



# Dynamics of glassy and jammed colloidal systems

L. Cipelletti

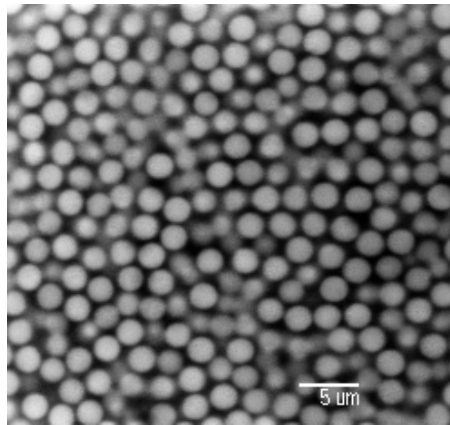
*L2C, Université Montpellier 2 and CNRS*

Collaborators: L. Berthier (LCVN), V. Trappe (Fribourg)

Students: P. Ballesta, G. Brambilla, A. Duri, D. El Masri

Postdocs: S. Maccarrone, M. Pierno

# Soft glassy/jammed materials

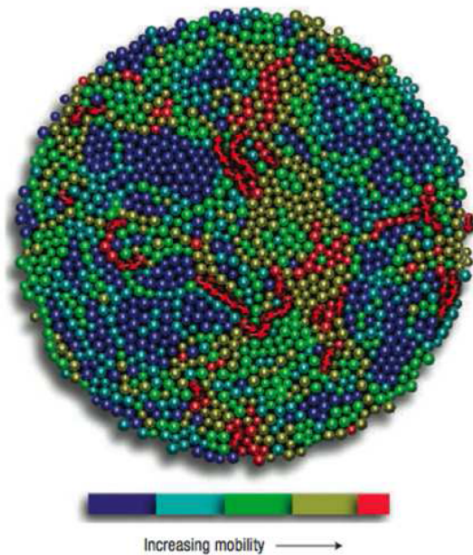


E. Weeks



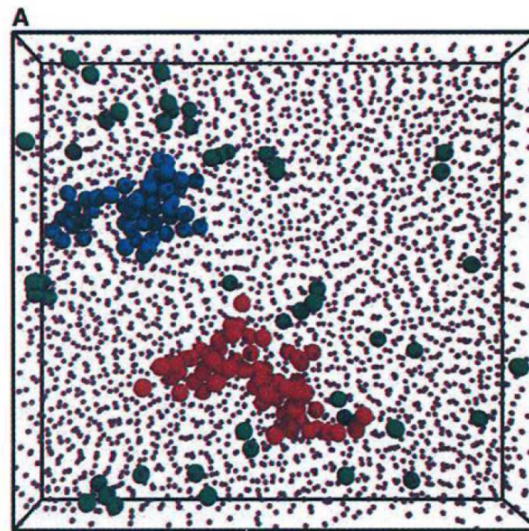
# Dynamical heterogeneity is ubiquitous!

Granular matter



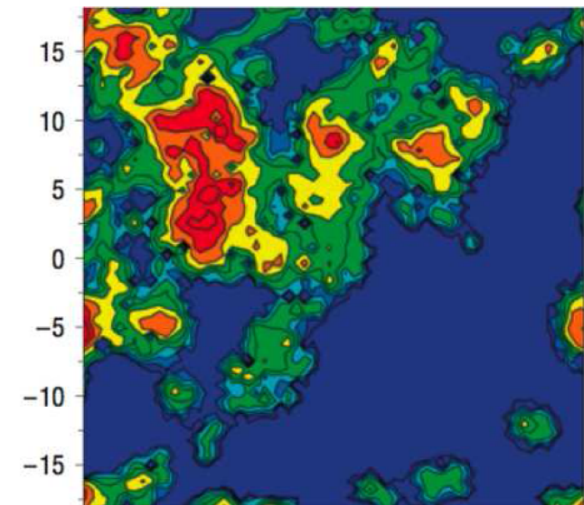
*Keys et al. Nat. Phys. 2007*

Colloidal Hard Spheres



*Weeks et al. Science 2000*

Repulsive disks



*A. Widmer-Cooper Nat. Phys. 2008*

# Outline

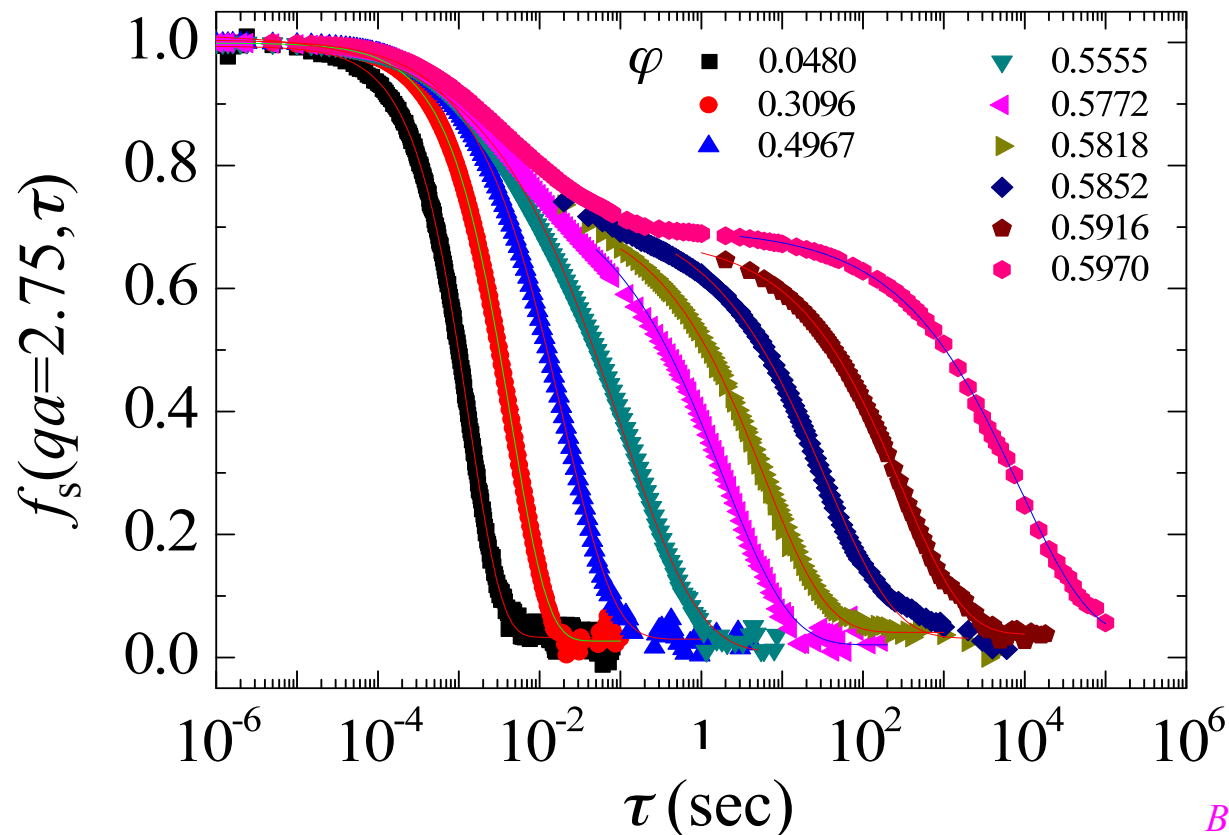
- **Average dynamics and dynamical heterogeneity of supercooled colloidal HS**
- **Temporal fluctuations of the dynamics in jammed/glassy materials**
- **Real-space measurements of correlations: ultra-long range spatial correlations and elasticity**

# Equilibrium dynamics of supercooled HS

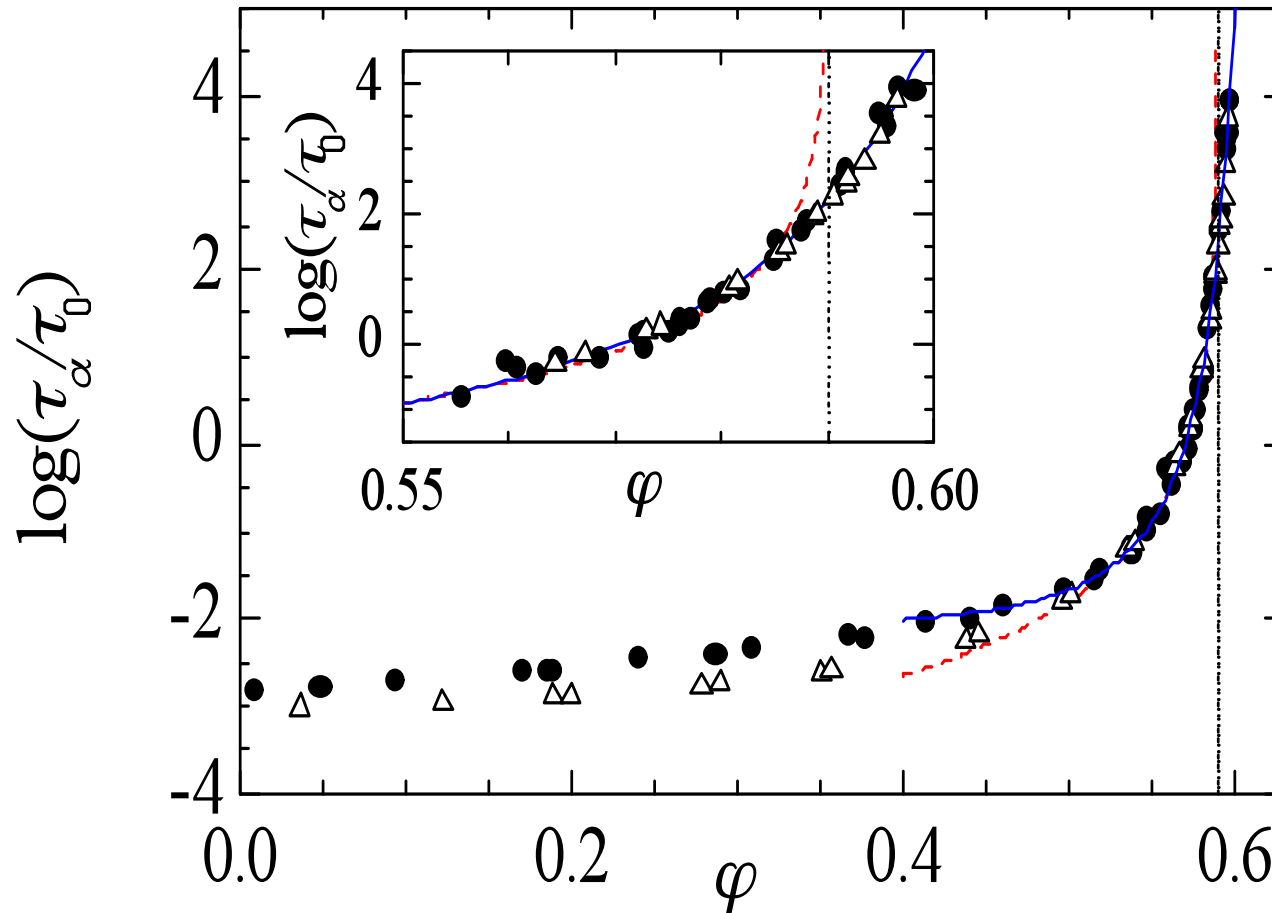
Multispeckle DLS

$$f_s(q, \tau) = \left\langle \frac{1}{N} \sum_{j=1}^N \exp[-i\mathbf{q} \cdot (\mathbf{r}_j(0) - \mathbf{r}_j(\tau))] \right\rangle$$

- PMMA in decalin/tetralin
- $a \sim 110$  nm
- polydispersity 12.2% (TEM)



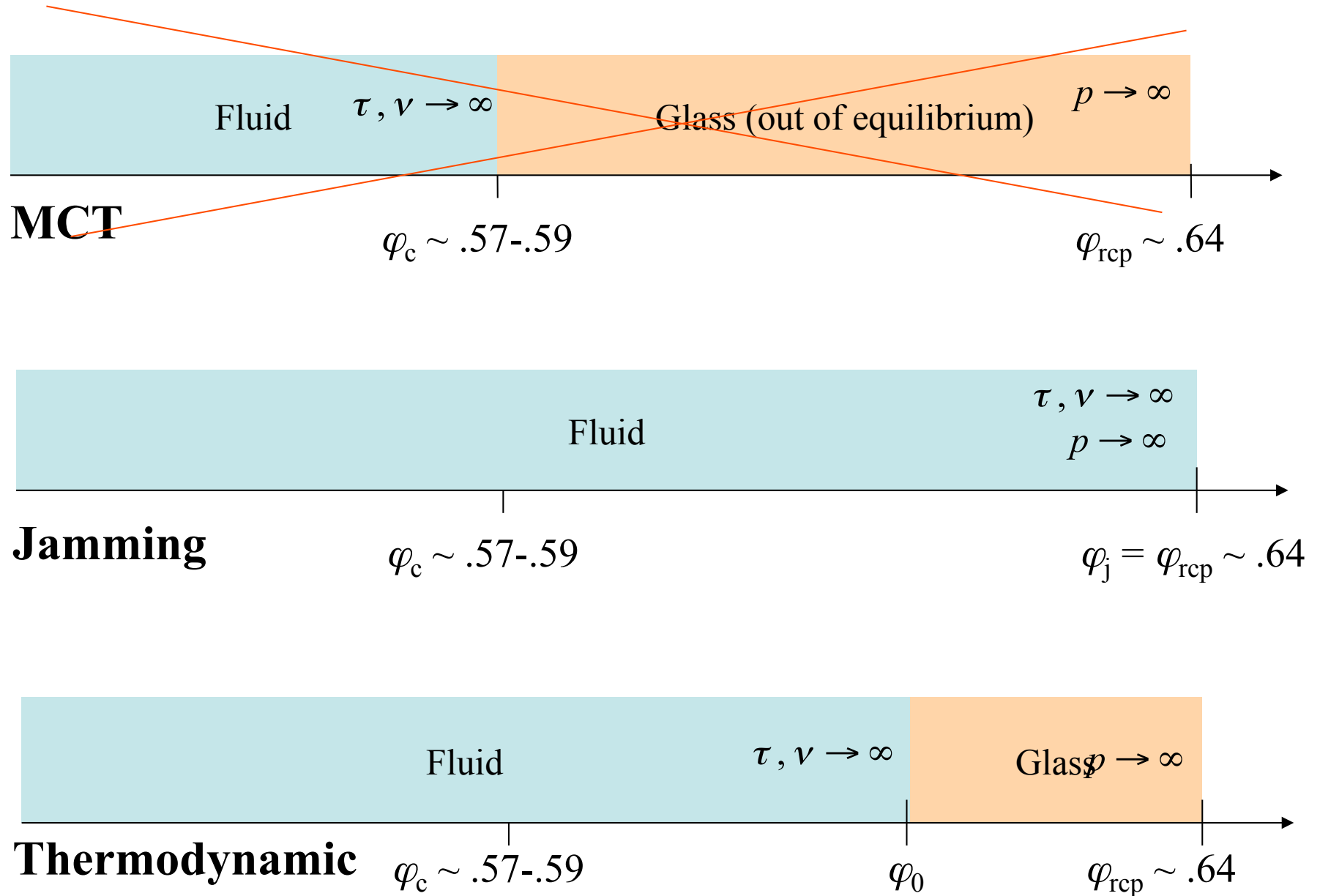
# $\varphi$ dependence of the relaxation time



MCT:  $\tau_\alpha \sim (\varphi_c - \varphi)^{-\gamma}$   
 $\varphi_c = 0.59, \gamma = 2.6$

VFT:  $\tau_\alpha \sim \exp[(\varphi_0 - \varphi)^{-\delta}]$   
 $\varphi_0 = 0.614, \delta = 1$   
 $\varphi_0 = 0.637, \delta = 2$

