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The Potts model on random graphs, loop models and random matrices

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St Petersbourg, june 2012

Based on joint work with V.Jones, D.Shlyakhtenko, P. Zinn Justin

Loop models, random matrices and planar algebras, to appear in CMP

and E. Maurel Segala

Combinatorial aspects of matrix models, Alea 2006

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2

Solving the matrix model

Outline

The Potts model on random maps

Enumerating planar maps

Random matrices

Solving the matrix model

Solving the matrix model

Outline

The Potts model on random maps

Enumerating planar maps

Random matrices

Solving the matrix model

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The Potts model

The partition function of the Potts model on a graph G = (V, E) is given by

$$Z_{G} = \sum_{\sigma: V \to \{1, ..., Q\}} \exp(K \sum_{\{i,j\} \in E} \delta_{\sigma_{i}, \sigma_{j}})$$
$$= \sum_{\sigma: V \to \{1, ..., Q\}} \prod_{\{i,j\} \in E} (1 + v \delta_{\sigma_{i}, \sigma_{j}})$$
$$= \sum_{E' \subset E} v^{\sharp \text{ bonds }} Q^{\sharp \text{ clusters}}$$

 $v = e^{K} - 1$, bonds=edges in E', clusters=connected components of the subgraph (V, E')



Definition of maps

We shall consider the Potts model on (random planar maps). A map is a connected graph which is embedded into a surface in such a way that edges do not cross and faces (obtained by cutting the surface along the edges) are homeomorphic to a disk. The genus of the map is the minimal genus of a surface in which it can be properly embedded.





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Potts model on planar maps

We shall consider the Potts model on random planar maps. We assume these graphs are rooted, that is are given a distinguished oriented edge. It is given by the partition function

$$Z = \sum_{G=(V,E)} x^{\#E} y^{\#V} Z_G$$

=
$$\sum_{G=(V,E)} x^{\#E} y^{\#V} \sum_{\sigma: V \to \{1,...,Q\}} \exp(K \sum_{\{i,j\} \in E} \delta_{\sigma_i,\sigma_j})$$

=
$$\sum_{G=(V,E)} x^{\#E} y^{\#V} \sum_{E' \subset E} v^{\sharp \text{ bonds }} Q^{\sharp \text{ clusters}}$$

 $v = e^{K} - 1$, bonds=edges in E', clusters=connected components of the subgraph (V, E')

The Potts model on random planar maps as a loop model If G is a planar map, there is dual (green) and a medial (blue) planar graph G_m



The Potts model on random planar maps as a loop model If G is a planar map, there is dual (green) and a medial (blue) planar graph G_m



Splitting a vertex

The Potts model on random planar maps as a loop model



If G is a planar map, there is a bijection

 $(G, E') \Leftrightarrow$ (loops, shaded vertices)

Moreover, writing Euler formula in each cluster gives the relation

$$\sharp$$
loops = 2 \sharp clusters + \sharp bonds - $\sharp V$

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The Potts model on random graphs as a loop model

The equivalence to the loop model allows to state that

$$Z = \sum_{G = (V,E)} x^{\#E} Q^{-\frac{1}{2} \#V} \sum_{E' \subset E} v^{\#} \text{ bonds } Q^{\#} \text{ clusters}$$
$$= \sum_{\Gamma} \delta^{\#} \log \alpha^{\#} \beta^{\#} \beta^{\#}$$

where the summation is restricted to 4-valent rooted planar maps, and

$$\delta = \sqrt{Q} \qquad \frac{\alpha}{\beta} = \frac{V}{\sqrt{Q}} \qquad \beta = x$$

 δ is called the fugacity.

The Potts model on planar maps and loop models Hence, the partition function Z of the Potts model on planar maps is a generating function for the number of possible matchings of the end points of n copies of the vertex

so that the resulting graph

- is planar, connected,
- has p loops,
- is checkerboard shaded.





The Potts model on random maps

Enumerating planar maps

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Solving the matrix model

Outline

The Potts model on random maps

Enumerating planar maps

Random matrices

Solving the matrix model

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Definition of maps

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it can be properly embedded.

- $2 2g = \ddagger$ vertices
 - + faces
 - # edges.



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"Baby combinatorial problem"

Be given p vertices of valence d drawn on a surface (that is given an orientation).

Problem : Count the number of maps with genus g that can be obtained by matching the end points of the half-edges of the vertices.



3

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Example : g = 0, p = 4, d = 3.



"Baby combinatorial problem"

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Example : g = 0, p = 4, d = 3

The counting is done up to homeomorphisms. How do you count symmetries? (between vertices, edges of the vertices etc)

Solving the matrix model

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Combinatorial approach

Tutte (60's) : count rooted planar maps. A root= A distinguished oriented edge. Prescribing a root reduces the number of symmetries;



A map M with n edges has $2n/\sharp$ Automorphism(M) possible roots. Later we shall count maps with labelled rooted vertices; a rooted map with n + 1 vertices has $n!3^n$ possible labellings and rooting of its vertices.

