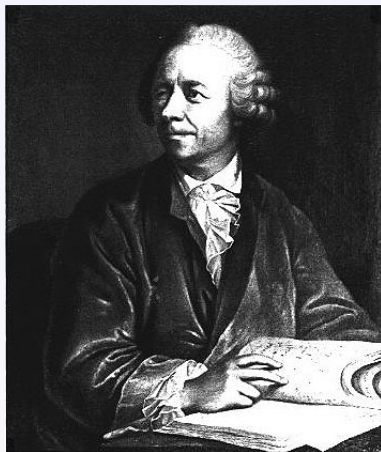


“Chirps” everywhere

Patrick Flandrin

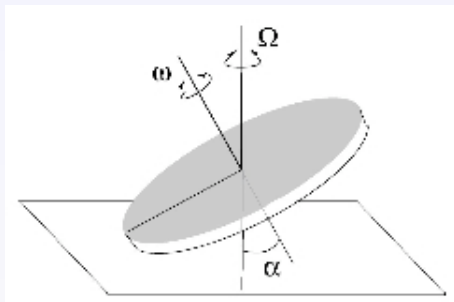
CNRS & École Normale Supérieure de Lyon, France





Leonhard Euler (1707-1783)

Euler's disk



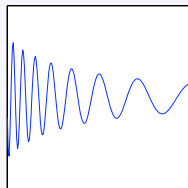
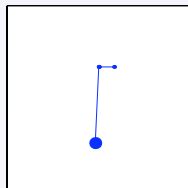
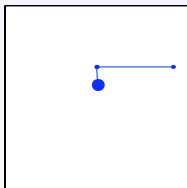


oscillations

oscillations

Definition

$$\ddot{\theta}(t) + (g/L)\theta(t) = 0$$



- *Constant length* $L = L_0$ — Small oscillations are sinusoidal, with *constant* period $T_0 = 2\pi\sqrt{L_0/g}$
- *"Slowly-varying" length* $L = L(t)$ — Small oscillations are almost-sinusoidal, with *varying* pseudo-period $T(t) \sim 2\pi\sqrt{L(t)/g}$



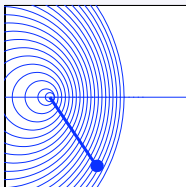
waves

waves

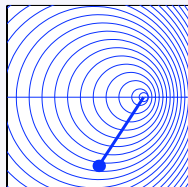
Observation

Moving monochromatic source \Rightarrow *differential* perception of the emitted frequency

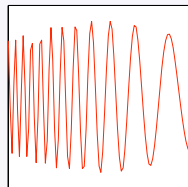
$f + \Delta f$



$f - \Delta f$



"chirp"



chirps in time

Definition

We will call "**chirp**" any complex-valued signal of the form $x(t) = a(t) \exp\{i\varphi(t)\}$, where $a(t) \geq 0$ is a low-pass amplitude whose evolution is slow as compared to the oscillations of the phase $\varphi(t)$

- Slow evolution ?
- Usual heuristic conditions assume that
 - 1 $|\dot{a}(t)/a(t)| \ll |\dot{\varphi}(t)|$: the amplitude is **almost-constant** at the scale of one pseudo-period $T(t) = 2\pi/|\dot{\varphi}(t)|$
 - 2 $|\ddot{\varphi}(t)|/\dot{\varphi}^2(t) \ll 1$: the pseudo-period $T(t)$ is itself **slowly varying** from one oscillation to the next

chirps in frequency

Theorem (Stationary phase principle)

Assuming that $\dot{\varphi}(t)$ has monotonic variation, with t_s such that $\dot{\varphi}(t_s) = 2\pi f$, one can approach the chirp spectrum

$$X(f) = \int a(t) e^{i(\varphi(t) - 2\pi ft)} dt$$

by its **stationary phase approximation** $\tilde{X}(f) \propto a^2(t_s)/|\ddot{\varphi}(t_s)|$