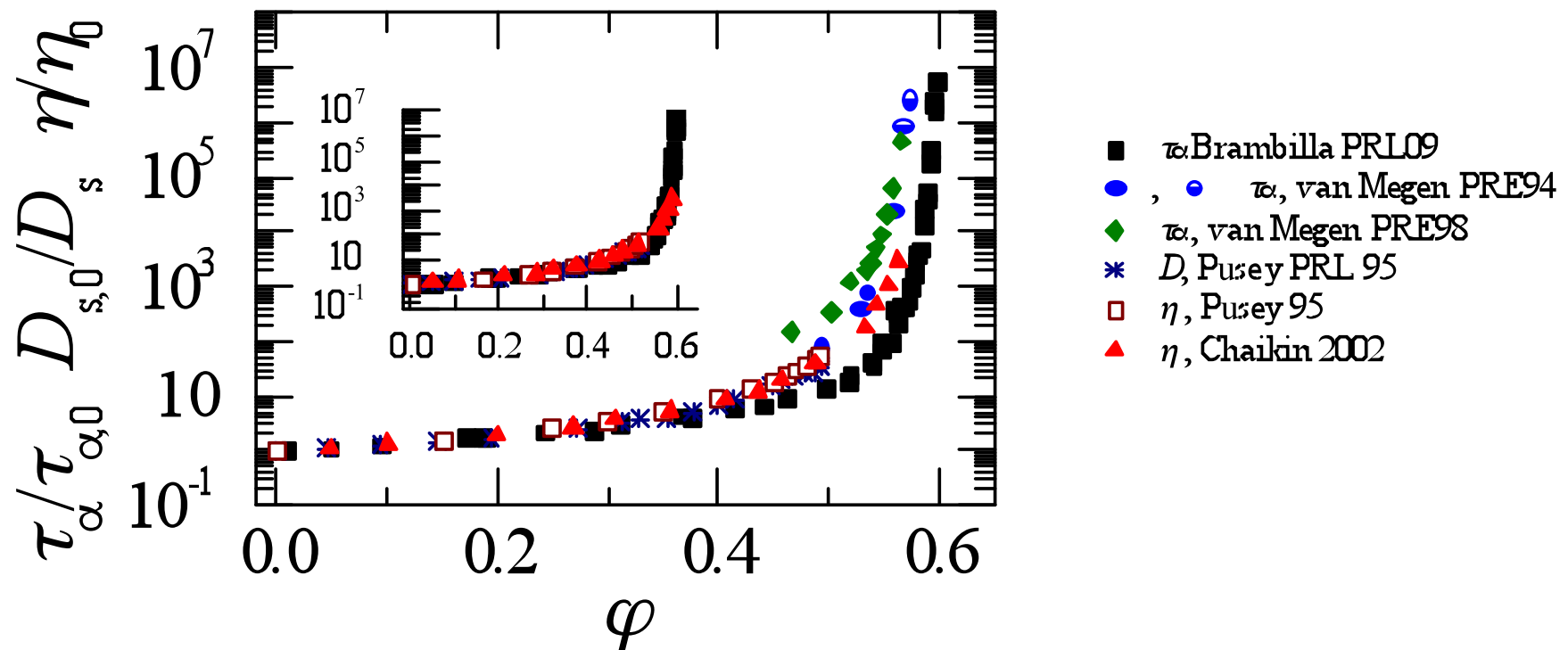
# Glass/jamming transition in colloidal HS





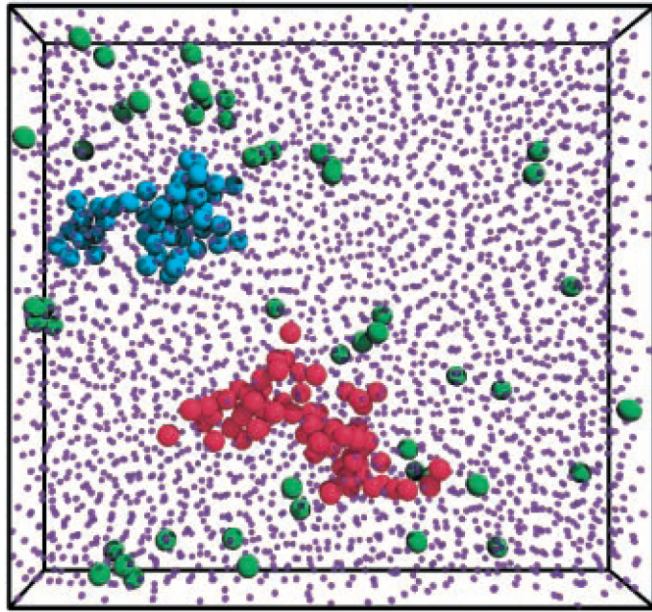
# Something experimentalists in colloids know well (but seldom tell you)...

*Absolute* determination of  $\varphi$  very difficult!

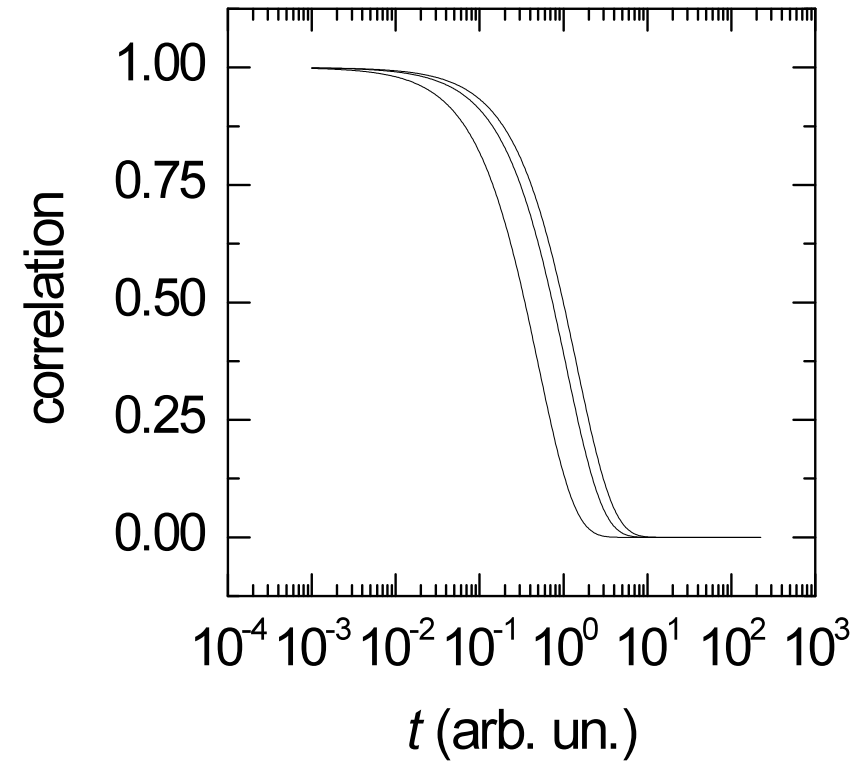




# Dynamical heterogeneity

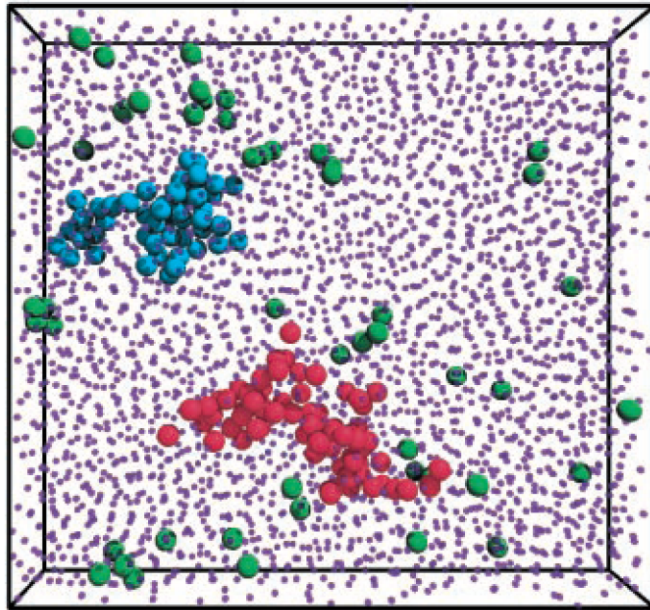


$\sigma$

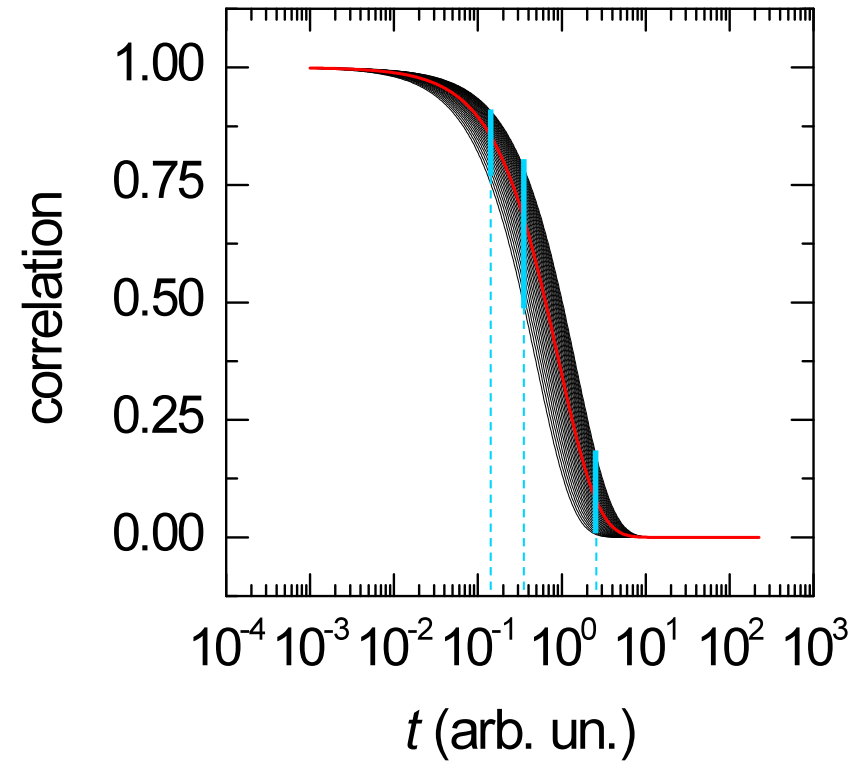


*Weeks et al. Science 2000*

# Dynamical heterogeneity

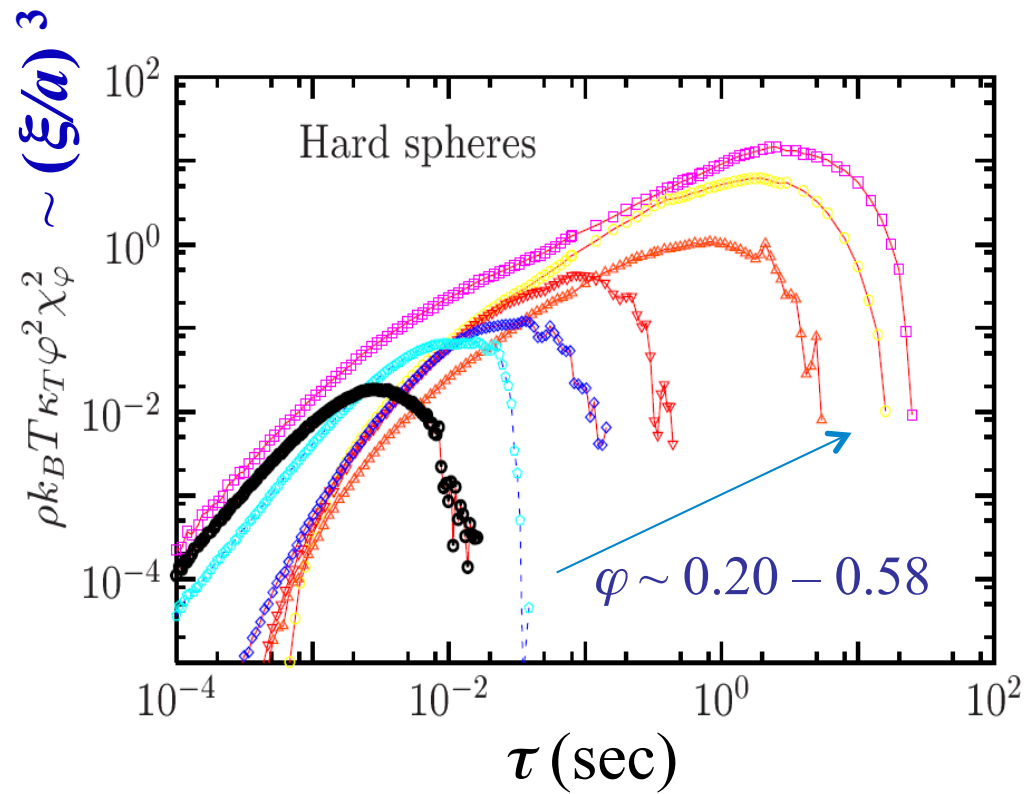


*Weeks et al. Science 2000*



$\text{var}[f(\tau)] \rightarrow \chi_4(\tau)$   
**dynamical susceptibility**

# DH in colloidal HS



dynamical susceptibility  $\chi_4(\tau)$

*Berthier et al., Science 2005*

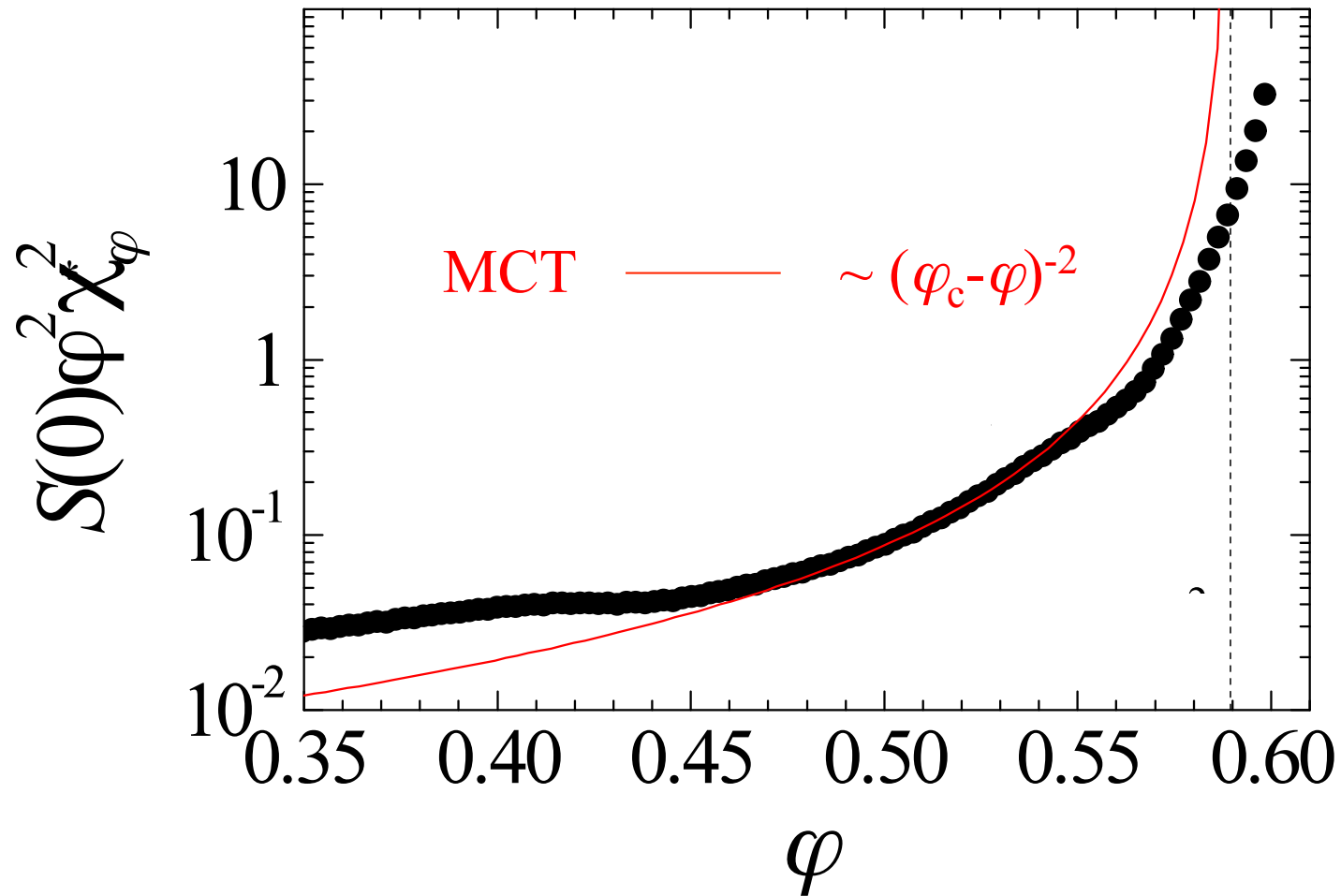
A useful relationship

$$\chi_4(\tau) \geq S(0)\varphi^2 \chi_\varphi^2(\tau)$$

- $\chi_\varphi(\tau) = \frac{\partial f(\tau)}{\partial \varphi}$
- equilibrium
- OK at high  $\varphi$