Random matrices

Solving the matrix model

Tutte's Theorem

The number A((p, 3)) of rooted planar maps with p vertices of valence 3 is equal to $2^{p+1} \frac{(3p)!}{p!(2p+2)!}$

Proof: Surgery on maps =Induction relations on number of maps. Let A((p,3), (1, n)) be the number of planar maps with p vertices of degree 3 and one of degree n.

$$A((p,3),(1,n)) = \# \{ Y \not\approx Y \}$$
$$= \# \{ Y \not\ll Y \} + \# \{ Y \not\approx Y \}$$

 $= A((p-1,3),(1,n+1)) + \sum_{k=0}^{n-2} \sum_{\ell=0}^{p} A((\ell,3),(1,k)) A((p-\ell,3),(1,n-k-2))$

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More general map enumeration

More general maps could be enumerated by using combinatorial techniques, including surgery but also bijections etc For instance

- The Ising model on random planar maps [Bousquet-Mélou, Schaeffer](2002)
- The Potts model on random planar maps [Bernardi, Bousquet-Mélou](2009)

However each case is treated separately and is not easy. The physicists have used since the seventies a rather indirect but somehow systematic approach : the matrix integrals.

The Potts model on random maps

Enumerating planar maps

Random matrices

Solving the matrix model

Outline

The Potts model on random maps

Enumerating planar maps

Random matrices

Solving the matrix model

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Combinatorics and Gaussian variables

Wick formula : If (G_1, \cdots, G_{2p}) is a centered Gaussian vector,

$$\mathbb{E}[G_1 G_2 \cdots G_{2p}] = \sum_{\substack{1 \leq s_1 < s_2 \cdots < s_p \leq 2p \\ r_i > s_j}} \prod_{j=1}^p \mathbb{E}[G_{s_j} G_{r_j}].$$

Example : If for all i, $G_i = G$ follow the standard Gaussian distribution,

The Gaussian Unitary Ensemble (GUE)

Let $\mathcal{H}_N = \{A \in \mathcal{M}_{N \times N}(\mathbb{C}); A = A^*\}$. The law μ_N of the GUE is the probability measure on \mathcal{H}_N

$$d\mu_N(A)=\frac{1}{Z_N}e^{-\frac{N}{2}\mathrm{Tr}\,(A^2)}dA.$$

In other words, $A_{lk} = \bar{A}_{kl}$ for $1 \le k < l \le N$ and

$$A_{kl} = (2N)^{-\frac{1}{2}}(g_{kl} + i\tilde{g}_{kl})$$
 for $k < l, A_{kk} = N^{-\frac{1}{2}}g_{kk}$

where the $(g_{kl}, \tilde{g}_{kl}, k \leq l)$ are i.i.d standard Gaussian variables;

$$P(d\tilde{g}_{kl}, dg_{kl}, k \leq l) = \prod_{1 \leq k \leq l \leq N} e^{-\frac{1}{2}(g_{kl})^2} \frac{dg_{kl}}{\sqrt{2\pi}} \prod_{1 \leq k < l \leq N} e^{-\frac{1}{2}(\tilde{g}_{kl})^2} \frac{d\tilde{g}_{kl}}{\sqrt{2\pi}}$$

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The enumeration of maps and one matrix-integrals

For all $d \in \mathbb{N}^*$, all $p \in \mathbb{N}$,

$$\int \left(N\mathrm{Tr} (A^d)\right)^p d\mu_N(A) = \sum_{F \ge 0} \frac{1}{N^{\frac{dp}{2} - p - F}} G(d, p, F)$$

G(d, p, F) = {Union of maps with F faces and p vertices of valence d }.

Recall that a connected graph can be embedded into a surface with Euler characteristic

$$\chi = 2 - 2g =$$
 \$\pmy vertices + \$\pmy faces - \$\pmy edges = p + F - $\frac{pd}{2}$.

Counting is done with labeled vertices and half-edges.

Solving the matrix mode

Application

Wigner (1958) already noticed that,

$$\lim_{N\to\infty}\int \frac{1}{N}\mathrm{Tr}(A^d)d\mu_N(A) = G(d,1,\frac{d}{2}+1) = \begin{cases} 0 & \text{if d is odd,} \\ C_{\frac{d}{2}} & \text{otherwise,} \end{cases}$$

where $C_{\frac{d}{2}}$ is the Catalan number, i.e. the number of rooted planar (g = 0) maps with one vertex of valence d.

(Here, p = 8, F = 5)



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 $G(d, 1, \frac{d}{2} + 1) = \int x^d d\sigma(x)$, where σ is the semi-circular law

The enumeration of colored maps and several matrices integrals

Let $m \in \mathbb{N}$. To any monomial $q(X_1, \dots, X_m) = X_{i_1} \dots X_{i_d}$, we associate (bijectively) a star of type q =

oriented vertex with half-edges of color i_1 , i_2 , ... i_d , ordered clockwise, the first half-edge being marked. Here $q(X) = X_1^2 X_2^2 X_1^4 X_2^2$.



The enumeration of colored maps and several matrices integrals

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For any monomial q, all $p \in \mathbb{N}$,

$$\int N \mathrm{Tr} \ (q(A_1, \cdots, A_m)) d\mu_N(A_1) \cdots d\mu_N(A_m) = \sum_{g \ge 0} \frac{M(q,g)}{N^{2g-2}},$$

 $M(q,g) = \sharp \{ \text{ rooted maps with genus } g \text{ build on one star of type } q \}.$ The matching is only allowed between half-edges of the same color.

