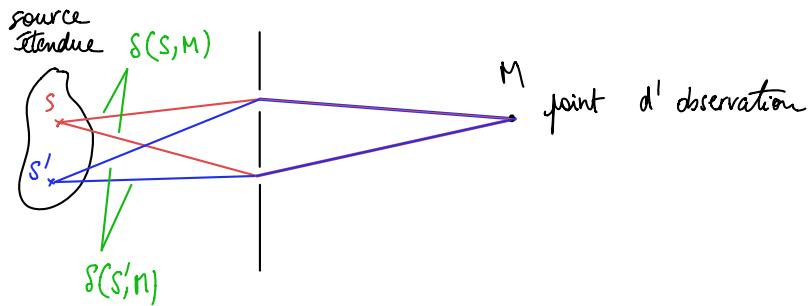


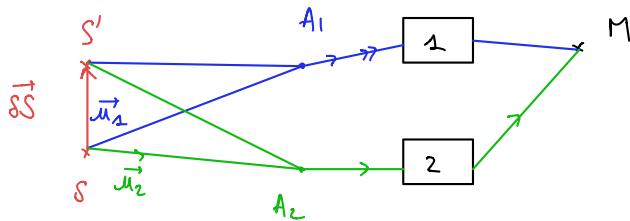
I)

1.



2.

interféromètre à deux voies stigmatiques



$$\cdot \delta(S, M) = (S\pi)_2 - (SM)_2 = SA_2 + L_2 - SA_1 - L_1 \\ \text{idem } S' \leftrightarrow S$$

$$\cdot \delta(S', M) - \delta(S, M) = (S'A_2 - SA_2) - (S'A_1 - SA_1)$$

$$\text{or } d(SA_n) = d(\vec{SA}_n \cdot \vec{\mu}_n) = \underbrace{(d\vec{SA}_n \cdot)}_{\vec{SA}'_n - \vec{SA}_n} \vec{\mu}_n + \underbrace{\vec{SA}_n d\vec{\mu}_n}_{\frac{1}{2} d(\vec{\mu}_n^2)}$$

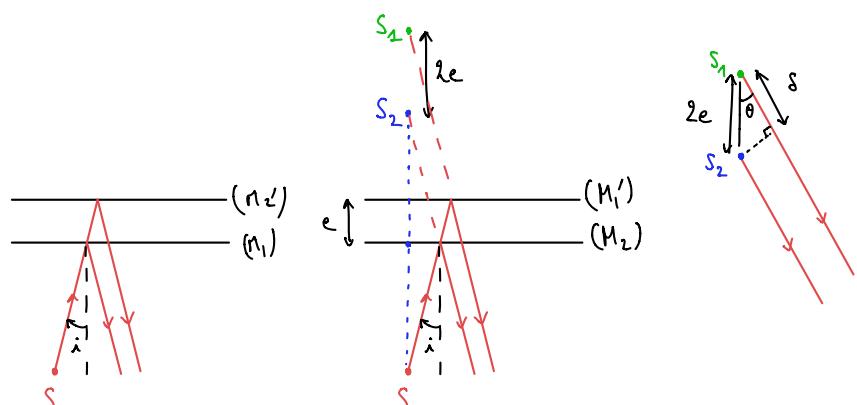
$$= \vec{SS}' + \underbrace{0}_{SA_n} \underbrace{((S'S)^2)}_{\vec{SA}_n \vec{\mu}_n \cdot d\vec{\mu}_n}$$

$$d(SA_n) = \vec{SS}' \cdot \vec{\mu}_n$$

$$\rightarrow \boxed{\delta(S', M) - \delta(S, M) = \vec{SS}' (\vec{\mu}_2 - \vec{\mu}_1) + O\left(\frac{(\vec{SS}')^2}{SA_n}\right)}$$

I)

2. lame d'air



déférence de marche : par stigmatisme des miroirs, tout se passe comme si les rayons émergents venaient de 2 sources ponctuelles  $S_1$  et  $S_2$