*Berthier et al., J. Phys. Chem. 2007*

# $\varphi$ dependence of the max of $\chi_4$

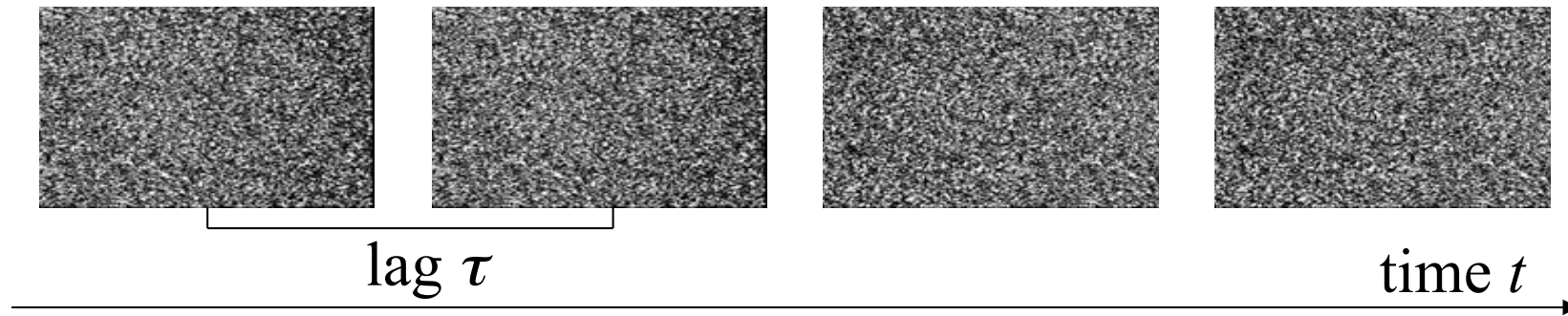


High  $\varphi$  :  
**deviation**  
**from MCT**

# Outline

- Average dynamics and dynamical heterogeneity of supercooled colloidal HS
- **Temporal fluctuations of the dynamics in jammed/glassy materials**
- Real-space measurements of correlations: ultra-long range spatial correlations and elasticity

# Time Resolved Correlation (TRC)

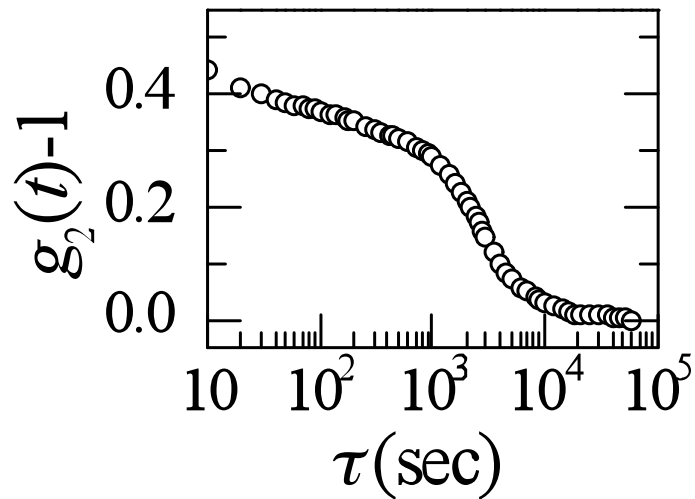


degree of correlation  $c_I(t, \tau) = \frac{\langle I_p(t) I_p(t + \tau) \rangle_p}{\langle I_p(t) \rangle_p \langle I_p(t + \tau) \rangle_p} - 1$

**degree of correlation**  $c_I(t, \tau) = \frac{\langle I_p(t) I_p(t + \tau) \rangle_p}{\langle I_p(t) \rangle_p \langle I_p(t + \tau) \rangle_p} - 1$

Average over  $t$  ↓

intensity correlation  
function  $g_2(\tau) - 1$



$g_2(\tau) - 1 \sim f(\tau)^2 \longrightarrow$  Average dynamics



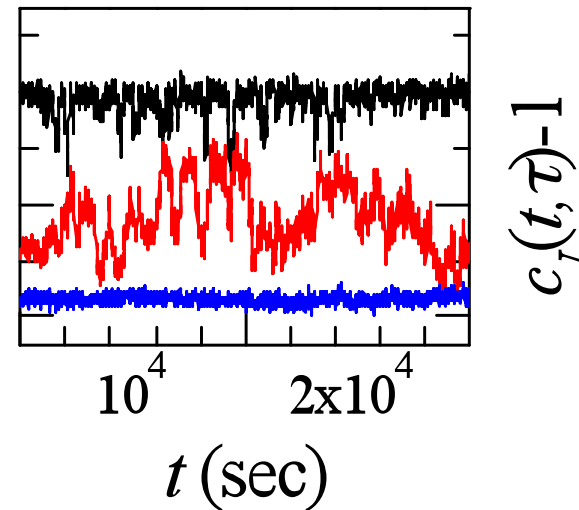
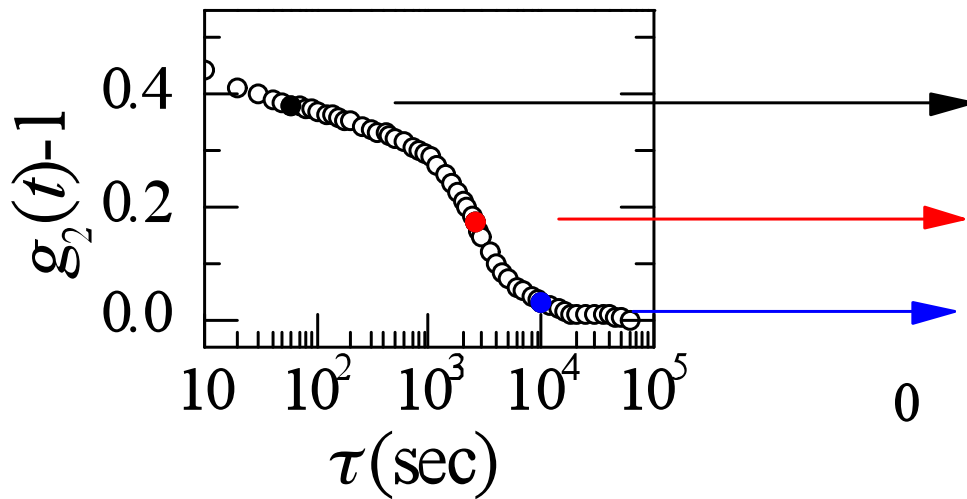
**degree of correlation**  $c_I(t, \tau) = \frac{\langle I_p(t) I_p(t + \tau) \rangle_p}{\langle I_p(t) \rangle_p \langle I_p(t + \tau) \rangle_p} - 1$

Average over  $t$  ↓

intensity correlation function  $g_2(\tau) - 1$

fixed  $\tau$ , vs.  $t$  ↓

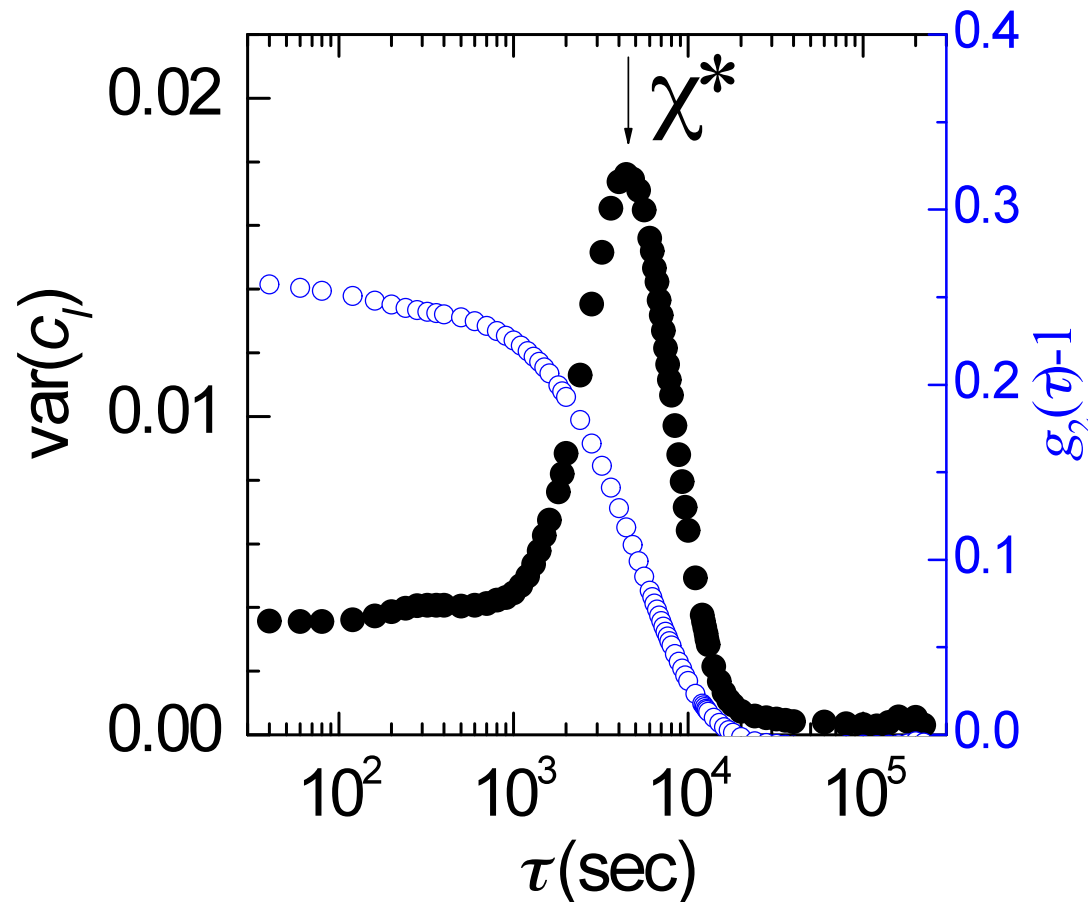
**fluctuations** of the dynamics



$g_2(\tau) - 1 \sim f(\tau)^2$  → Average dynamics

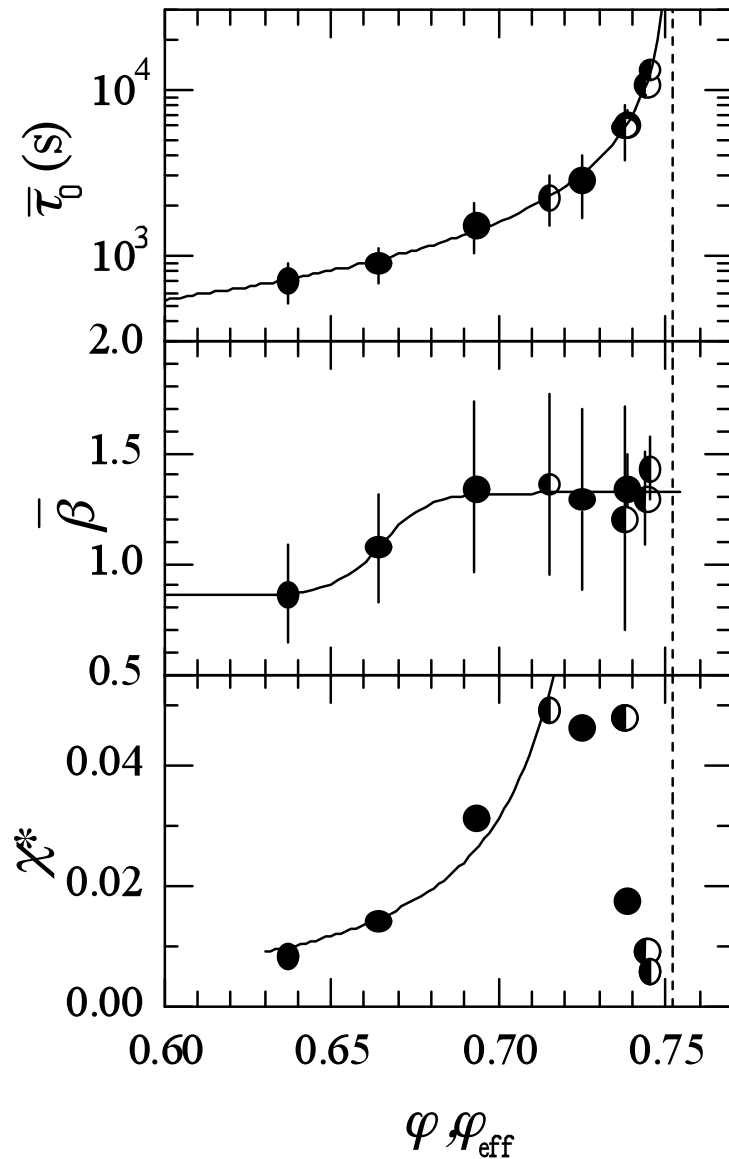
$\text{var}[c_I(\tau)] \sim \chi(\tau)$

# (Polydisperse) colloids near rcp



- PVC in DOP
- $a \sim 5 \mu\text{m}$
- polydispersity  $\sim 33\%$
- slightly soft
- $\varphi$  close to rcp
- multiple scattering (DWS)  
 $\Lambda \sim 10 \text{ nm} \ll a$