Application

Voiculescu (1984) Let $m \in \mathbb{N}$ and $q(X_1, \dots, X_m) = X_{i_1} \cdots X_{i_d}$ for $i_1, \dots, i_d \in \{1, \dots, m\}$

$$\lim_{N\to\infty}\int \frac{1}{N} \operatorname{Tr}\left(q(A_1,\cdots,A_m)\right) d\mu_N(A_1)\cdots d\mu_N(A_m) = G_c(q,1,0)$$
$$:= \sigma_m(q)$$

where $\sigma_m(q)$ is the number of planar maps drawn with a star of type qby gluing half-edges of the same color (Here, $q(\mathbf{X}) = X_1^2 X_2^2 X_1^4 X_2^2$)

 σ_m = law of *m* free semi-circular variables.



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Solving the matrix model

Application to loop models

The Temperley-Lieb elements are boxes with boundary points connected by nonintersecting strings, equipped with a shading and a marked boundary point.

Problem : Take your favorite Temperley-Lieb element. Count the number of planar matching of the end points of the Temperley-Lieb elements so that there are exactly n loops.





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Random matrices and loop enumeration For a Temperley-Lieb element *B*, we denote $p \stackrel{B}{\sim} \ell$ if a string joins the *p*th boundary point with the ℓ th boundary point in *B*, then we associate to *B* with *k* strings the polynomial

$$q_B(X) = \sum_{\substack{i_j = i_p \text{ if } j \stackrel{B}{\sim}_p \\ 1 \leq i_j \leq n}} X_{i_1} \cdots X_{i_{2k}}.$$

$$q_B(X) = \sum_{i,j,k=1}^n X_i X_j X_j X_i X_k X_k \Leftrightarrow$$



Theorem

If ν^{M} denotes the law of n independent GUE matrices,

$$\lim_{M\to\infty}\int\frac{1}{M}\mathrm{Tr}\left(q_B(X)\right)\nu^M(dX)=\sum n^{\sharp\mathrm{loops}}$$

where we sum over all planar maps that can be built on B.

Solving the matrix model



By Voiculescu's theorem, if B =

$$\lim_{M \to \infty} \int \frac{1}{M} \operatorname{Tr} (q_B(X)) \nu^M (dX)$$

=
$$\sum_{i,j,k=1}^n \lim_{M \to \infty} \int \frac{1}{M} \operatorname{Tr} (X_i X_j X_j X_i X_k X_k) \nu^M (dX)$$

$$= \sum \sum_{i,j,k=1}^{n} \sum_{i,j,k=1}^{n}$$
$$= \sum n^{\sharp \text{loops}}$$

because the indices have to be constant along loops.

Problem ; this only gives the generating function at integer values of the fugacity and therefore does not permit to identify it.

Non integer fugacities, $\beta = 0$ [cf Jones 99'] Recall $p \stackrel{B}{\sim} j$ if a string joins the *p*th boundary point with the *j*th boundary point in the TL element B

$$q_B(X) = \sum_{i_j=i_p \text{ if } j \stackrel{B}{\sim} p} X_{i_1} \cdots X_{i_{2k}} \Rightarrow q_B^v(X) = \sum_{e_j=e_p^o \text{ if } j \stackrel{B}{\sim} p} \sigma_B(w) X_{e_1} \cdots X_{e_{2k}}$$

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• e_i edges of a bipartite graph $\Gamma = (V = V_+ \cup V_-, E)$ so that the adjacency matrix of Γ has eigenvalue δ with eigenvector $(\mu_v)_{v \in V}$ with $\mu_v \ge 0$ (\exists for any $\delta \in \{2 \cos(\frac{\pi}{n})\}_{n \ge 3} \cup [2, \infty[)$

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e_i edges of a bipartite graph Γ = (V = V₊ ∪ V₋, E) so that the adjacency matrix of Γ has eigenvalue δ with eigenvector (μ_v)_{v∈V} with μ_v ≥ 0 (∃ for any δ ∈ {2 cos(π/n)}_{n≥3} ∪ [2,∞[)
The sum runs over loops w = e₁ ··· e_{2k} in Γ which starts at v. v ∈ V₊ iff * is in a white region.

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e_i edges of a bipartite graph Γ = (V = V₊ ∪ V₋, E) so that the adjacency matrix of Γ has eigenvalue δ with eigenvector (μ_ν)_{ν∈V} with μ_ν ≥ 0 (∃ for any δ ∈ {2 cos(π/n)}_{n≥3} ∪ [2,∞[)
The sum runs over loops w = e₁ ··· e_{2k} in Γ which starts at v. v ∈ V₊ iff * is in a white region.
If σ(e) := √(μ_{t(e)}/μ_{s(e)}, e = (s(e), t(e)), we set σ_B(e₁ ··· e_{2p}) to be the product Π_{i∈j}/μ_{s(e)} of the σ(e) so that each string of B brings σ(e) with e the edge which labels the start of the string.

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Non integer fugacities

For $e \in E$, e = (s(e), t(e)), X_e^M are independent (except $X_{e^o} = X_e^*) [M\mu_{s(e)}] \times [M\mu_{t(e)}]$ matrices with i.i.d centered Gaussian entries with variance $1/(M_{\sqrt{\mu_{s(e)}\mu_{t(e)}}})$.