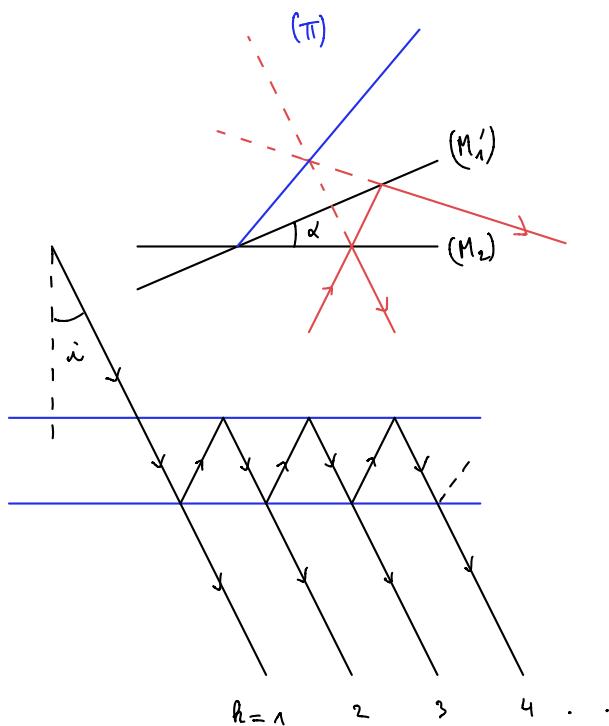
$$\rightarrow \boxed{\delta = 2ne \cos i}$$

- formule de Fresnel :  $I(i) = \frac{I_{\max}}{2} (1 + \cos(\frac{2\pi}{\lambda} e_m \cos(i)))$

- doublet du sodium :  $I(i) = \frac{I_{\max}}{2} (1 + \cos(\frac{2\pi}{\lambda_1} \delta)) + \frac{I_{\max}}{2} (1 + \cos(\frac{2\pi}{\lambda_2} \delta))$   
 $= I_{\max} (1 + \frac{1}{2} (\cos(\frac{2\pi}{\lambda_1} \delta) + \cos(\frac{2\pi}{\lambda_2} \delta))$   
 $= I_{\max} (1 + \underbrace{\cos(\frac{(k_1+k_2)}{2} \delta)}_{\text{franges}} \underbrace{\cos(\frac{(k_1-k_2)}{2} \delta)}_{\text{contraste}})$   
 où  $\frac{k_1 - k_2}{2} = \frac{\Delta k}{2} = \frac{1}{2} \Delta \left( \frac{2\pi}{\lambda} \right) \approx \frac{1}{2} \cdot \frac{2\pi \Delta \lambda}{\lambda^2}$

au centre,  $\delta=2e$ , contraste  $C(e) = \cos\left(2\pi \frac{\Delta \lambda}{\lambda^2} e\right)$   $e_n = \frac{1}{2} \left( \frac{1}{2} + m \right) \frac{\Delta \lambda^2}{\Delta \lambda}$

### 3. Caisson d'air



3. 1.

amplitude complexe

$E_1 = t_{12} t_{21} E_0 \quad \leftarrow \text{référence de phase idem Richardson}$

$E_2 = t_{12} r_{12} e^{i\Delta\ell} r_{12} t_{21} E_0 \quad \text{où } \Delta\ell = \int \frac{2\pi}{\lambda} \cdot 2m \cos i$

$E_k = t_{12} (r_{12} e^{i\Delta\ell} r_{12})^{k-1} t_{21} E_0$

$E_{\text{Tot}} = t_{12} t_{21} E_0 \cdot \sum_{k=1}^{\infty} (r_{12} e^{i\Delta\ell} r_{12})^{k-1}$

$= \frac{t_{12} t_{21} E_0}{1 - r_{12}^2 e^{i\Delta\ell}}$

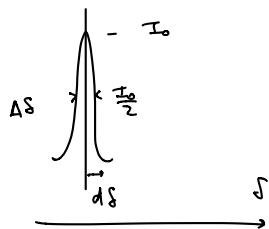
$I_{\text{Tot}} = \frac{(1-R)^2 I_0}{1 + R^2 - 2R \cos \Delta\ell} = \frac{I_0}{1 + \frac{2R(1-\cos \Delta\ell)}{(1-R)^2}}$

$$I_{TOT} = \frac{I_0}{1 + \frac{4R}{(1-R)^2} \cdot \sin^2\left(\frac{\Delta\Phi}{2}\right)}$$

$$\cos 2x = \cos^2 - \sin^2 \\ < 1 - 2 \sin^2$$

2. Pouvoir de résolution :  $PR = \frac{\lambda}{\Delta\lambda}$  où  $\delta = p\lambda$  à un pic

on cherche  $d\delta$  tel que  $I = \frac{I_0}{2}$  ie  $m \sin^2(p\pi + \frac{k d\delta}{2}) = 1$



$$\sin^2(p\pi + \varepsilon) = \sin^2(\varepsilon) \underset{\varepsilon \rightarrow 0}{\sim} \varepsilon$$

$$d\delta \sim \frac{2}{k} \cdot \frac{1}{\sqrt{m}} = \frac{2}{\pi \sqrt{m}} = p d\lambda$$

$$\hookrightarrow PR = p \frac{\pi \sqrt{m}}{2} = p F$$

$$\Delta\delta = 2 d\delta = p \Delta\lambda$$

autre résonance  
si  $\lambda \rightarrow \lambda + d\lambda$