# Non-monotonic behavior of $\chi^*(\varphi)$



Fit:  $g_2(\tau)-1 \sim \exp[-(\tau/\tau_0)^\beta]$

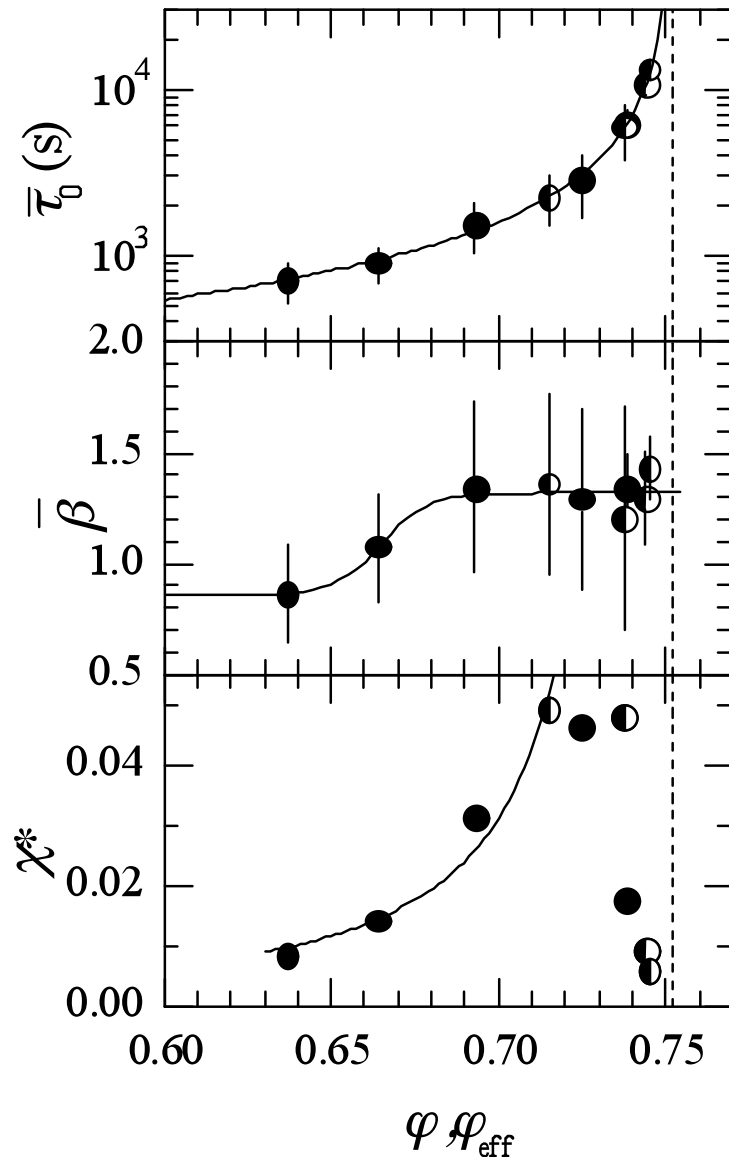
← relaxation time

← stretching exponent

← peak of dynamical susceptibility

**non monotonic!**

# Non-monotonic behavior of $\chi^*(\varphi)$



## Competition between:

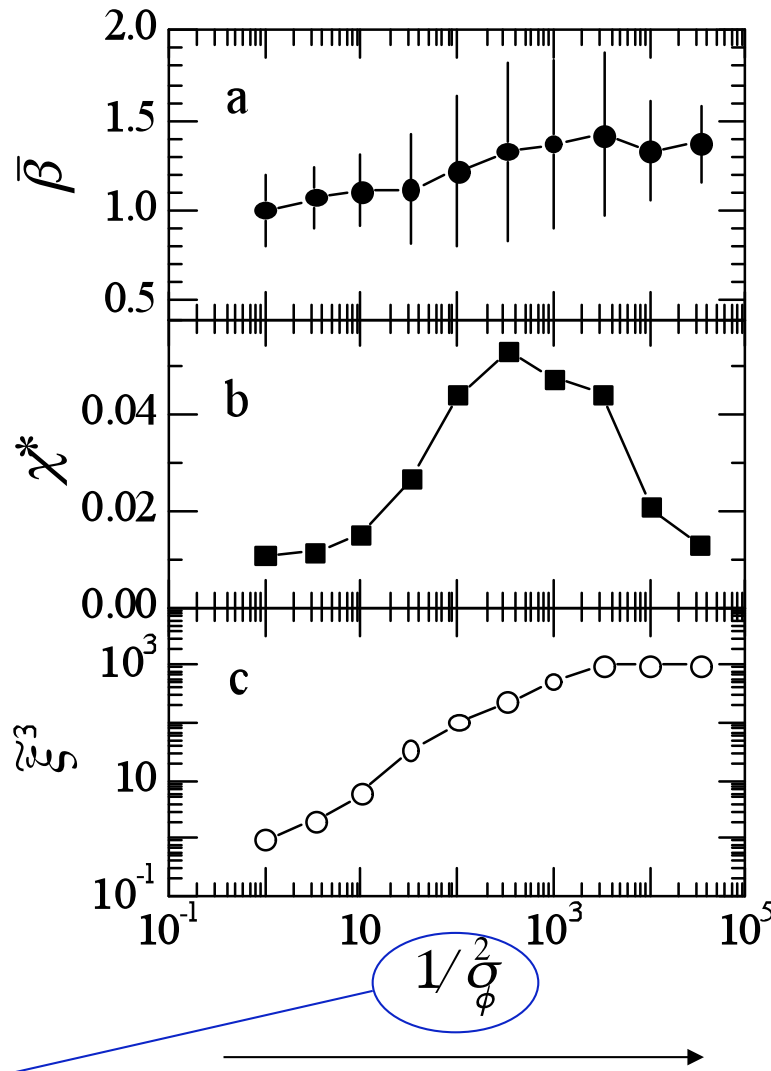
increasing  $\xi$  ( $\chi^* \nearrow$  as  $\varphi \nearrow$ )

increasingly restrained displacement  
(many events over  $\tau_0$ ,  $\chi^* \searrow$  as  $\varphi \nearrow$ )

## Model:

- random events of size  $\xi$
- Poissonian statistics
- Quasi-ballistic motion  
( $\langle \delta r^2 \rangle \sim n^p$ ,  $p = 1.65$ )

# Simulating the model



← cell thickness!

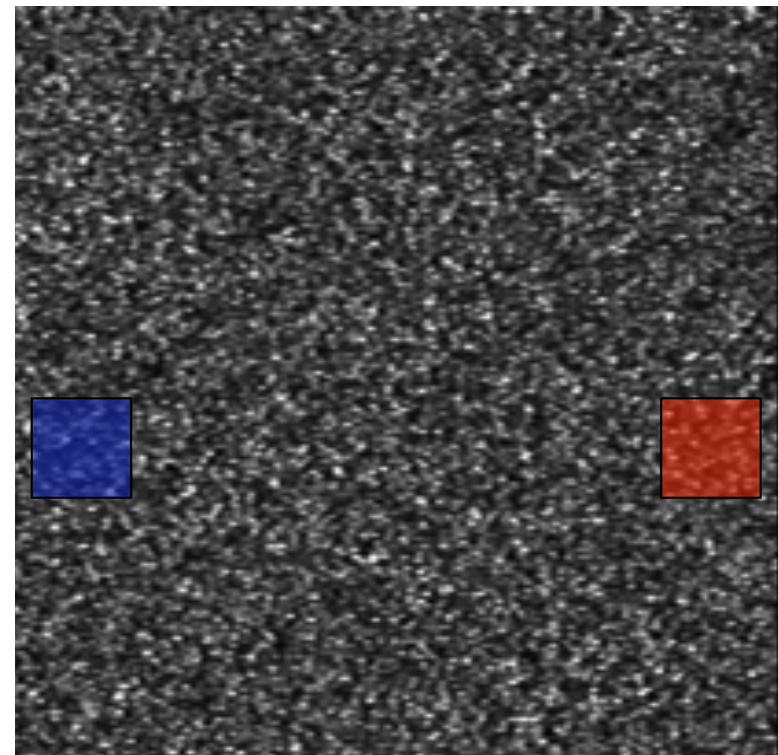
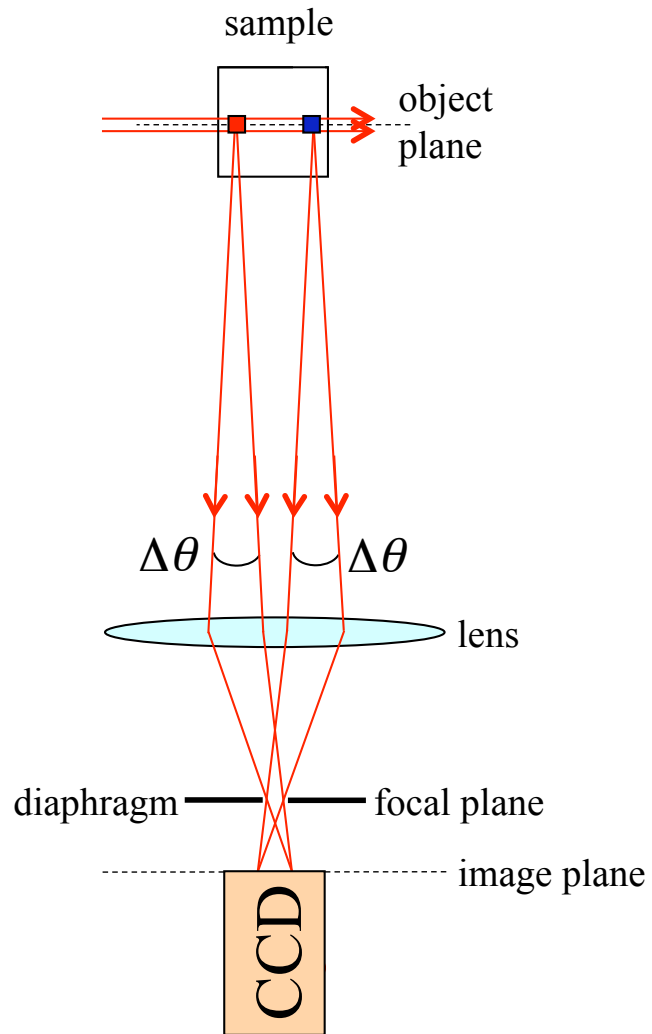
inverse jump size

volume fraction

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# Measuring $\xi$ by Photon Correlation Imaging



2.3 mm

*Duri et al., Phys. Rev. Lett. 2009*

$$\theta = 90^\circ \longrightarrow 1/q \sim 45 \text{ nm}$$

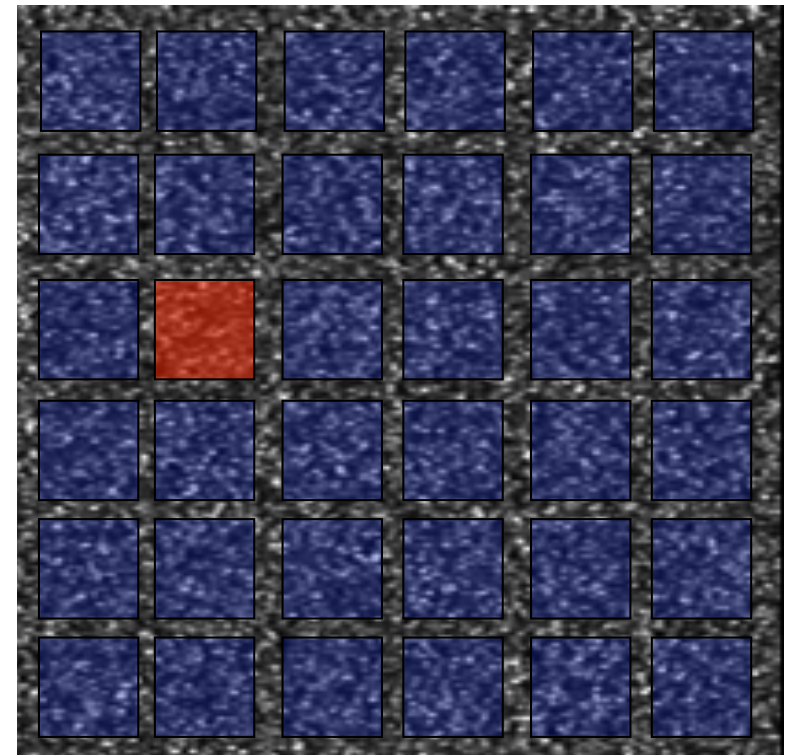


# Local, instantaneous dynamics: $c_I(t, \tau, \mathbf{r})$

$$c_I(t, \tau, \mathbf{r}) = \frac{\langle I_p(t) I_p(t + \tau) \rangle_{p(\mathbf{r})}}{\langle I_p(t) \rangle_p \langle I_p(t + \tau) \rangle_{p(\mathbf{r})}} - 1$$

Note:  $\langle \langle c_I(t, \tau, \mathbf{r}) \rangle_t \rangle_r = g_2(\tau) - 1$

$[g_2(\tau) - 1] \sim f(\tau)^2$



2.3 mm

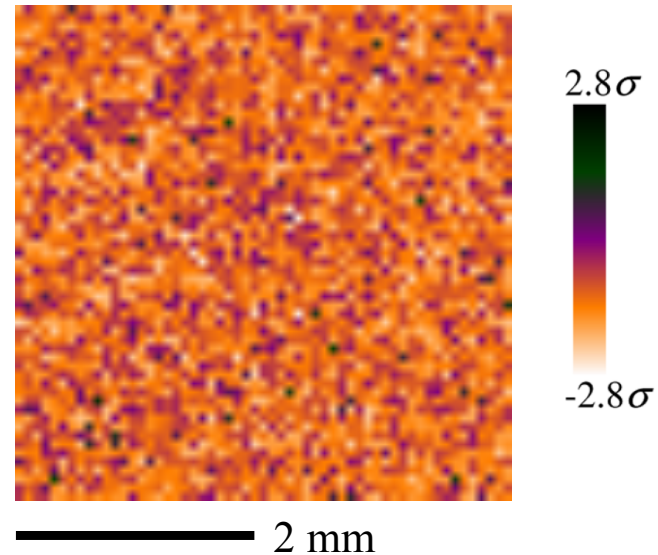
# Dynamic Activity Maps

## Brownian particles

$$g_2(\tau)-1 \sim \exp[-\tau/\tau_r], \tau_r = 40 \text{ s}$$

$$c_I(t_0, \tau_r/200, \mathbf{r})$$

Movie accelerated 10x

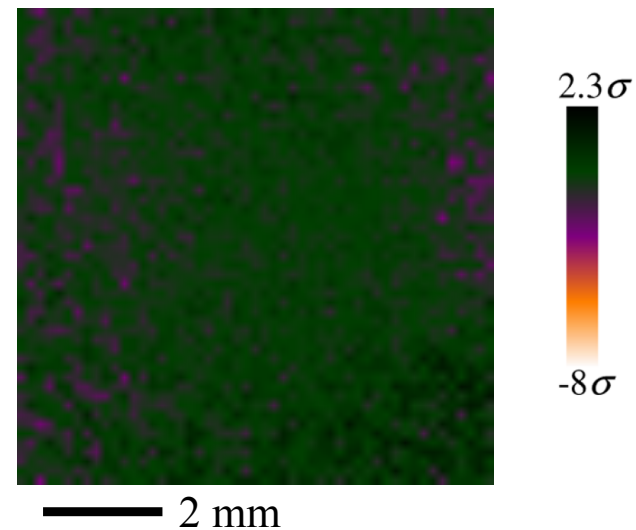


## Colloidal gel

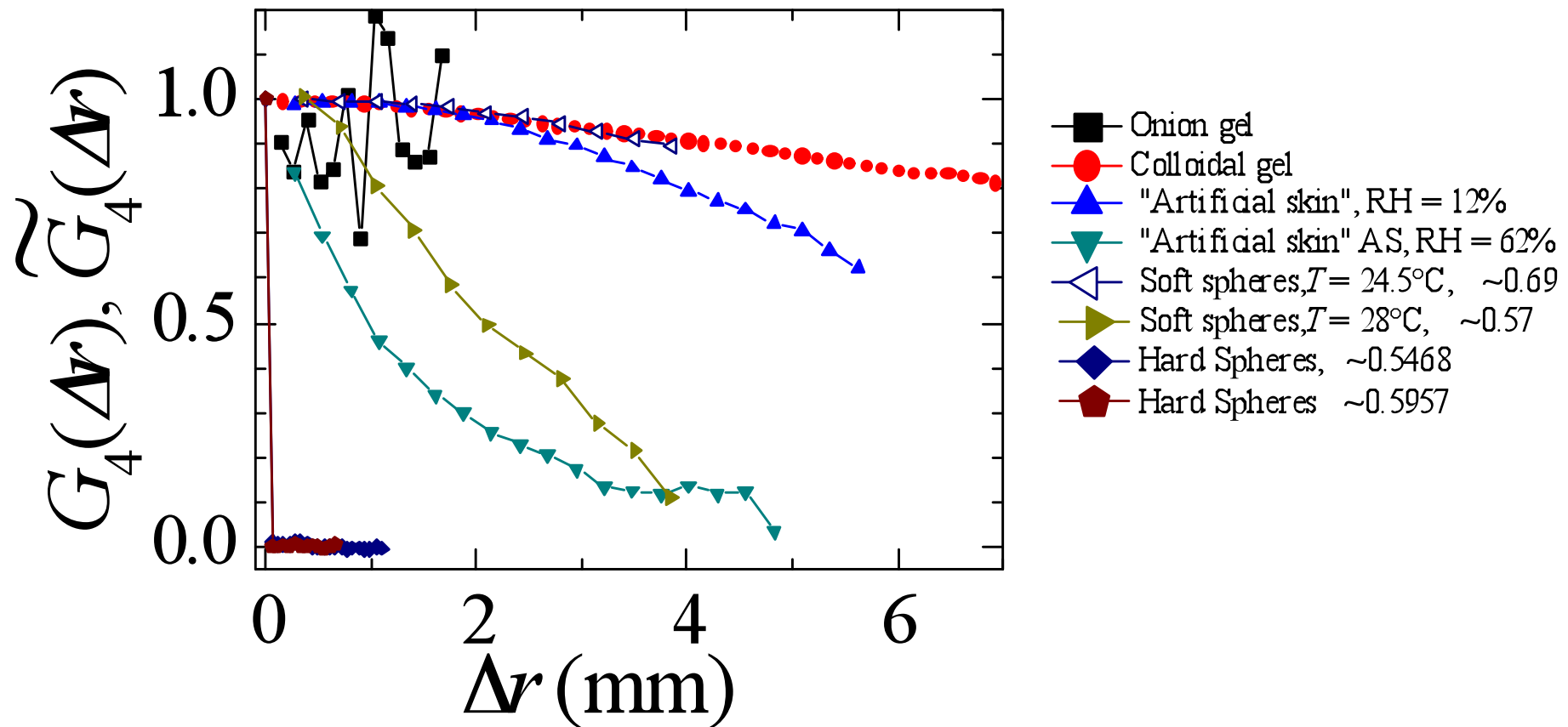
$$g_2(\tau)-1 \sim \exp[-(\tau/\tau_r)^{1.5}], \tau_r = 5000 \text{ s}$$