Recall
$$q_B^v(X^M) = \sum_{w=e_1\cdots e_{2k}\in L_B \atop s(e_1)=v} \sigma_B(w) X_{e_1}^M \cdots X_{e_{2k}}^M$$

Theorem (G-Jones-Shlyakhtenko 07')

Let Γ be a bipartite graph whose adjacency matrix has δ as Perron-Frobenius eigenvalue. Let B be Temperley-Lieb element so that * is in an unshaded region. Then, for all $v \in V^+$

$$\lim_{M\to\infty} E[\frac{1}{M\mu_{\nu}} \operatorname{Tr}(q_B^{\nu}(X^M))] = \sum \delta^{\sharp \text{loops}}$$

where the sum runs above all planar maps built on B.

Non integer fugacities

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$$\lim_{M\to\infty} E[\frac{1}{M\mu_{\nu}} \operatorname{Tr}(q_B^{\nu}(X^M))] = \sum \delta^{\sharp \text{loops}}$$

where the sum runs above all planar maps built on B. Based on $\sum_{e \in E: s(e) = v} \mu_{t(e)} = \delta \mu_{v}$.

Solving the matrix model

If
$$B =$$

$$\mathbb{E}\left[\frac{1}{M\mu_{v}}\operatorname{Tr}\left(\sum_{e:s(e)=v}\sigma(e)X_{e}X_{e^{0}}\right)\right] = \frac{1}{M\mu_{v}}\sum_{e:s(e)=v}\sqrt{\frac{\mu_{t(e)}}{\mu_{v}}}\frac{M\mu_{v}M\mu_{t(e)}}{M\sqrt{\mu_{t(e)}\mu_{s(e)}}}$$

$$= \frac{1}{\mu_{v}}\sum_{e:s(e)=v}\mu_{t(e)} = \delta$$
If $B =$

$$\lim_{M \to \infty} \mathbb{E}\left[\frac{1}{M\mu_{v}}\operatorname{Tr}\left(\sum_{e:s(e)=v}\sigma(e)\sigma(f)X_{e}X_{e^{0}}X_{f}X_{f^{0}}\right)\right]$$

$$s_{v}^{2} = \frac{1}{2}\sum_{e:s(e)=v}\mu_{t(e)} = s_{v}^{2} = s_{v}^{2}$$

$$= \delta^{2} + \frac{1}{\mu_{\nu}} \sum_{e=f}^{L} \frac{\mu_{t}(e)}{\mu_{\nu}} \frac{\mu_{\nu}\mu_{t}(e)}{\mu_{t}(e)\mu_{\nu}} = \delta^{2} + \delta$$

More generally, the edge is constant along the loop and brings the contribution $\mu_{t(e)}/\mu_{v}$ hence leading after summation to $\delta_{e} = 0$

Random matrices and the enumeration of maps with several vertices

't Hooft noticed in 1974 that matrix integrals are generating functions for the enumeration of maps. If μ_N is the Gaussian law on $N \times N$ Hermitian matrices, we have at list formally

$$\log \int e^{-N \operatorname{Tr}\left(\sum_{i=1}^{p} t_i q_i(A_1, \cdots, A_m)\right)} d\mu_N^{\otimes m}(A_1, \cdots, A_m)$$

$$= \sum_{k_1,...,k_p \in \mathbb{N}} \sum_{g \ge 0} \frac{1}{N^{2g-2}} \prod_{j=1}^p \frac{(-t_j)^{k_j}}{k_j!} M((k_i, q_i)_{1 \le i \le p}; g)$$



where $M((k_i, q_i)_{1 \le i \le p}; g)$ counts maps with genus g build on k_i stars of type q_i so that only half-edges of the same color are matched.

The formal derivation

To prove that formally $\log \int e^{-N \operatorname{Tr}\left(\sum_{i=1}^{p} t_i q_i(A_1, \cdots, A_m)\right)} d\mu_N^{\otimes m}(A_1, \cdots, A_m) \text{ is a generating function for maps : Expand the exponential together with the remarks}$

• For all monomials q_1, \ldots, q_p we have

$$\int \prod_{i=1}^{p} (N \operatorname{Tr}(q_i(A_1, \cdots, A_m))) d\mu_N^{\otimes m}(A_1, \cdots, A_m) = \sum_{g \ge 0} \frac{M(q_1, \ldots, q_p; g)}{N^{2-2g}}$$

where $M(q_1, \ldots, q_p; g)$ is the number of union of maps with genus g which can be build upon one star of type q_1, \ldots , one star of type q_p .

• Taking the logarithm reduces the sum to connected graphs.