Interpretation

The “instantaneous frequency” curve $\dot{\varphi}(t)$ defines a **one-to-one** correspondence between one time and one frequency. The spectrum follows by weighting frequencies with durations

towards AM-FM

Observation

Given the *harmonic* model $x(t) = a \cos(2\pi f_0 t + \varphi_0)$,
unambiguous definition of amplitude a and frequency f_0

Aim

Switch to an *evolutive* model $x(t) = a(t) \cos \varphi(t)$, with $a(t)$
 time-varying and $\varphi(t)$ non linear

Problem

Given one observation, *no unicity* anymore for the
 representation since, for any function $0 < b(t) < 1$,

$$a(t) \cos \varphi(t) = [a(t)/b(t)] [b(t) \cos \varphi(t)] =: \tilde{a}(t) \cos \tilde{\varphi}(t)$$

from monochromatic waves...

Observation

The *real-valued* harmonic model can indeed be written

$$x(t) = a \cos(2\pi f_0 t + \varphi_0) = \operatorname{Re} \{ a \exp i(2\pi f_0 t + \varphi_0) \},$$

with

$$a \exp i(2\pi f_0 t + \varphi_0) = x(t) + i(\mathbf{H}x)(t)$$

and \mathbf{H} the *Hilbert transform* (quadrature)

Interpretation

A monochromatic wave (prototype of a deterministic "stationary" signal) is described, in the complex plane, by a *rotating vector* whose magnitude and rotation speed are *time-invariant* quantities

... to AM-FM

Idea

Go to the complex plane and describe an AM-FM waveform by a *rotating* (Fresnel) vector whose magnitude and rotation speed are *time-varying* quantities, while mimicking the monochromatic construction

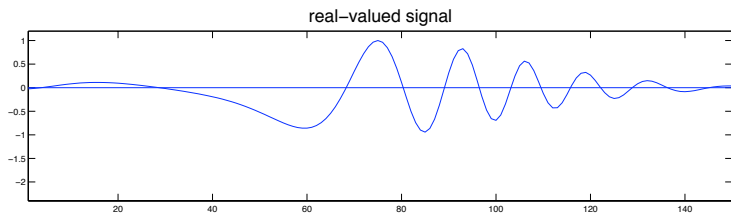
$$x(t) \rightarrow z_x(t) := x(t) + i(\mathbf{H}x)(t)$$

Definition (Gabor, '46; Ville, '48)

The instantaneous amplitude and frequency follow from this complex-valued representation — referred to as *analytic signal* — as

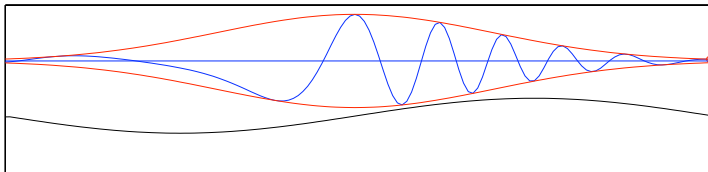
$$a_x(t) := |z_x(t)| \quad ; \quad f_x(t) := \frac{1}{2\pi} \frac{d}{dt} \arg z_x(t)$$

example

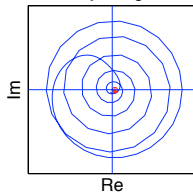


example

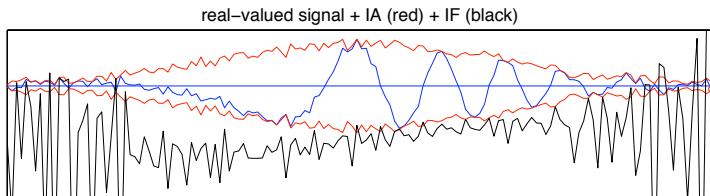
real-valued signal + IA (red) + IF (black)



