mathbf{k}}^{\dagger})^2 = (\hat{b}_{\mathbf{k}})^2 = 0$.

1.2

Introducing the operators (Anderson's pseudospins)

$$\begin{aligned} \hat{S}_{\mathbf{k}}^z &= \frac{1}{2} (\hat{n}_{\mathbf{k},\uparrow} + \hat{n}_{-\mathbf{k},\downarrow} - 1) \\ \hat{S}_{\mathbf{k}}^+ &= \hat{b}_{\mathbf{k}}^{\dagger} \\ \hat{S}_{\mathbf{k}}^- &= \hat{b}_{\mathbf{k}} \end{aligned} \quad (4)$$

show that they satisfy the commutation relations of angular momentum

$$[\hat{S}_{\mathbf{k}}^+, \hat{S}_{\mathbf{k}}^-] = 2\hat{S}_{\mathbf{k}}^z \quad [\hat{S}_{\mathbf{k}}^+, \hat{S}_{\mathbf{k}}^z] = -\hat{S}_{\mathbf{k}}^+ \quad (5)$$

Given that $(\hat{S}_{\mathbf{k}}^+)^2 = (\hat{S}_{\mathbf{k}}^-)^2 = 0$, what is the spin length S ?

2 BCS Hamiltonian as a spin Hamiltonian

The BCS Hamiltonian for (quasi-)electrons interacting via an effective phonon-mediated interaction reads

$$\hat{\mathcal{H}} - \mu\hat{N} = \sum_{\mathbf{k}} (\epsilon_{\mathbf{k}} - \mu) (\hat{n}_{\mathbf{k},\uparrow} + \hat{n}_{\mathbf{k},\downarrow}) - \frac{V_0}{\mathcal{V}} \sum'_{\mathbf{k},\mathbf{q}} \hat{b}_{\mathbf{k}}^\dagger \hat{b}_{\mathbf{q}} \quad . \quad (6)$$

Here \mathcal{V} is the volume, V_0 is the strength of the interaction, and the sum $\sum'_{\mathbf{k},\mathbf{q}}$ runs over momenta \mathbf{k} and \mathbf{q} such that $|\epsilon_{\mathbf{k}(\mathbf{q})} - \mu| \leq \epsilon_c$, where $\epsilon_c \approx \hbar\omega_D$ is the characteristic energy cutoff of the interaction.

2.1

Rewrite the above Hamiltonian in terms of the pseudo-spin operators. You should find

$$\hat{\mathcal{H}} - \mu\hat{N} = \sum_{\mathbf{k}} \left[2(\epsilon_{\mathbf{k}} - \mu) - \frac{V_0}{\mathcal{V}} \vartheta(\epsilon_c - |\epsilon_{\mathbf{k}} - \mu|) \right] \hat{S}_{\mathbf{k}}^z - \frac{V_0}{\mathcal{V}} \sum'_{\mathbf{k} \neq \mathbf{q}} \left(\hat{S}_{\mathbf{k}}^x \hat{S}_{\mathbf{q}}^x + \hat{S}_{\mathbf{k}}^y \hat{S}_{\mathbf{q}}^y \right) + \text{const.} \quad (7)$$

It might be useful to remember the following relationships

$$\hat{S}_{\mathbf{k}}^+ \hat{S}_{\mathbf{k}}^- = \frac{1}{2} + \hat{S}_{\mathbf{k}}^z \quad \frac{1}{2} \left(\hat{S}_{\mathbf{k}}^+ \hat{S}_{\mathbf{q}}^- + \hat{S}_{\mathbf{q}}^+ \hat{S}_{\mathbf{k}}^- \right) = \hat{S}_{\mathbf{k}}^x \hat{S}_{\mathbf{q}}^x + \hat{S}_{\mathbf{k}}^y \hat{S}_{\mathbf{q}}^y \quad . \quad (8)$$

Now, think of \mathbf{k} -space as a lattice (it is discretized after all, due to the boundary conditions), and put a $S = 1/2$ (pseudo-)spin on each lattice site. The above Hamiltonian effectively describes a system of interacting spins on a lattice. Taking the spin at lattice site \mathbf{k} , which are the sites that this spin is interacting with? And what is the value of the local magnetic field?