First order asymptotics

$$F_{N}(\mathbf{t}) = \frac{1}{N^{2}} \log \int e^{-N \operatorname{Tr}(V_{\mathbf{t}}(A_{1}, \cdots, A_{m}))} d\mu_{N}(A_{1}) ... d\mu_{N}(A_{m}), \ V_{\mathbf{t}} = \sum_{i=1}^{p} t_{i} q_{i}.$$

Theorem Hypothesis : $\phi_{V_t} : (A_1(ij), \dots, A_m(ij))_{i \leq j} \to \operatorname{Tr} (V_t(\mathbf{A}))$ is real-valued. ϕ_{V_t} is convex (or we add a cutoff). For all $\ell \geq 0$, $\exists \varepsilon_{\ell} > 0$ so that if $|\mathbf{t}| = \sum_{i=1}^{p} |t_i| \leq \varepsilon_{\ell}$,

$$F_{N}(\mathbf{t}) = \sum_{g=0}^{\ell} \frac{1}{N^{2g}} \sum_{k_{1}, \cdots, k_{p}} \prod \frac{(-t_{i})^{k_{i}}}{k_{i}!} M((q_{i}, k_{i})_{1 \leq i \leq p}; g) + o\left(\frac{1}{N^{2\ell}}\right)$$

 $M((q_i, k_i)_{1 \le i \le p}; g) = \#\{\text{maps with genus } g \text{ with } k_i \text{ stars of type } q_i\}$

-m = 1: Ambjórn et al. (95), Albeverio-Pastur-Scherbina (01), Ercolani-McLaughlin (03) $-m \ge 2$: G.-Maurel-Segala ($\ell \le 1$ (05)(06)), Maurel-Segala (for all ℓ (06))

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Small t_i 's expansion, more results

Take $V_{\mathbf{t}} = \sum_{i=1}^{p} t_i q_i$. Let $\hat{\mu}_{\mathbf{A}}^{N}$ be the (mean) empirical distribution

$$\hat{\mu}_A^N(P) := \int \frac{1}{N} \operatorname{Tr} \left(P(A_1, \cdots, A_m) \right) d\mu_{V_t}^N(A_1, \cdots, A_m)$$

 $d\mu_{V_{\mathbf{t}}}^{N}(A_{1}, \cdots, A_{m}) = e^{-N^{2}F_{N}(\mathbf{t}) - N\operatorname{tr}(V_{\mathbf{t}}(A_{1}, \cdots, A_{m}))}d\mu_{N}(A_{1})..d\mu_{N}(A_{m})$ Then (same hypothesis as before)

$$\hat{\mu}_{A}^{N}(P) = \sum_{g=0}^{\ell} \frac{1}{N^{2g}} \tau_{g}^{(t_{i})_{1 \le i \le p}}(P) + o(\frac{1}{N^{2\ell}})$$

with

$$\tau_g^{(t_i)_{1\leq i\leq p}}(P) = \sum \prod \frac{(-t_i)^{k_i}}{k_i!} M((q_i, k_i)_{1\leq i\leq p}, (P, 1); g).$$

Proof. g = 0, $V_{\mathbf{t}} = \sum_{i=1}^{n} t_i q_i$

All the limit points $\tau_{\mathbf{t}}$ of

$$\hat{\mu}_A^N(P) := \int \frac{1}{N} \operatorname{tr} \left(P(A_1, \cdots, A_m) \right) d\mu_{V_{\mathbf{t}}}^N(A_1, \cdots, A_m)$$

satisfy Schwinger-Dyson's (or loop) equation

 $\tau_{\mathbf{t}}(X_{i}Q) = \tau_{\mathbf{t}} \otimes \tau_{\mathbf{t}}(\partial_{i}Q) - \tau_{\mathbf{t}}(D_{i}VQ) \quad \forall Q \in \mathbb{C}\langle X_{1}, \cdots, X_{m} \rangle \quad \forall i \in [1, m]$

with $\partial_i P = \sum_{P=P_1X_iP_2} P_1 \otimes P_2$, $D_i P = \sum_{P=P_1X_iP_2} P_2P_1$. By convexity (or cutoff), $\exists R < \infty$ so that $|\tau_t(q)| \le R^{\text{deg } q}$. If $|\mathbf{t}| \le \varepsilon_0$, there is a unique solution

$$\tau_{\mathbf{t}}(q) = \sum_{k_1, \cdots, k_p \in \mathbb{N}} \prod_{i=1}^{p} \frac{(-t_i)^{k_i}}{k_i!} M((q, 1), (q_i, k_i)_{1 \le i \le p}; 0)$$

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How to get Schwinger-Dyson's equation By integration by parts, if $q = q(A_1, \dots, A_m)$

$$\begin{split} \int \operatorname{Tr}(A_i q) d\mu_{V_{\mathbf{t}}}^N &= \sum_{k,\ell} \int [\frac{1}{N} \partial_{A_i(\ell k)} q_{\ell k} - q_{\ell k} \partial_{A_i(\ell k)} \operatorname{tr}(V)] d\mu_{V_{\mathbf{t}}}^N \\ &= \int [\frac{1}{N} \operatorname{Tr} \otimes \operatorname{Tr}](\partial_i q) - \operatorname{Tr}(q \mathcal{D}_i V)] d\mu_{V_{\mathbf{t}}}^N \end{split}$$

By concentration of measure for all monomial P

$$\mu_{V_{\mathbf{t}}}^{N}\left(\left|\frac{1}{N}\mathrm{Tr}\boldsymbol{P}-\mu_{V_{\mathbf{t}}}^{N}(\frac{1}{N}\mathrm{Tr}\boldsymbol{P})\right|\geq\delta\right)\leq Ce^{-cN\delta}$$

and therefore as N goes to infinity any limit point $\tau(P)$ of $\mu_{V_{t}}^{N}|\frac{1}{N}\operatorname{Tr} P, P \in \mathbb{C}\langle X_{1}, \ldots, X_{n} \rangle$ satisfies the Schwinger equation

$$\tau_{\mathbf{t}}(X_i q) = \tau_{\mathbf{t}} \otimes \tau_{\mathbf{t}} \left(\partial_i q \right) - \tau_{\mathbf{t}}(D_i V q)$$

