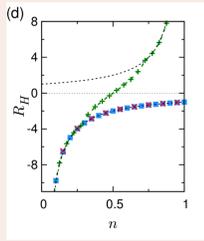


Hall effect in the two leg ladder geometry

$R_H = -1/n$ dotted line

Fermions: green Bosons: blue and purple



Greschner et al. Phys. Rev. Lett. **122**, 083402 (2019)

Tomonaga-Luttinger liquid in a single chain of interacting bosons

$$H = \sum_j -t(b_j^\dagger b_{j+1} + b_{j+1}^\dagger b_j) + \frac{U}{2} n_j (n_j - 1)$$

$$H_{TLL} = \int \frac{dx}{2\pi} \left[uK (\nabla\theta)^2 + \frac{u}{K} (\nabla\phi)^2 \right]$$

$$\frac{n_j}{a} = \rho_0 - \frac{\partial_x \phi}{\pi} + \sum_{m \neq 0} A_m e^{i2m(\phi(x) - \pi\rho_0 x)}$$

$$b_j = e^{i\theta(x)} \sum_m B_m e^{i2m(\phi(x) - \pi\rho_0 x)}$$

$$[\phi(x), \nabla\theta(y)] = i\pi\delta(x-y)(x=ja)$$

Non-perturbative definition of Tomonaga-Luttinger parameters

stiffness: $\mathcal{D} = uK$

compressibility: $\chi = K/(\pi\rho_0^2 u)$

F. D. M. Haldane Phys. Rev. Lett. **47**, 1840 (1981).

Tomonaga-Luttinger liquid : Band curvature corrections

$$H_{\text{curv.}} = \int dx \left[\alpha (\partial_x \phi)^3 + \gamma (\nabla\theta)^2 \partial_x \phi \right]$$

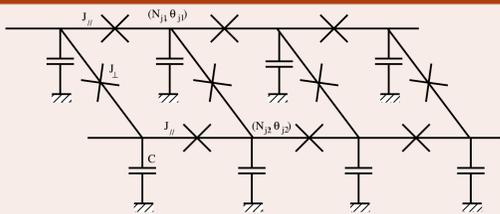
$$\alpha = -\frac{\partial}{\partial \rho_0} \left(\frac{u}{6\pi^2 K} \right)$$

$$\gamma = -\frac{\partial}{\partial \rho_0} \left(\frac{uK}{2\pi^2} \right)$$

Present unless there is particle-hole symmetry ($\phi \rightarrow -\phi, \theta \rightarrow -\theta$).

Matveev JETP 117, 508 (2013)

Josephson junction ladder



$$H = \sum_j \left\{ \left[\sum_{p=1,2} \frac{(e^* N_{j,p})^2}{2C} - J_\parallel \cos(\theta_{j+1,p} - \theta_{j,p}) \right] - J \cos(\theta_{j,1} - \theta_{j,2}) \right\}$$

In magnetic field: Peierls substitution

$$b_j^\dagger b_k \rightarrow e^{ie \int_j^k \vec{A} d\vec{l}} b_j^\dagger b_k$$

$$\cos(\theta_j - \theta_k) \rightarrow \cos \left[\theta_j - \theta_k - e \int_j^k \vec{A} d\vec{l} \right]$$

Landau gauge $A_x = -By$.

Bosonized Hamiltonian with band curvature

$$H = \int \frac{dx}{2\pi} \left[uK (\nabla\theta_s)^2 + \frac{u}{K} (\partial_x \phi_s)^2 \right]$$

$$+ \int \frac{dx}{2\pi} \left[uK (\nabla\theta_a - eBa/\sqrt{2})^2 + \frac{u}{K} (\partial_x \phi_a)^2 - \frac{2t_\perp}{a_0} \cos(\sqrt{2}\theta_a) \right]$$

$$+ \int \frac{dx}{\sqrt{2}} \left[\gamma \partial_x \phi_s (\nabla\theta_s)^2 + \alpha (\partial_x \phi_s)^3 \right]$$

$$+ \int \frac{dx}{\sqrt{2}} \left\{ \partial_x \phi_s \left[\gamma (\nabla\theta_a - eBa/\sqrt{2})^2 + 3\alpha (\partial_x \phi_a)^2 \right] \right.$$

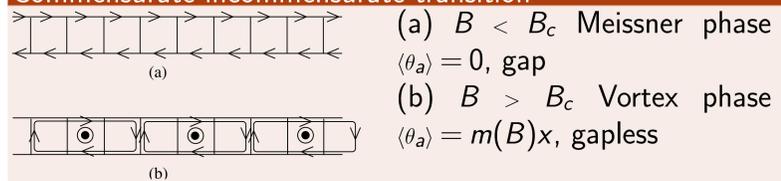
$$\left. + 2\gamma (\nabla\theta_a - eBa/\sqrt{2}) (\nabla\theta_s) \partial_x \phi_a \right\}$$

Hall imbalance from first order perturbation theory in α, γ .

$$\langle P_H \rangle = -\frac{\sqrt{2}}{\pi} \int_{-\infty}^{+\infty} \langle \partial_x \phi_a \rangle$$

$$= \frac{2\gamma L \langle \nabla\theta_s \rangle}{\pi} \int d\tau dx \langle T_\tau \left[(\nabla\theta_a - eBa/\sqrt{2}) \partial_x \phi_a \right] (x, \tau) \partial_x \phi_a(0, 0) \rangle + O(\gamma^2)$$

Commensurate-incommensurate transition



Hall imbalance in Meissner phase

$$\langle P_H \rangle^{(1)} = \frac{\pi Ba \gamma \langle j_s \rangle}{uK} \chi_{\rho_a \rho_s}(q=0, \omega_n=0)$$

Hall voltage obtained by cancelling $\langle P_H \rangle$

$$V_H = -\frac{\pi^2 \gamma \langle j_s \rangle}{euK} Ba$$

$$= -\langle j_s \rangle \frac{Ba}{2e} \frac{1}{uK} \frac{\partial}{\partial \rho_0} (uK)$$

$$= -\langle j_s \rangle \frac{Ba}{2e} \frac{\partial (\ln \mathcal{D})}{\partial \rho_0}$$

Hall coefficient near the Metal-insulator transition

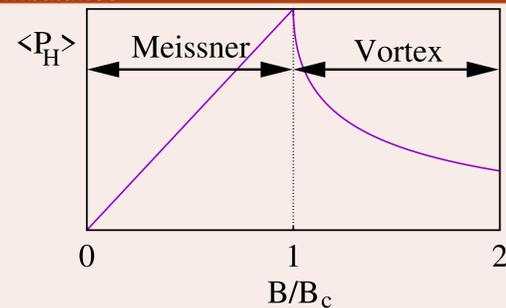
At $1/n$ filling ($n \in N^*$)

$$H_s = \int \frac{dx}{2\pi} \left[u_s K_s (\nabla\theta_\rho)^2 + \frac{u_s}{K_s} (\partial_x \phi_s)^2 \right]$$

$$- \frac{2gU}{(2\pi\alpha)^2} \int dx \cos(n\sqrt{2}\phi_s) + \frac{\sqrt{2}}{\pi} \mu \int dx \partial_x \phi_s$$

$$u_s^* \sim (\rho_0 - \rho_{0,c}) \quad \& \quad K_s^* \sim \frac{2}{n^2} \Rightarrow R_H \sim \frac{1}{e(\rho_0 - \rho_{0,c})}$$

Behavior of Hall imbalance



Perspectives

- Beyond perturbation theory
- Systems of many coupled chains
- Quantum Hall phases
- Vortex lattice phases