2.2

We take for the moment $V_0 = 0$. Write down the ground state for each pseudo-spin $\hat{S}_{\mathbf{k}}$. If you report the \mathbf{k} points on the energy axis $\epsilon_{\mathbf{k}}$ (a simply sketch is sufficient) together with the associated pseudo-spin, can you find a “domain wall” at a given energy in the spin configuration? And can you anticipate qualitatively what happens when the interaction is turned on?

3 BCS variational wavefunction

We now look for the ground state of the interacting system ($V_0 \neq 0$) in a *factorized* form, namely in the form

$$|\Psi_0\rangle = \prod_{\mathbf{k}} |\vartheta_{\mathbf{k}}, \phi_{\mathbf{k}}\rangle \quad (9)$$

where

$$|\vartheta, \phi\rangle = \cos(\vartheta/2) |\uparrow\rangle + \sin(\vartheta/2) e^{i\phi} |\downarrow\rangle \quad (10)$$

3.1

Show that the above wavefunction is equivalent to the so-called BCS wavefunction

$$|\Psi_0\rangle = \prod_{\mathbf{k}} \left(u_{\mathbf{k}} + v_{\mathbf{k}} \hat{b}_{\mathbf{k}}^\dagger \right) |0\rangle \quad (11)$$

where the coefficients $u_{\mathbf{k}}$ and $v_{\mathbf{k}}$ are to be identified. (Suggestion: write the vacuum $|0\rangle$ in terms of the pseudo-spin states).

3.2

Show that the expectation value of the spin operator $\hat{\mathbf{S}}$ on the $|\vartheta, \phi\rangle$ state behaves like a classical spin of length $S = 1/2$:

$$\langle \vartheta, \phi | \hat{\mathbf{S}} | \vartheta, \phi \rangle = \frac{1}{2} (\cos \phi \sin \vartheta, \sin \phi \sin \vartheta, \cos \vartheta) \quad (12)$$

and, moreover

$$\langle \vartheta_{\mathbf{k}}, \phi_{\mathbf{k}} | \langle \vartheta_{\mathbf{q}}, \phi_{\mathbf{q}} | \left(\hat{S}_{\mathbf{k}}^x \hat{S}_{\mathbf{q}}^x + \hat{S}_{\mathbf{k}}^y \hat{S}_{\mathbf{q}}^y \right) | \vartheta_{\mathbf{k}}, \phi_{\mathbf{k}} \rangle | \vartheta_{\mathbf{q}}, \phi_{\mathbf{q}} \rangle = \frac{1}{4} \cos(\phi_{\mathbf{k}} - \phi_{\mathbf{q}}) \sin \vartheta_{\mathbf{k}} \sin \vartheta_{\mathbf{q}} \quad (13)$$

Show that the variational energy takes the form

$$\langle \Psi_0 | \hat{\mathcal{H}} - \mu \hat{N} | \Psi_0 \rangle = \sum_{\mathbf{k}} \left[\epsilon_{\mathbf{k}} - \mu - \frac{V_0}{2\mathcal{V}} \vartheta (\epsilon_c - |\epsilon_{\mathbf{k}} - \mu|) \right] \cos \vartheta_{\mathbf{k}} - \frac{V_0}{4\mathcal{V}} \sum'_{\mathbf{k} \neq \mathbf{q}} \cos(\phi_{\mathbf{k}} - \phi_{\mathbf{q}}) \sin \vartheta_{\mathbf{k}} \sin \vartheta_{\mathbf{q}} \quad (14)$$

Given that $V_0 > 0$ and $\vartheta_{\mathbf{k}} \in [0, \pi]$, what value of the $\phi_{\mathbf{k}}$ angles minimizes the energy?