$$c_I(t_0, \tau_r/10, \mathbf{r})$$

Movie accelerated 500x



# Spatial correlation of the dynamics: $\xi \sim$ system size in jammed soft matter!

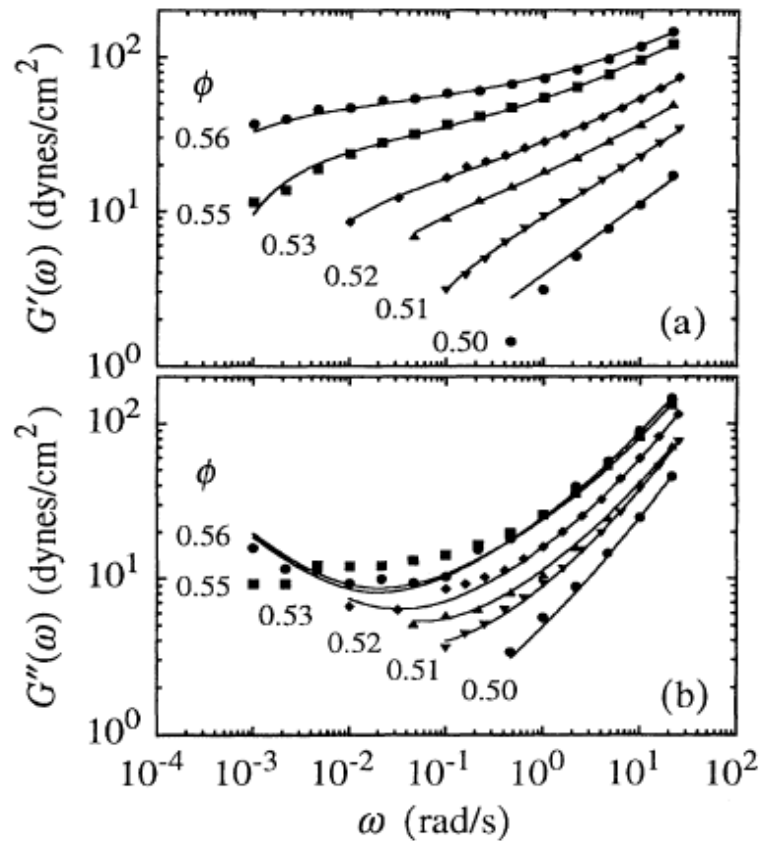


# When does $\xi \longrightarrow$ infinity?

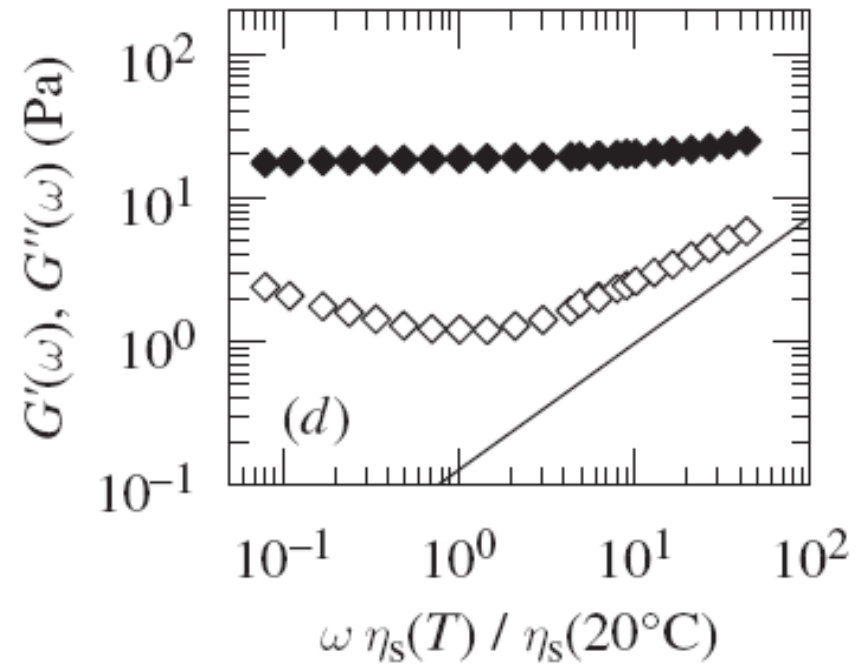
- $G' > G''$  : **strain propagation** in a predominantly **solid-like** material
- **Magnitude** of  $G'$  **unimportant** ( $G'_{\text{onions}}/G'_{\text{colloidal gel}} \sim 6 \times 10^4$  !!)

# Rheology of hard and soft spheres

## Hard spheres



## Soft spheres, $\phi_{\text{nom}} = 0.69$



*Sessoms et al., Phil. Mag. A 2009*

*Mason & Weitz., Phys. Rev. Lett. 1995*

# When does $\xi \longrightarrow$ infinity?

- $G' > G''$  : **strain propagation** in a predominantly **solid-like** material
- **Magnitude** of  $G'$  **unimportant** ( $G'_{\text{onions}}/G'_{\text{colloidal gel}} \sim 6 \times 10^4$  !!)
- **Origin of elasticity:**
  - *entropic* (e.g. hard spheres)  $\xi$  very **small**
  - *enthalpic* (attractive gels, squeezed particles...)  $\xi$  **system-size**

# Conclusions

- DH a **general feature** of glassy/jammed dynamics
- Supercooled hard spheres:
  - **equilibrium dynamics** above the (apparent) MCT divergence
  - $\xi$  **limited** to  $5-10a$
- Jammed materials:
  - $\xi$  and  $\chi$  **may decouple!**
  - $\xi \sim$  **system size**
  - role of the **origin of elasticity**



# Some open questions...

- Structural signature of Dynamical Heterogeneity?
- Dynamical Heterogeneity and aging
- Origin of the dynamics and its long range correlation in (undriven) jammed soft materials. Internal stress ?
- DH in the spontaneous dynamics and plasticity/shear banding/fracture in driven systems

# Thanks to...

L. Berthier (LCVN), V. Trappe (Fribourg)

## *Hard spheres*

G. Brambilla, M. Pierno, D. El Masri (LCVN), A. Schofield (Edinburgh),  
G. Petekidis (FORTH)

## *TRC*

A. Duri, P. Ballesta (LCVN), D. Sessoms, H. Bissig (Fribourg)

## *PCI*

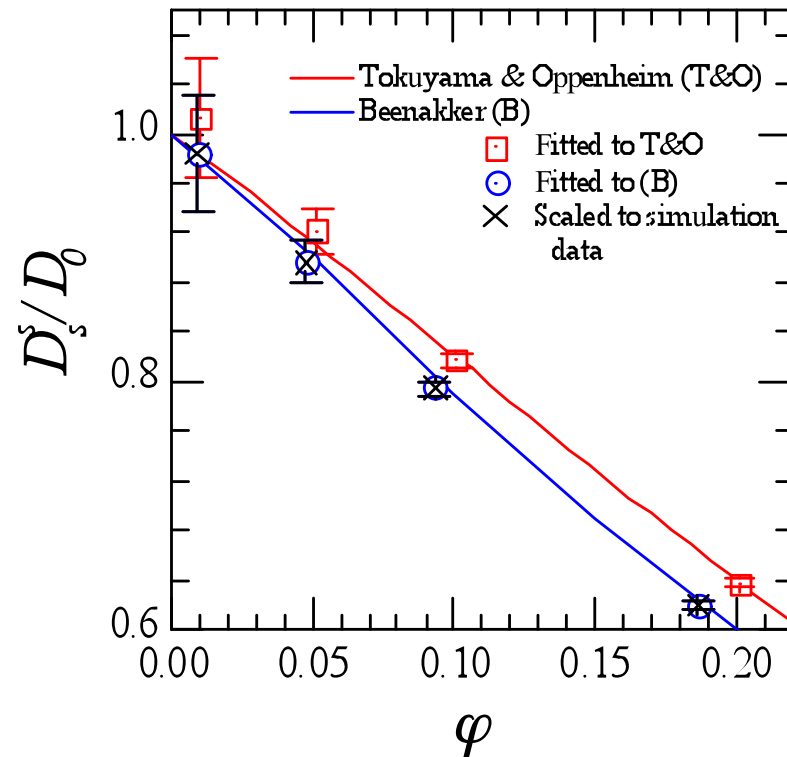
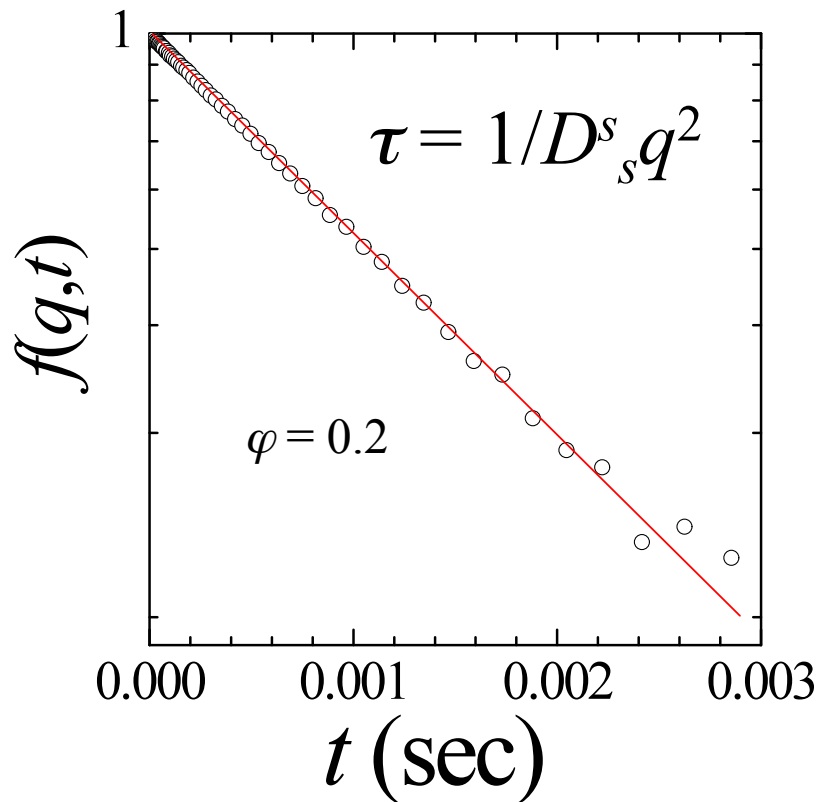
A. Duri, S. Maccarrone (LCVN), D. Sessoms (Fribourg), E. Pashkowski  
(Unilever)

## *Funding*

CNES, ACI, ANR, PICS, Unilever

# Determining the volume fraction

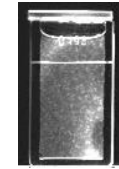
$\varphi$  dependence of the short-time diffusing coefficient



***Absolute* uncertainty:  $\sim 4\%$**

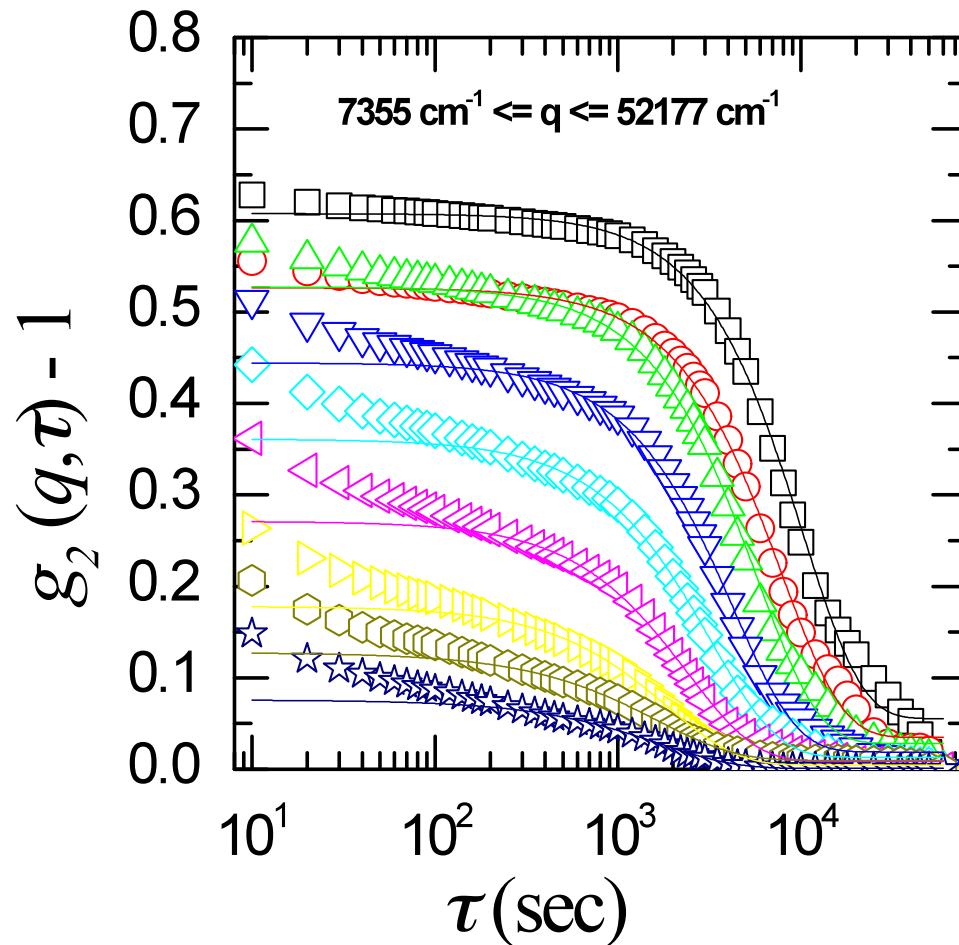
***Relative* uncertainty:  $\sim 10^{-4}$**