Schwinger-Dyson are equivalent to Tutte equation

 $M((q_1, k_1) \cdots (q_p, k_p), q)$ denotes the number of planar maps build on k_i stars of type q_i and 1 of type $q = X_{i_1} \cdots X_{i_n}$. Tutte surgery shows that $M((q_1, k_1) \cdots (q_p, k_p), q)$ equals

$$\sum_{q=X_{i_1}p_1X_{i_1}p_2}\sum_{\ell_i=0}^{k_i}\frac{k_i!}{\ell_i!(k_i-\ell_i)!}M((q_i,\ell_i)_{1\leq i\leq p},p_1)M((q_i,k_i-\ell_i)_{1\leq i\leq p},p_2)$$

$$+\sum_{i=1}^{p}k_{i}\sum_{q_{i}=q_{i}^{1}X_{i_{1}}q_{i}^{2}}M((q_{1},k_{1})\cdots(q_{i},k_{i}-1)\cdots(q_{p},k_{p}),(qq_{i}^{2}q_{i}^{1},1))$$

Put $\tau_{\mathbf{t}}(q) = \sum \prod \frac{(-t_i)^{k_i}}{k_i!} M((q_1, k_1) \cdots (q_p, k_p), (q, 1))$. Then, summation shows that $\tau_{\mathbf{t}}$ satisfies the Schwinger-Dyson equation

$$\tau_{\mathbf{t}}(X_{i}Q) = \tau_{\mathbf{t}} \otimes \tau_{\mathbf{t}}(\partial_{i}Q) - \tau_{\mathbf{t}}(D_{i}VQ)$$

Advantage : we know that τ_t is a tracial state!

Matrix models for loop models

Let B_i be Temperley Lieb elements with * with color $\sigma_i \in \{+, -\}$, $1 \leq i \leq p$. Let Γ be a bipartite graph whose adjacency matrix has eigenvalue δ as before. Let ν^M be the law of the previous independent rectangular Gaussian matrices and set

$$d\nu^{M}_{(B_{i})_{i}}(X_{e}) = \frac{1_{\|X_{e}\|_{\infty} \leq L}}{Z^{M}_{B}} e^{\mathsf{Mtr}(\sum_{i=1}^{p} \beta_{i} \sum_{v \in V_{\sigma_{i}}} \mu_{v} q^{v}_{B_{i}}(X))} d\nu^{M}(X_{e}).$$

Theorem (G-Jones-Shlyakhtenko-Zinn Justin 10') For any L > 2, for β_i small enough real numbers, for any Temperley-Lieb element B with color σ , any $v \in V_{\sigma}$,

$$\lim_{M\to\infty}\int \frac{1}{M\mu_{\nu}}tr(q_{B}^{\nu}(X))d\nu_{(B_{i})_{i}}^{M}(X)=\sum_{n_{i}\geq0}\sum\delta^{\sharp \text{loops}}\prod_{i=1}^{p}\frac{\beta_{i}^{n_{i}}}{n_{i}!}$$

where we sum over the planar maps build on n_i TL elements B_i and one B.

Matrix model for the Potts model

Let $\delta \in \{2\cos(\frac{\pi}{n})\}_{n\geq 3} \cup [2,\infty[\text{ and } \Gamma = (V_+ \cup V_-, E) \text{ be a bipartite graph with eigenvalue } \delta \text{ and Perron-Frobenius eigenvector } \mu$.

$$\nu_{\beta\pm}^{M}(dX_{e}) = \frac{1_{\|X_{e}\|_{\infty} \leq L}}{Z_{\beta\pm}^{M}} e^{\mathsf{Mtr}\left(\sum_{v \in V} \mu_{v} \sum_{\sigma=\pm} \beta_{\sigma} \mathbf{1}_{v \in V_{\sigma}} \left(\sum_{e:s(e)=v} \sigma(e) X_{e} X_{e}^{*}\right)^{2}\right)}$$
$$\prod_{e} e^{-\frac{M}{2} \left(\mu_{s(e)} \mu_{t(e)}\right)^{\frac{1}{2}} \mathsf{tr}(X_{e} X_{e}^{*})} dX_{e} dX_{e}^{*}$$

Theorem (G-Jones-Shlyakhtenko-Zinn Justin 10') Then, for L large enough, β_{\pm} small enough, for all TL B, $v \in V_{\sigma_B}$,

$$\lim_{M \to \infty} \frac{1}{M\mu_{\nu}} \int tr(q_{B}^{\nu}(X)) \nu_{\beta_{\pm}}^{M}(dX)$$
$$= \sum_{n_{+}, n_{-} \ge 0} \sum \prod_{i=\pm} \frac{\beta_{i}^{n_{i}}}{n_{i}!} \delta^{\sharp \text{loops}} =: \text{Tr}_{\beta_{\pm}, \delta}(B)$$