3.3

Minimize the variational energy with respect to $\vartheta_{\mathbf{k}}$ for $|\epsilon_{\mathbf{k}} - \mu| < \epsilon_c$, to find the condition

$$\left(\epsilon_{\mathbf{k}} - \mu - \frac{V_0}{2\mathcal{V}} \right) \sin \vartheta_{\mathbf{k}} = -\frac{V_0}{2\mathcal{V}} \sum'_{\mathbf{q}} \sin \vartheta_{\mathbf{q}} \cos \vartheta_{\mathbf{k}} \quad (15)$$

This condition defines a set of coupled equations.

Making an error of order $\mathcal{O}(1/\mathcal{V})$, we can neglect the term $V_0/(2\mathcal{V})$ on the left-hand side and add the term with $\mathbf{q} = \mathbf{k}$ in the sum on the right-hand side.

Introducing then the symbol

$$\Delta = \frac{V_0}{2\mathcal{V}} \sum'_{\mathbf{q}} \sin \vartheta_{\mathbf{q}} \quad (16)$$

rewrite Eq. (15) in terms of Δ , $\epsilon_{\mathbf{k}}$, μ and $\vartheta_{\mathbf{k}}$.

3.4

Solve for $\vartheta_{\mathbf{k}}$, to find

$$\sin \vartheta_{\mathbf{k}} = \frac{\Delta}{E_{\mathbf{k}}} \quad \cos \vartheta_{\mathbf{k}} = \frac{\mu - \epsilon_{\mathbf{k}}}{E_{\mathbf{k}}} \quad E_{\mathbf{k}} = \sqrt{\Delta^2 + (\epsilon_{\mathbf{k}} - \mu)^2} \quad (17)$$

What is the pseudo-spin orientation at the Fermi energy $\epsilon_F = \mu$? Draw schematically how the pseudo-spin orientation evolve when the energy goes from $\mu - \epsilon_c$ to $\mu + \epsilon_c$.

3.5

From Eq. (16) recover the (implicit) gap equation for $|\Delta|$ as seen in the lecture.

4 Average particle number

The BCS wavefunction, Eq. (11), does not contain a well defined particle number. In particular, $|\Psi_0\rangle$ contains all possible *even* particle numbers from 0 to ∞ . But let us have a look at how well defined the *average* particle number is

4.1

Express the average particle number

$$\langle \hat{N} \rangle = \sum_{\mathbf{k}} \langle \hat{n}_{\mathbf{k},\uparrow} + \hat{n}_{\mathbf{k},\downarrow} \rangle \quad (18)$$

and the average square number

$$\langle \hat{N}^2 \rangle = \sum_{\mathbf{k},\mathbf{q}} \langle (\hat{n}_{\mathbf{k},\uparrow} + \hat{n}_{\mathbf{k},\downarrow}) (\hat{n}_{\mathbf{q},\uparrow} + \hat{n}_{\mathbf{q},\downarrow}) \rangle \quad (19)$$

in terms of the $\vartheta_{\mathbf{k}}$ angles.

4.2

Show that

$$\langle \delta^2 \hat{N} \rangle = \sum_{\mathbf{k}} (1 - \langle \cos \vartheta_{\mathbf{k}} \rangle^2) \sim \mathcal{O}(N) \quad (20)$$

How can one conclude that the sum is $\mathcal{O}(N)$? Which \mathbf{k} states are contributing to it?

4.3

Conclude on the importance of the relative particle-number fluctuations.

5 Angle $\vartheta_{\mathbf{k}}$ vs. angle $\theta_{\mathbf{k}}$

Consider the angle $\theta_{\mathbf{k}}$ given in the lectures on BCS theory (such that $u_{\mathbf{k}} = \cos(\theta_{\mathbf{k}}/2)$, $v_{\mathbf{k}} = \sin(\theta_{\mathbf{k}}/2)$): what is the relationship with the angle $\vartheta_{\mathbf{k}}$ introduced in this exercise?