# PS gel: time-averaged dynamics



$$g_2(q, \tau) - 1 \sim [f(q, \tau)]^2$$

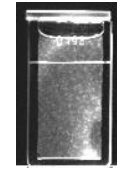
$$f(q, \tau) = \sum_{j,k} \langle \exp[iq \cdot (\mathbf{r}_j(\tau) - \mathbf{r}_k(0))] \rangle$$



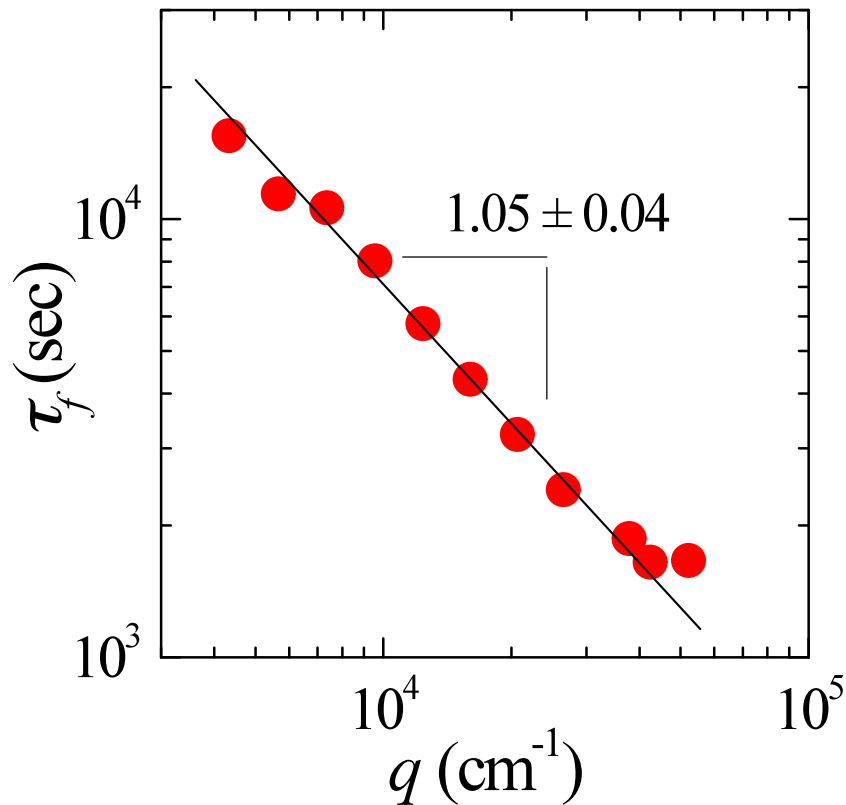
- **Fast dynamics:** overdamped vibrations (~ 500 nm) *Krall & Weitz PRL 1998*
- **Slow dynamics:** rearrangements

$$g_2(q, \tau) - 1 \sim \exp[-(\tau/\tau_c)^2]$$

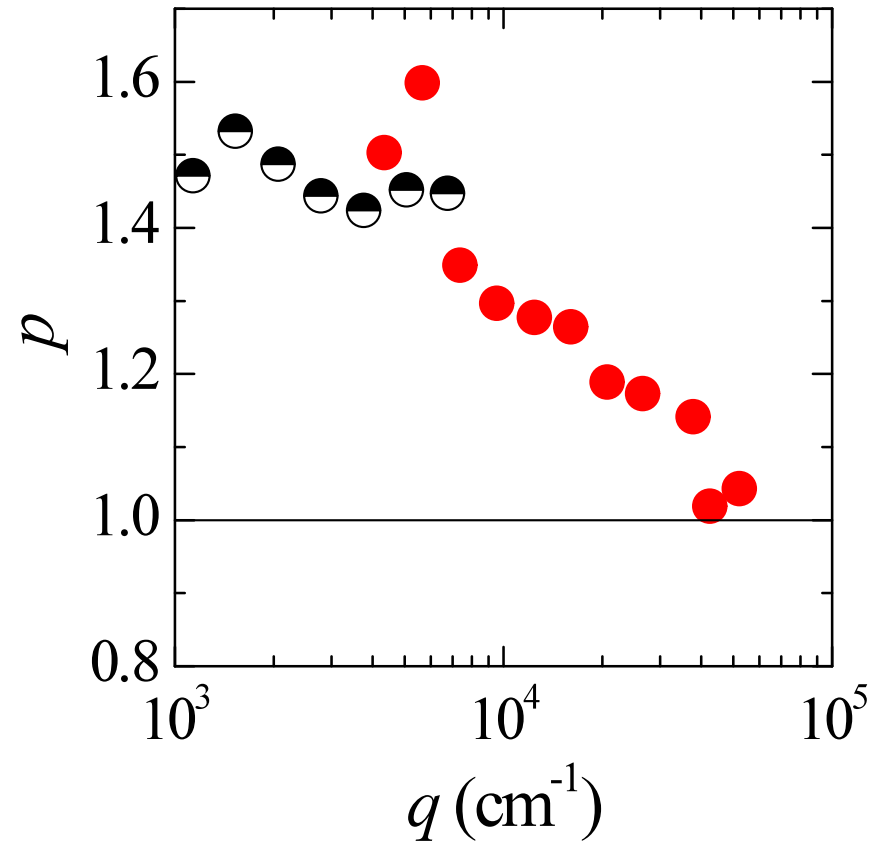
# PS gels: $q$ dependence of $\tau_f$ and $p$



$$g_2(q, \tau) - 1 \sim \exp\left[-(\tau/\tau_f)^p\right]$$

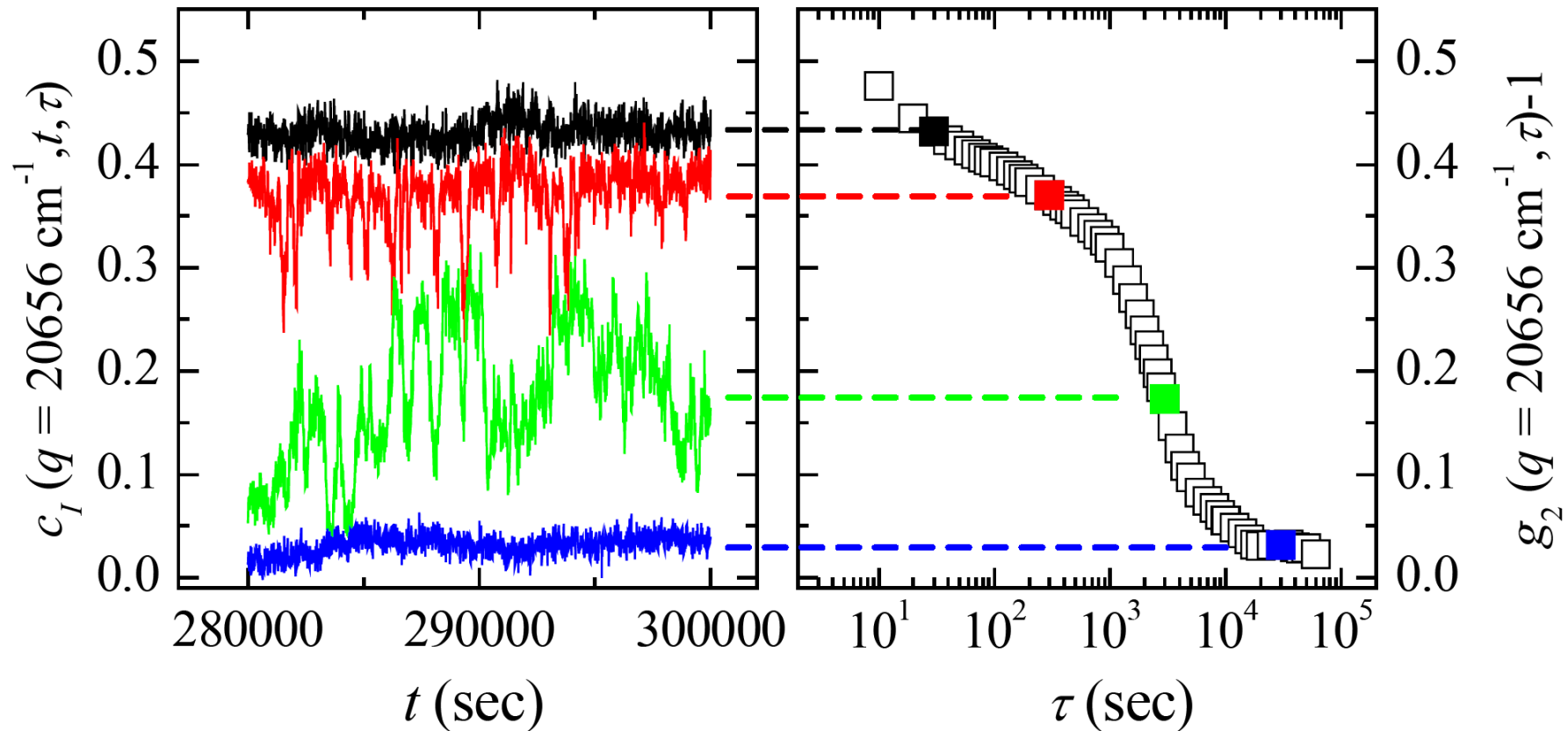
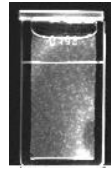


« ballistic » motion



« compressed » exponential

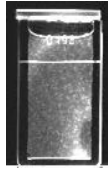
# PS gel: temporally heterogeneous dynamics



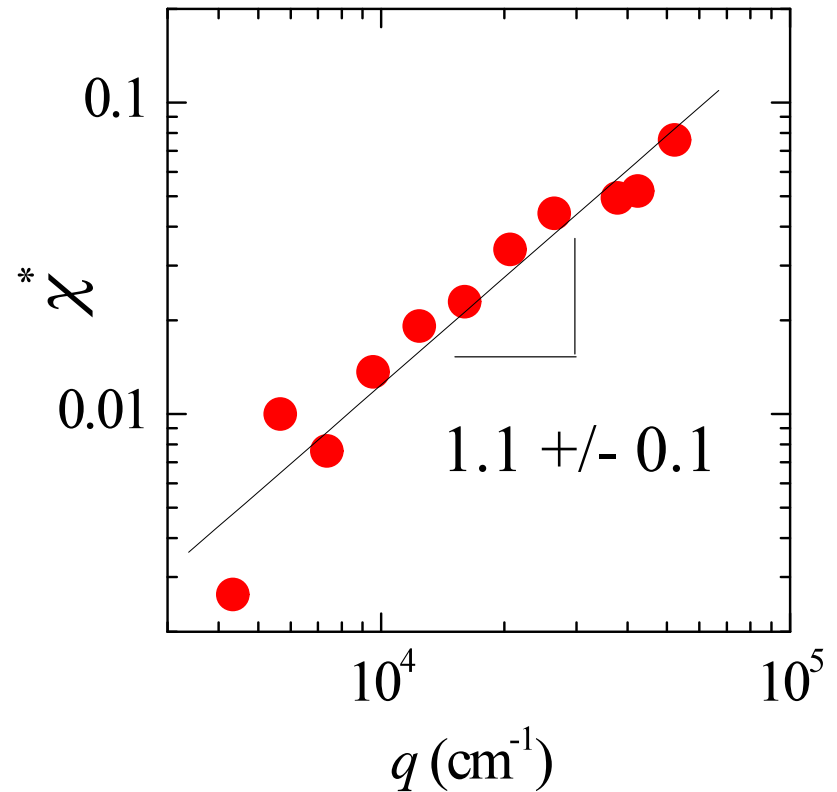
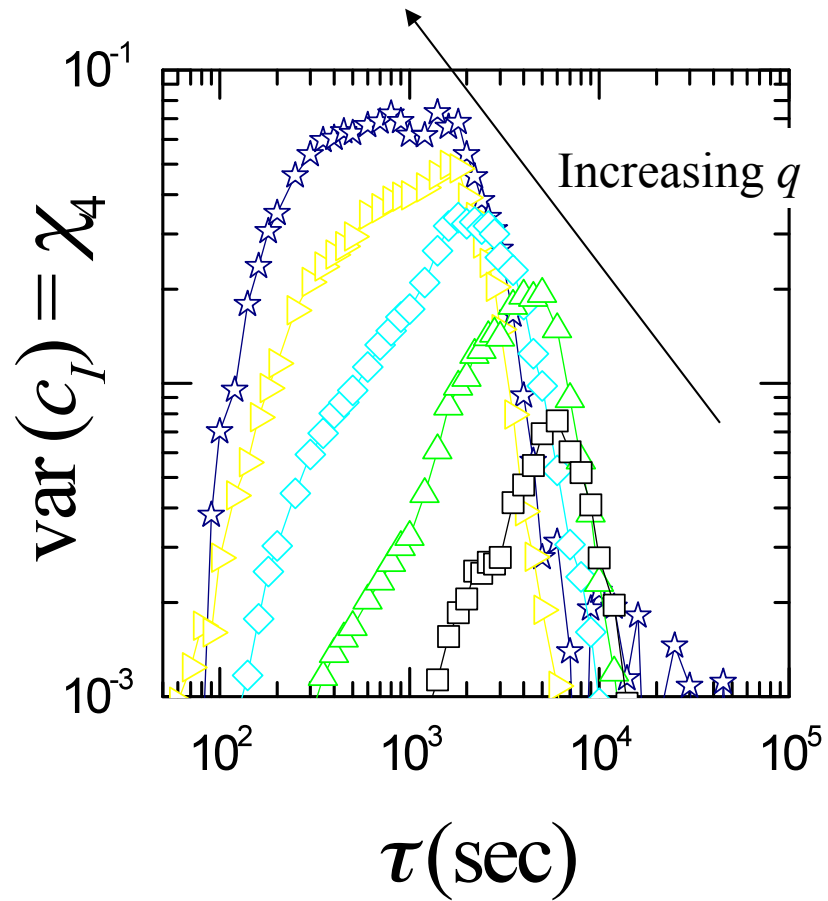
$c_I(t_w, \tau)$  @ fixed  $\tau$ :  
temporal **fluctuations**