The Potts model on random maps

Enumerating planar maps

Random matrices

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Solving the matrix model

Outline

The Potts model on random maps

Enumerating planar maps

Random matrices

Solving the matrix model

Solving the matrix model

Recap about the Potts model

Assume for simplicity Γ finite hereafter (which includes $\delta = 2\cos(\pi/n), n \ge 3$). By construction



equals

$$\lim_{M\to\infty}\frac{1}{M^2\sum \mu(v)^2}\log Z^M_{\beta\pm}$$

with for L large enough $Z^M_{\beta\pm}$ equals to

$$\int_{\|X_e\| \le L} e^{M \operatorname{Tr} \left(\sum_{v \in V} \mu_v (\beta_+ 1_{v \in V_+} + \beta_- 1_{v \in V_-}) \left(\sum_{e: s(e) = v} \sqrt{\mu_v \mu_{t(e)}} X_e X_e^* \right)^2 \right)} \nu^M (dX)$$

if ν^{M} is the law of X_{e} , $[M\mu_{s(e)}] \times [M\mu_{t(e)}]$ matrices with iid centered Gaussian entries with covariance $(M^{2}\mu_{s(e)}\mu_{t(e)})^{-\frac{1}{2}}$

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Solving the matrix model :HS transformation With $G_{\nu} [M\mu_{\nu}] \times [M\mu_{\nu}]$ independent matrices from the GUE and $X_e [M\mu_{s(e)}] \times [M\mu_{t(e)}]$ matrices with covariance $(M^2\mu_{s(e)}\mu_{t(e)})^{-\frac{1}{2}}$ under ν^M , with $\alpha_{\pm} = \sqrt{2\beta_{\pm}}$, $\Delta(\lambda) = \prod_{i \neq j} (\lambda_i - \lambda_j)$, $Z^M_{\beta_{\pm}} =$

$$= \int_{\|X_{e}\| \leq L} e^{M \operatorname{Tr} \left(\sum_{v \in V} \mu_{v} (\beta_{+} 1_{v \in V_{+}} + \beta_{-} 1_{v \in V_{-}}) \left(\sum_{e:s(e)=v} \sigma(e) X_{e} X_{e}^{*} \right)^{2} \right)} \nu^{M}(dX)$$

$$= \int_{\|X_{e}\| \leq L} e^{M \operatorname{Tr} \left(\sum_{\sigma=\pm} \sum_{v \in V_{\sigma}} \alpha_{\sigma} G_{v} \left(\sum_{e:s(e)=v} \sqrt{\mu_{v} \mu_{t(e)}} X_{e} X_{e}^{*} \right) \right)} \nu^{M}(dX, dG)$$

$$\approx \int_{\|G_{v}\| \leq L'} \prod_{e \in E_{+}} e^{-\operatorname{Tr} \otimes \operatorname{Tr} \left(\log(I + \alpha_{+} I \otimes G_{s(e)} + \alpha_{-} G_{t(e)} \otimes I) \right)} \nu^{M}(dG)$$

$$= \int_{|\lambda_{i}^{e}| \leq L'} \prod_{e \in E_{+}} e^{-\sum_{i=1}^{[M \mu_{s(e)}]} \sum_{j=1}^{[M \mu_{t(e)}]} \log(1 + \alpha_{+} \lambda_{i}^{e} + \alpha_{-} \eta_{j}^{e})}}$$

$$\Delta(\eta^{e}) \Delta(\lambda^{e}) e^{-\frac{[M \mu_{s(e)}]}{2} \sum_{i=1}^{[M \mu_{s(e)}]} (\lambda_{i}^{e})^{2} - \frac{[M \mu_{t(e)}]}{2} \sum_{j=1}^{[M \mu_{t(e)}]} (\eta_{j}^{e})^{2}} d\lambda^{e} d\eta^{e}.$$

The auxiliary matrix model :large deviations Let $P_{\alpha}^{M,L}$ be absolutely cont. wrt Lebesgue with density

$$\frac{1_{|\lambda^{\nu}|\leq L}}{Z_{\alpha}^{M,L}}\prod_{e\in E_{+}}\prod_{\substack{1\leq i\leq [M\mu_{s(e)}]\\1\leq j\leq [M\mu_{t(e)}]}}\frac{1}{1+\alpha_{+}\lambda_{i}^{s(e)}+\alpha_{-}\lambda_{j}^{t(e)}}\prod_{\nu\in V}\Delta(\lambda^{\nu})e^{-\frac{M\mu_{\nu}}{2}\sum(\lambda_{i}^{\nu})^{2}}.$$

By large deviation analysis,

$$P_{\alpha}^{M,L}\left(d(\frac{1}{M\mu_{\nu}}\sum \delta_{\lambda_{i}^{\nu}},\nu_{\nu})<\epsilon\,\forall\,\nu\right)\approx e^{-M^{2}[I(\nu_{\nu},\nu\in V)-\inf I]}$$

with if $\Sigma(\nu) = \int \log |x - y| d\nu(x) d\nu(y)$

$$I(\nu_{\nu},\nu\in V)=\sum_{\nu}\frac{\mu_{\nu}^{2}}{2}(\int x^{2}d\nu_{\nu}(x)-2\Sigma(\nu_{\nu}))$$

$$-\sum_{e\in E_+}\mu_v\mu_{t(e)}\int \log|1+\alpha_+x+\alpha_-y|d\nu_v(x)d\nu_{t(e)}(y).$$

Moreover $\lim_{M\to\infty} M^{-2} \log Z^{M,L}_{\alpha} = -\inf I$

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The auxiliary matrix model (Kostov, 95')

