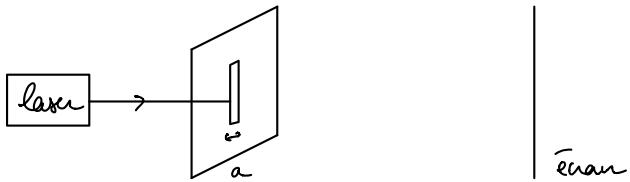


Intro



diffraction : $\lambda \gtrsim a$

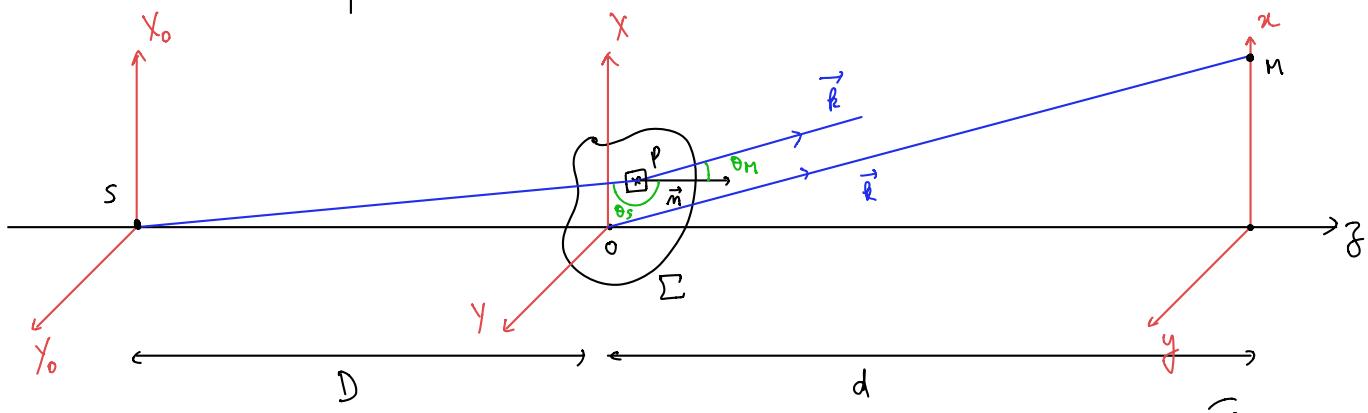
I)

1.

onde
incidente



$$t(x, y) = \frac{A(x, y, z=0^+)}{A(x, y, z=0^-)}$$



$$s(M) = \sum_{\text{objet}} d^{2p} \cdot K(\theta_s, \theta_M) \cdot \underbrace{s_0 \frac{e^{ikSP}}{SP}}_{s_0(p)} \cdot \underbrace{t(p)}_{\text{objet}} \cdot \underbrace{\frac{e^{ikMP}}{MP}}_{H-F}$$

2. Cadre : $\left. \begin{array}{l} \text{source lointaine} \\ \text{écran lointain} \end{array} \right\} D, d \gg x, y, X, Y$

(i) rayons peu inclinés $\rightarrow K(\theta_s, \theta_M) = K$ constante

$$(ii) MP = \sqrt{(x-X)^2 + (y-Y)^2 + d^2} \approx d \left(1 + \frac{1}{2} \frac{(x-X)^2}{d^2} + \frac{1}{2} \frac{(y-Y)^2}{d^2} + O\left(\frac{(x-X)^4}{d^4}, \frac{(y-Y)^4}{d^4}\right) \right)$$

$$\rightarrow \text{diminution } \frac{1}{MP} \approx \frac{1}{d}$$

\rightarrow dans l'exponentielle, il faut comparer les variations de MP à la longueur d'onde

(iii) idem avec SP

$$\begin{aligned} x &\leftrightarrow X_0 = 0 \\ y &\leftrightarrow Y_0 = 0 \\ d &\leftrightarrow D \end{aligned}$$

$$s(x,y) = \underbrace{\frac{ikS_0}{dD} e^{ikd}}_C \cdot \underbrace{e^{ikd} e^{ikD} \cdot e^{ik\left(\frac{x^2+y^2}{2d}\right)}}_{\text{ignoré car on observe } |s|^2} \left| \int \int dX dY t(X,Y) \cdot e^{-ik\left(\frac{xX+yY}{d}\right)} + ik \frac{(x^2+y^2)}{2} \left(\frac{1}{d} + \frac{1}{D}\right) \right|$$

(a)

Diffracton du Fraunhofer: le terme quadratique (a) de phase peut être négligé.