$\langle c_I(t_w, \tau) \rangle_{t_w}$  : **average** dynamics

# Length scale dependence of $\chi = \text{var}(c_I)$

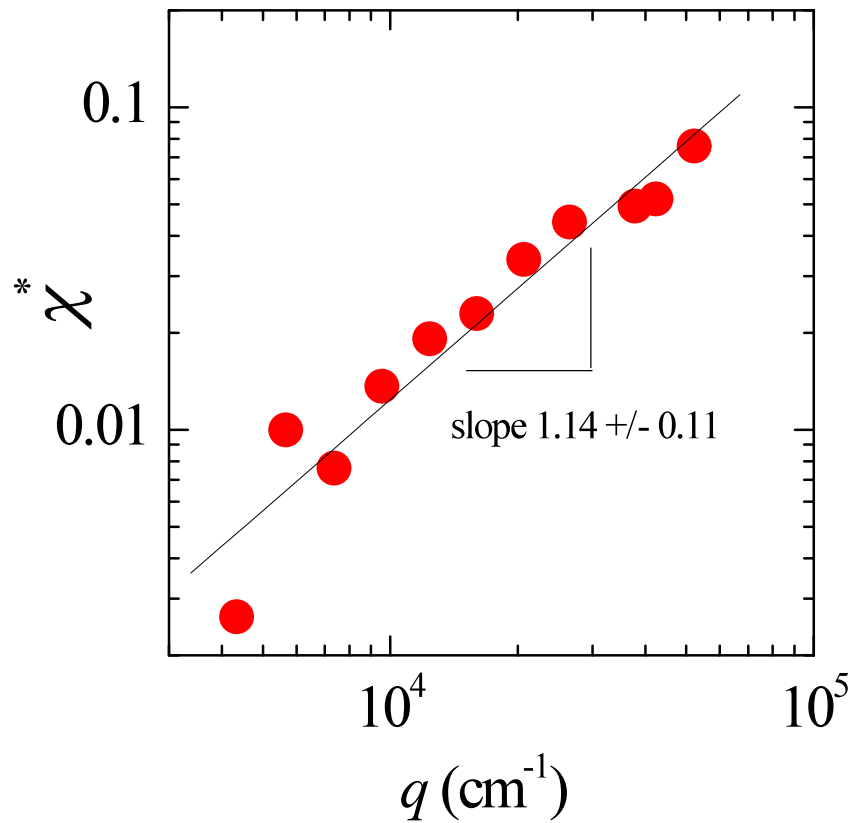


PS gel



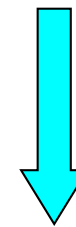


# Scaling of $\chi^*$



$$\chi^* \sim \text{var}(n)/\langle n \rangle \sim 1/\langle n \rangle$$

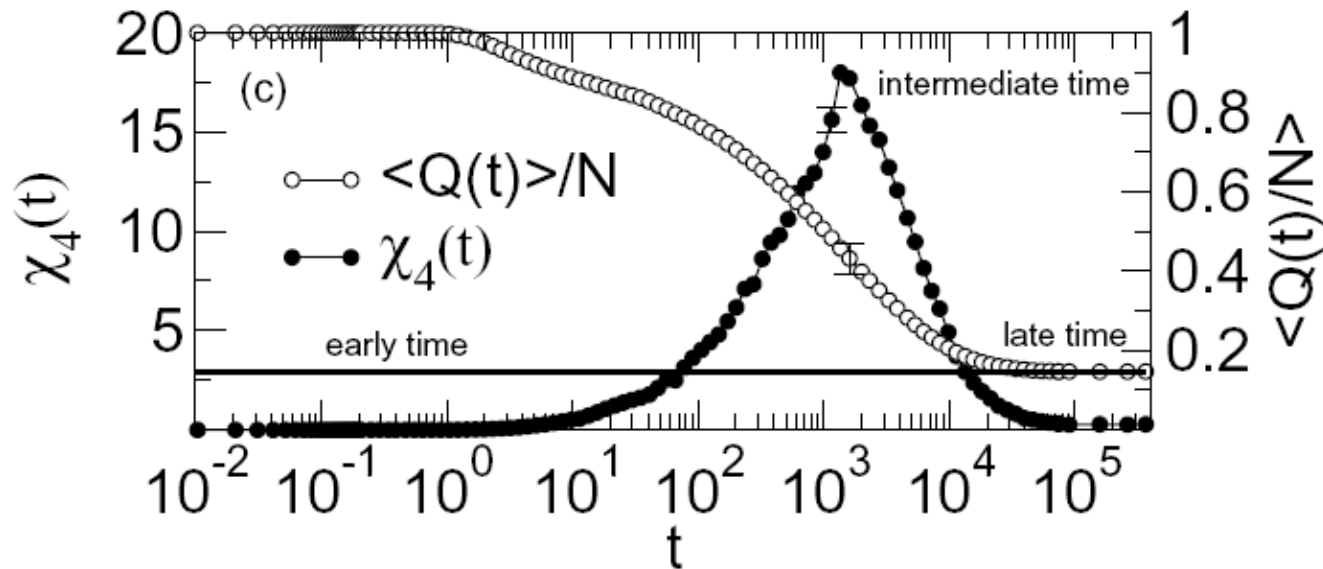
$$\langle n \rangle \sim \tau_f \sim 1/q$$



$$\chi^* \sim q$$

# Dynamical susceptibility in glassy systems

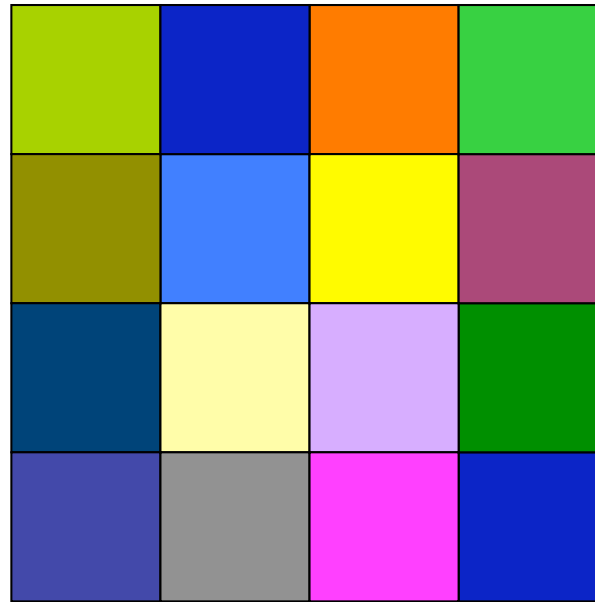
## Supercooled liquid (Lennard-Jones)



*Lacevic et al., Phys. Rev. E 2002*

$$\chi_4 = N \text{var}[Q(t)]$$

# Dynamical susceptibility in glassy systems



$N$  regions

$$\chi_4 = N \text{ var}[Q(t)]$$

$\chi_4$   $\longleftrightarrow$  dynamics spatially correlated

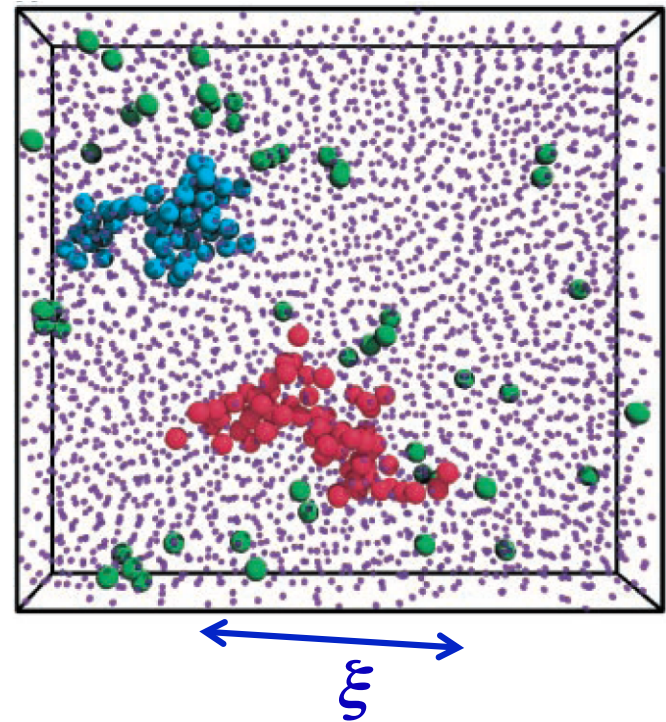
# Dynamical susceptibility in glassy systems

$$G_4(r; t) = \langle c(r; t, 0)c(0; t, 0) \rangle - \langle c(0; t, 0) \rangle^2$$

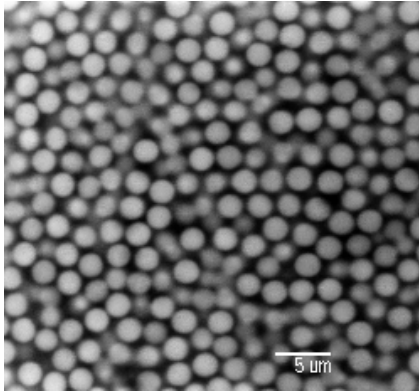
Spatial correlation of the dynamics

$$G_4(r; 0, t) \sim \frac{A(t)}{r^p} e^{-r/\xi_4(t)}$$

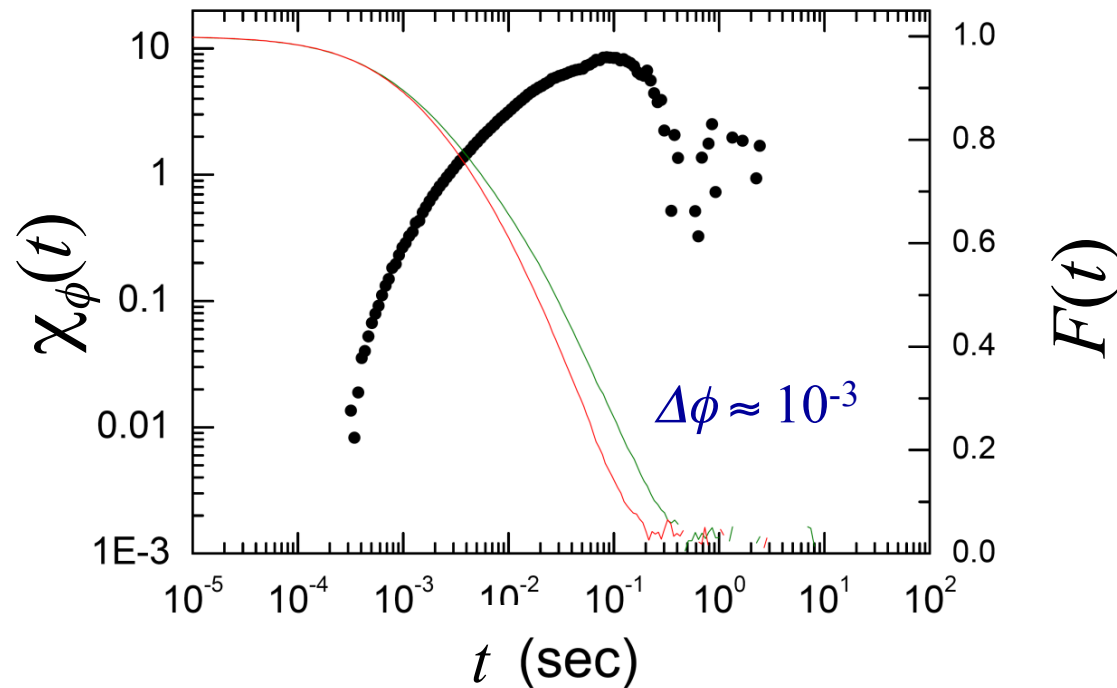
$$\chi_4(t) = \int dr G_4(r; t)$$



# The smart trick applied to colloidal HS



Define  $\chi_\phi(t) = \frac{\partial F(t)}{\partial \phi}$



# Dynamical heterogeneity: the theoreticians' trick

**Goal:** calculate 4-point dynamical susceptibility  $\chi_4 \sim$  size of rearranged region

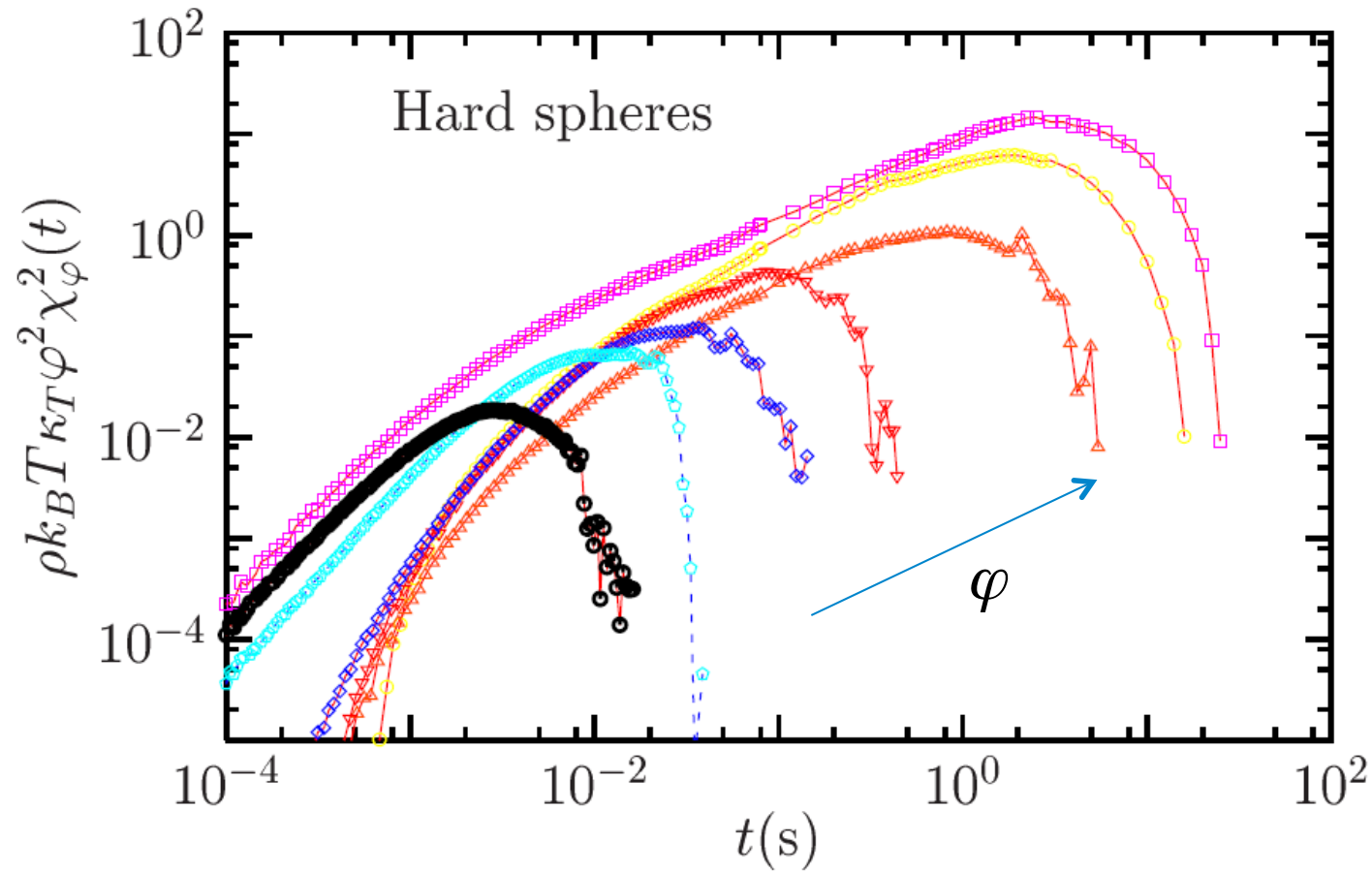
For colloidal HS at high  $\varphi$

$$\chi_T(t) = \frac{\partial f(t)}{\partial T}$$

$$\chi_\varphi(t) = \frac{\partial f(t)}{\partial \varphi}$$

$$\chi_4^{NPT}(t) = \chi_4^{NVE}(t) + \frac{k_B T^2}{c_V} \chi_T^2(t) + S(0) \varphi^2 \chi_\varphi^2$$