Let Γ be as before and consider the Gibbs measure $P_{\alpha}^{M,L}$ absolutely cont. wrt Lebesgue with density

$$\frac{1_{|\lambda^{\nu}|\leq L}}{Z_{\alpha}^{M,L}}\prod_{e\in E_{+}}\prod_{\substack{1\leq i\leq [M\mu_{\mathfrak{s}(e)}]\\1\leq j\leq [M\mu_{\mathfrak{t}(e)}]}}\frac{1}{1+\alpha_{+}\lambda_{i}^{\mathfrak{s}(e)}+\alpha_{-}\lambda_{j}^{\mathfrak{t}(e)}}\prod_{\nu\in V}\Delta(\lambda^{\nu})e^{-\frac{M\mu_{\nu}}{2}\sum(\lambda_{i}^{\nu})^{2}}.$$

Theorem [G-J-S-ZJ 10']• I achieves its minimal value at a unique set of probability measures $\nu_v, v \in V$.

- $\exists \nu_+, \nu_- \in P(\mathbb{R})$, so that $\nu_v = \nu_{\pm}$ if $v \in V_{\pm}$.
- For all L > 2 not too large,

$$\lim_{M\to\infty} E[\frac{1}{M\mu_{v}}\sum_{i=1}^{[M\mu_{v}]} (\lambda_{i}^{v})^{p}] = \int x^{p} d\nu_{v}(x) \quad \forall \quad p \in \mathbb{N}, \quad v \in V.$$

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More about the relation between Potts model and its auxiliary model

Let $M(z) = \int \sum_{n\geq 0} z^n x^n d\nu_+(x)$ and put

 $\gamma(z) = \alpha_+ z / (1 - z^2 M(z))$

and

$$C(z, A, B) = \sum_{n \ge 0} z^n \sum \delta^{\ell} \frac{\beta_+^{n_+}}{n_+!} \frac{\beta_-^{n_-}}{n_-!}$$

where we sum over the planar maps build over n_+ (resp. n_-) vertices with two strings and two black (resp. white) regions and one vertex of type $B_n =$ Then, for small z,

$$C(z,A,B) = \frac{\alpha_+}{z} \left[1 - \frac{\alpha_+ \gamma^{-1}(z)}{z}\right].$$

Study of u_-, u_+

 $(
u_-,
u_+)$ are the unique minimizers of,

$$\sum_{\varepsilon=\pm} \left(\frac{1}{2} \int x^2 d\nu_{\varepsilon}(x) - \int \log |x - y| d\nu_{\varepsilon}(x) d\nu_{\varepsilon}(y) \right)$$

$$+\delta\int \log|1+lpha_+x+lpha_-y|d
u_+(x)d
u_-(y).$$

Let p_+ (resp. p_-) be the law of $1 + \alpha_+ x$ and $-\alpha_- y$ under ν_+ (resp. ν_-).

•For α_{\pm} small enough, p_{\pm} has a connected support $[a_{\pm}, b_{\pm}]$ around 1 (resp. 0) and $a_{-} < b_{-} < a_{+} < b_{+}$.

• Set $G_{\pm}(z) = \int (z - x)^{-1} dp_{\pm}(x)$. Then

 $G_{\pm}(z+i0) + G_{\pm}(z-i0) = P_{\pm}(z) + \delta G_{\mp}(z)$ $z \in [a_{\pm}, b_{\pm}]$

with $P_{-}(z) = z/\alpha_{-}$, $P_{+}(z) = (1 - z)/\alpha_{+}$. • If $q = 2 \cos \pi/n$, $n \ge 3$, G_{\pm} satisfy an algebraic equation. Exact solution for $G_{\pm}(z) = \int (z-x)^{-1} d\nu_{\pm}(z)$ Introduce

$$u(z) = \int_{b_-}^{z} \frac{1}{\sqrt{(v-a_+)(v-a_-)(v-b_+)(v-b_-)}} dv,$$

with inverse z(u). With $\delta = q + q^{-1}$ set

$$\omega_{\pm}(u) = q^{\pm 1}G_{+}(z(u)) - G_{-}(z(u)) \pm \frac{1}{q - q^{-1}}(P_{+} + q^{\pm 1}P_{-})z(u).$$

Then $\omega_{\pm}(u + 2K) = \omega_{\pm}(u)$ $\omega_{\pm}(u + 2iK') = q^{\pm 2}\omega_{\pm}(u)$

Theorem

 ω_\pm are meromorphic with only poles at $\pm u_\infty$.Set

$$\Theta(u) = 2\sum_{k=0}^{\infty} e^{i\pi\frac{\omega_2}{\omega_1}(k+1/2)^2} \sin(2k+1)\frac{\pi u}{\omega_1}$$
$$\omega_+(u) = c_+ \frac{\Theta(u-u_\infty - \nu\omega_1)}{\Theta(u-u_\infty)} + c_- \frac{\Theta(u+u_\infty - \nu\omega_1)}{\Theta(u+u_\infty)}$$

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Conclusion

- Loop models can be represented by matrix models for a continuum of fugacities δ .
- Therefore the computation of the matrix model allows to identify the generating function for all δ by analyticity.
- The matrix model for the shaded O(n) model (or the Potts model on random graphs) can be computed.