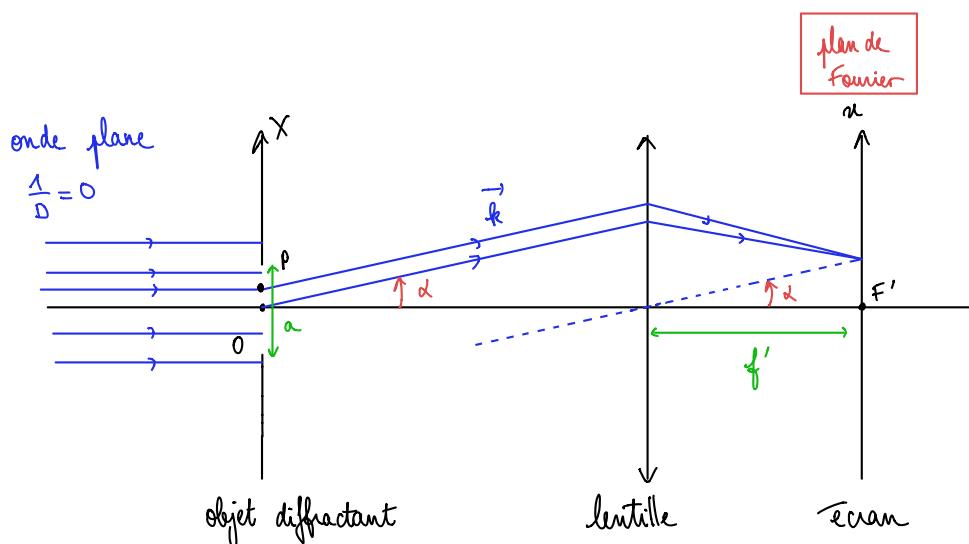
$$\begin{aligned} \text{Dans ce cas, } s(x,y) &= C \cdot \left| \int \int dX dY t(X,Y) \cdot e^{-ik(\theta_x X + \theta_y Y)} \right| \\ &= C \cdot \text{TF}[t](q_x, q_y) \end{aligned}$$

$$\theta_x = \frac{x}{d}, \quad \theta_y = \frac{y}{d}$$

angles sous lesquels P voit l'objet diffractant

$$q_x = k\theta_x \quad q_y = k\theta_y \quad \text{fréquences spatiales}$$

3. (i) onde plane ($\frac{1}{D} = 0$) à l'infini ($\frac{1}{d} = 0$) au foyer d'une lentille.

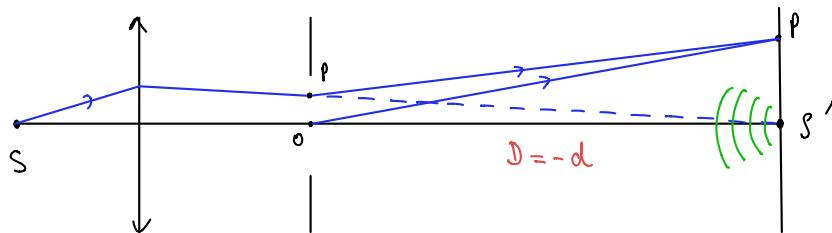


(ii) onde plane ($\frac{1}{D} = 0$) en régime de Fraunhofer

$$N_F = \frac{\alpha^2}{\lambda d} \ll 1$$

(iii) au voisinage de l'image géométrique

$$\frac{1}{D} + \frac{1}{d} = 0$$



translation - modulation

$$t(x+x_0) \xrightarrow{\text{TF}} e^{jfx_0} \text{TF}[t](f)$$

similitude

$$t(ax) \xrightarrow{\text{TF}} \frac{1}{|a|} \text{TF}[t]\left(\frac{f}{a}\right)$$

linéarité \$\Rightarrow\$

$$1-t \xrightarrow{\text{TF}} S(f) - \text{TF}[t](f)$$

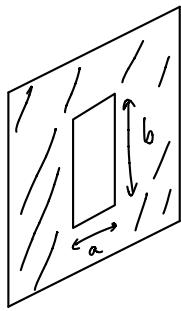
produit - convolution

$$f * g \xrightarrow{\text{TF}} \text{TF}[f] * \text{TF}[g]$$

itération

$$f(x) \xrightarrow{\text{TF}} \xrightarrow{\text{TF}} \text{TF}[\text{TF}[f]](x) = f(-x)$$

2. 1.



éclairage par une onde plane

$$\begin{aligned}
 A(x, y) &= C \int_{-\alpha/2}^{\alpha/2} dx \int_{-\beta/2}^{\beta/2} dy e^{-iq_x x} e^{-iq_y y} \\
 &= C \left[\frac{e^{-iq_x x}}{-iq_x} \right]_{-\alpha/2}^{\alpha/2} \left[\frac{e^{-iq_y y}}{-iq_y} \right]_{-\beta/2}^{\beta/2} \\
 &= C \left(\frac{e^{-iq_x a} - e^{+iq_x a}}{-iq_x} \right) \left(\frac{e^{-iq_y b} - e^{+iq_y b}}{-iq_y} \right) \\
 &= C \operatorname{sinc}\left(\frac{q_x a}{2}\right) \operatorname{sinc}\left(\frac{q_y b}{2}\right)
 \end{aligned}$$

$$I(x, y) = I_0 \operatorname{sinc}^2\left(\frac{k_x a}{2d}\right) \operatorname{sinc}^2\left(\frac{k_y b}{2d}\right)$$

$$I=0 \text{ lorsque } \pi = \frac{k_x a}{2d} \text{ ie } \boxed{\theta_n = \frac{2\pi}{ka} = \frac{\lambda}{a}}$$

on retrouve la formule du lysée

2. $s(x, y) = C \operatorname{TF}(t)[q_x, q_y]$

TF)

1.

$$t_{\text{bifente}} = t_{\text{fente}} * \left(S_{-\frac{a}{2}} + S_{\frac{a}{2}}\right) \rightarrow \hat{t}_{\text{bifente}}(q_x) \propto \hat{t}_{\text{fente}}(q_x) \cdot \cos\left(\frac{a}{2} q_x\right)$$

modulé par la fig
diffaute 1 fente

$$2\pi \frac{a}{2} \cdot \frac{\lambda}{2d}$$

interférage, $i = \frac{\lambda d}{a}$
car $|\hat{f}|^2$ a la période
double de t