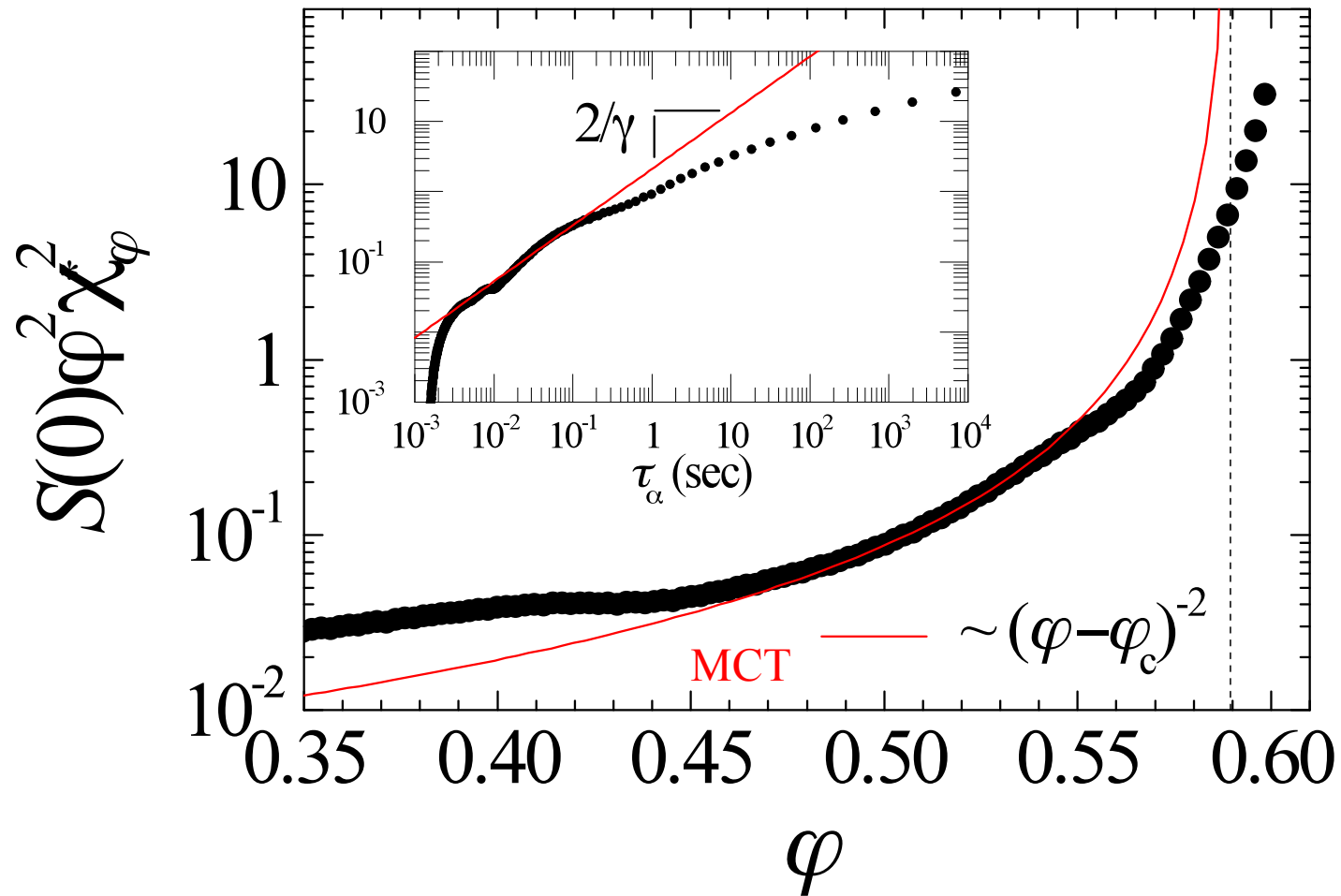
# Evidence of a growing dynamic length scale



$$\phi \sim 0.20 - 0.58$$

*Berthier et al., Science 2005*

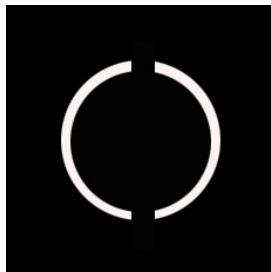
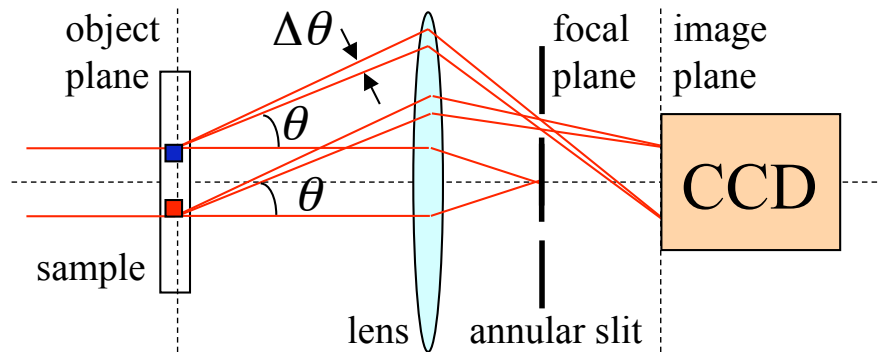
# $\varphi$ dependence of the max of $\chi_4$



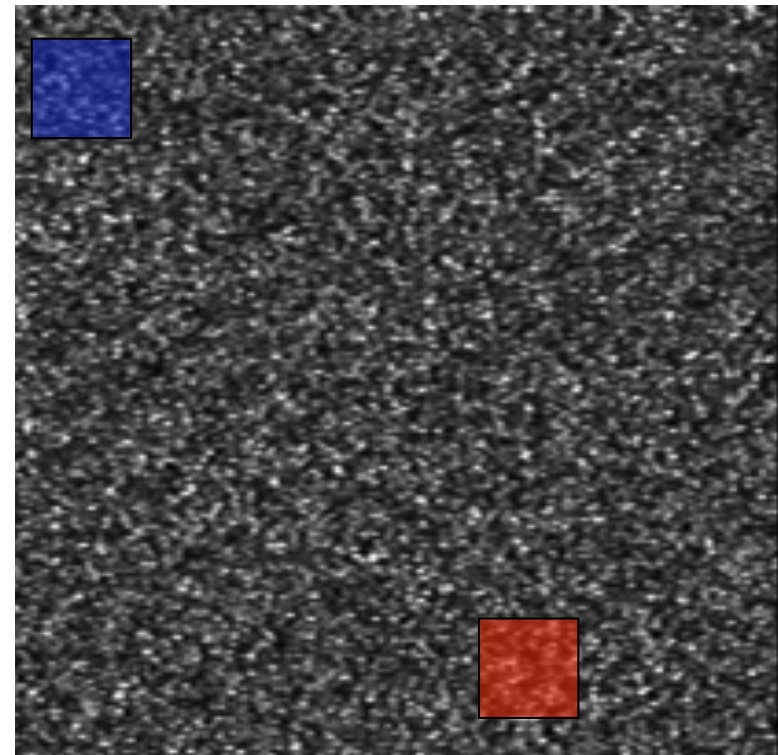
High  $\varphi$  :  
**deviation**  
**from MCT**



# Photon Correlation Imaging (PCIm)



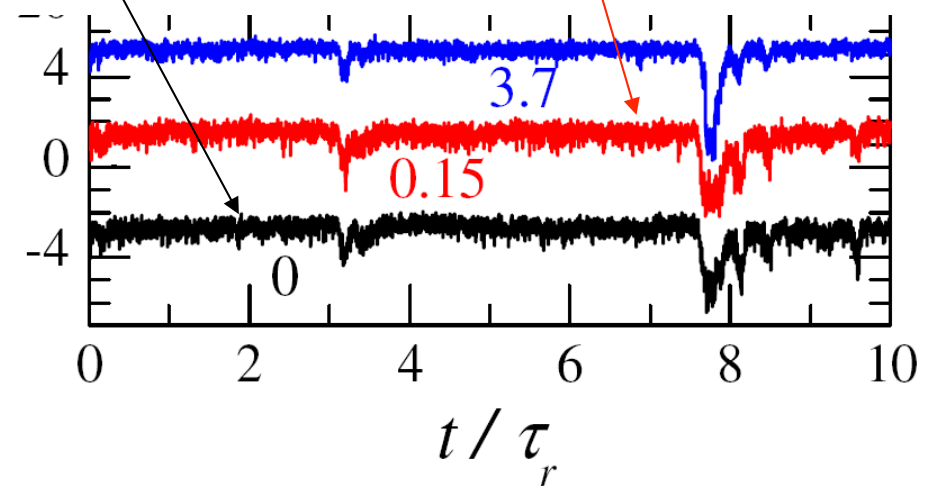
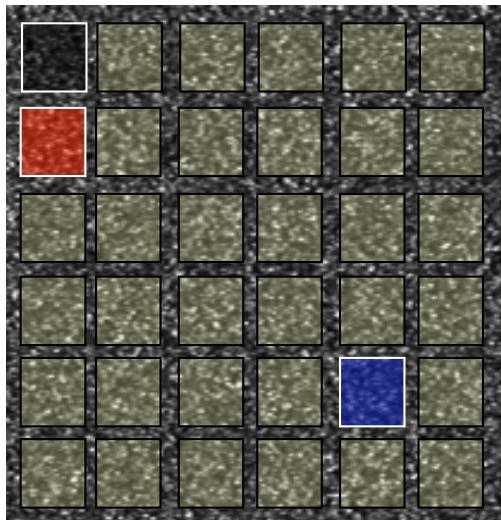
$$\theta = 6.4^\circ \longrightarrow q = 1 \mu\text{m}^{-1}.$$



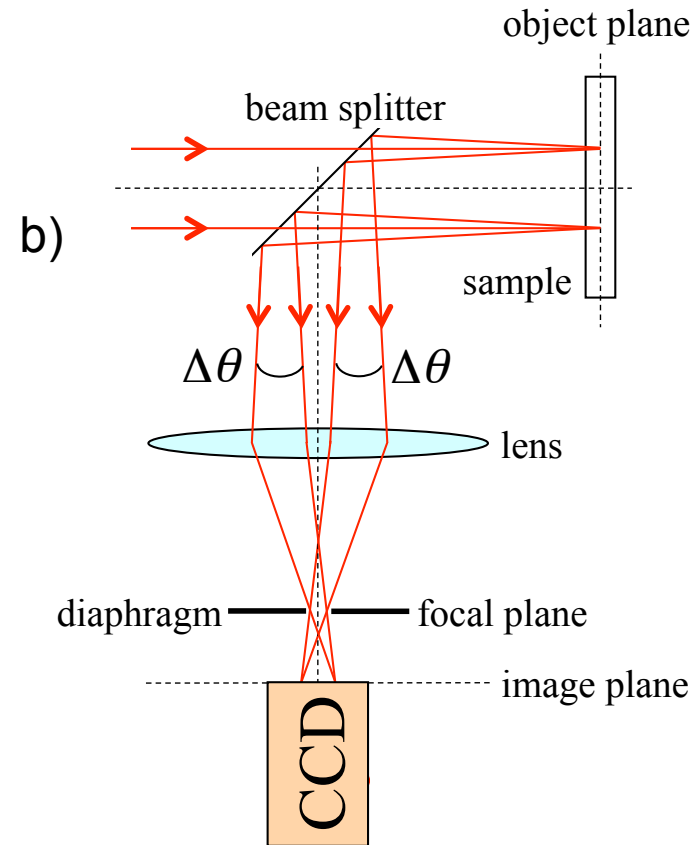
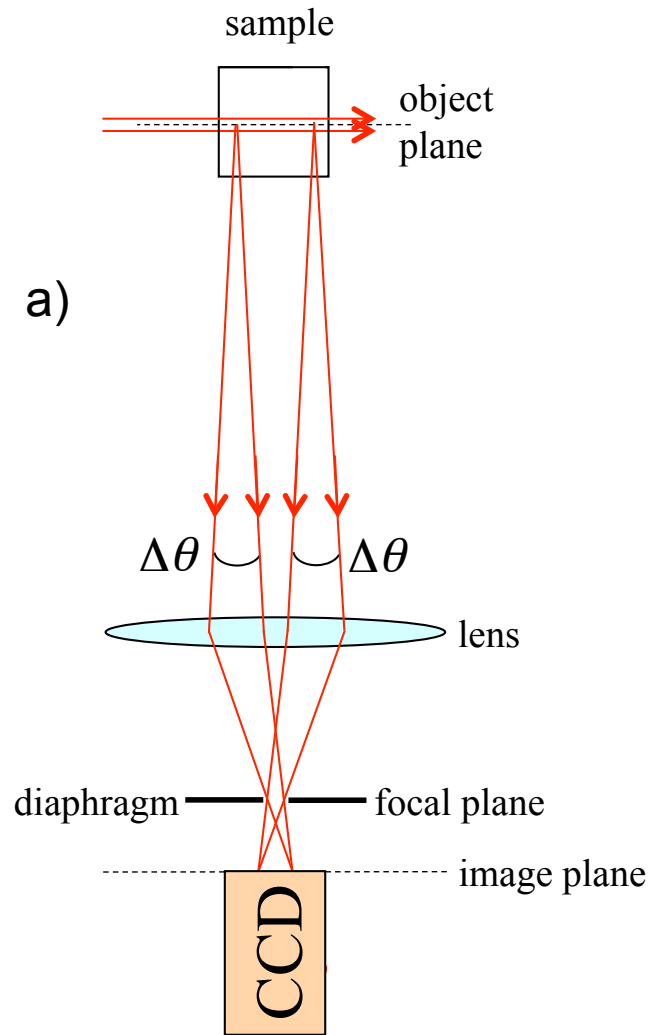
2.3 mm

# Spatial correlation of the dynamics

$$\tilde{G}_4(\Delta r, \tau) = \left\langle \frac{\langle \delta c_I(t, \tau; \mathbf{r}_1) \delta c_I(t, \tau; \mathbf{r}_2) \rangle_t}{\sigma(\tau, \mathbf{r}_1) \sigma(\tau, \mathbf{r}_2)} \right\rangle_{|\mathbf{r}_1 - \mathbf{r}_2| = \Delta r}$$



# Other PCIm geometries



# Additional stuff on diverging $\xi$

TABLE I: Rheological parameters of most of the system shown in Figs. 5 and 6. The sample names are as in the caption of Fig. 5.

	$\nu$ (Hz)	$G'(\nu)$ (Pa)	$G'(\nu)/G''(\nu)$	Ref.
On	1	600	15	[40]
CG	1	$\sim 0.9 \times 10^{-3}$	10	[52]
SoS, $\varphi = 0.57$	1.6	0.6	0.3	[27]
SoS, $\varphi = 0.69$	1.6	20	8	[27]
HS, $\varphi = 0.5468$	1	210	1.25	[53, 54]
HS, $\varphi = 0.5957$	1	$> 400$	$> 1.4$	[53, 54]
Laponite	0.7	$\gtrsim 300$	20	[48, 49]

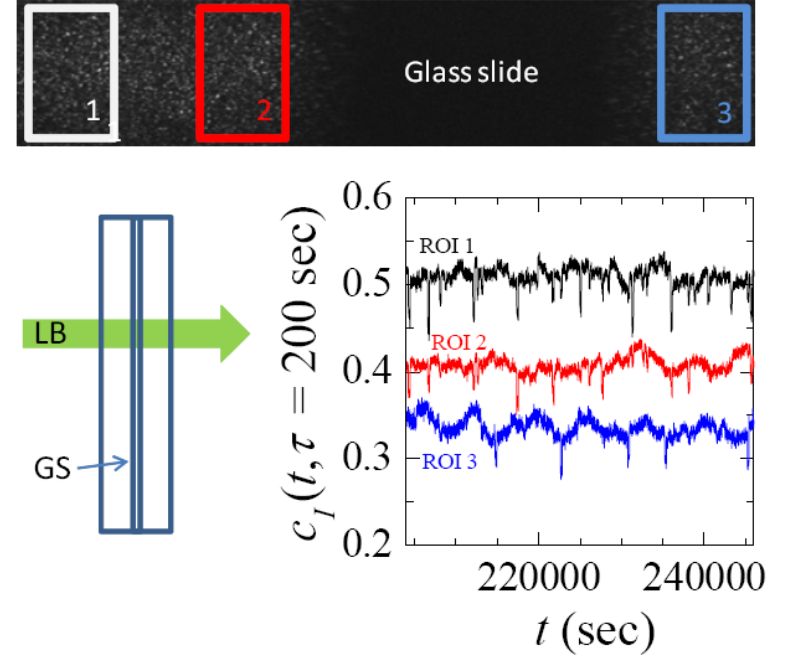


FIG. 6: Bottom left: schematic side view of the cell used for the experiment on Laponite. LB: laser beam, GS: glass slide. Top: typical CCD image of the scattering volume. The dark region corresponds to the thickness of the glass slide, view from the side. The three ROIs for which the degree of correlation is shown in the bottom right plot are highlighted. The size of the imaged region is  $2.62 \times 0.52 \text{mm}^2$ . Bottom right: time dependence of  $c_I$ , for a delay time  $\tau = 200$  sec, for the three ROIs shown above. For the sake of clarity, the curves of ROIs 1 and 2 have been offset vertically by 0.2 and 0.1, respectively. Note that the signals measured on the same side of the glass slide (ROIs 1 and 2) are correlated, while signals from opposite sides are uncorrelated.