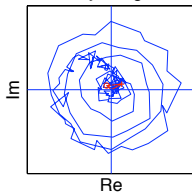
analytic signal



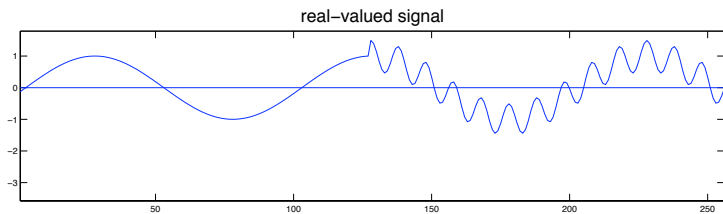
limitation: noise



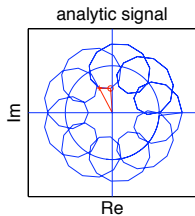
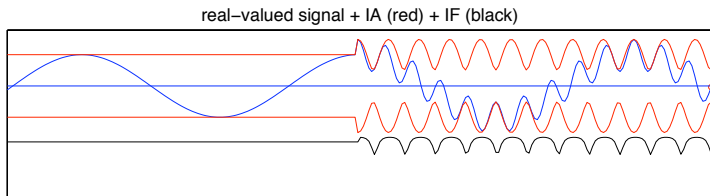
analytic signal



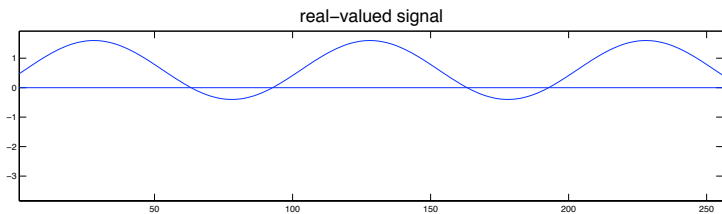
assumption 1: monocomponent



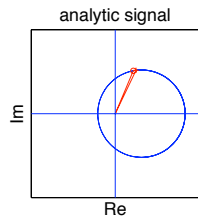
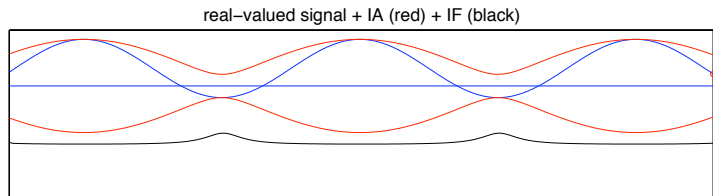
assumption 1: monocomponent



assumption 2: zero-mean



assumption 2: zero-mean



alternatives

Ideas

- ① *Teager, '86 & Kaiser, '90: define IA-IF from a local **energy operator***
- ② *Huang, '98: pre-process the observation by an **Empirical Mode Decomposition** so as to get well-behaved modes for IA-IF extraction*
- ③ *Equis et al., '11: consider rotations relatively to a local, moving center and deduce IA-IF from the estimation of an **osculating circle***

wedding time and frequency

Aim

t or f (Fourier) $\rightarrow f(t)$ (Fresnel) $\rightarrow t$ **and** f (time-frequency)

Problem (Heisenberg, '25; Gabor, '46)

Localization trade-off, classically based on a second order (variance-type) measure: $\Delta t_x \Delta f_x \geq \|x\|/4\pi (> 0)$, with $\Delta t_x = (\int t^2 |x(t)|^2 dt)^{1/2}$ and $\Delta f_x = (\int f^2 |X(f)|^2 df)^{1/2}$

Interpretation

No perfect pointwise localization

Remark

Same limitation with other spreading measures, e.g., entropy (Hirschman, '57). Common denominator: minimum achieved with **Gaussians**

chirps as uncertainty minimizers

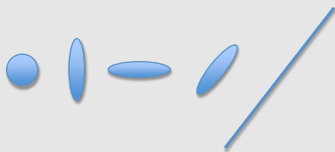
Remark

No pointwise localization does not mean no localization

Stronger uncertainty relation (Schrödinger, 1935)

$$\Delta t_x \Delta f_x \geq \frac{\|x\|}{4\pi} \sqrt{1 + 16\pi^2 \left(\int t (\partial_t \arg x(t)) |x(t)|^2 dt \right)^2}$$

bound achieved for **“squeezed states”** of the form $\{\exp(\alpha t^2 + \beta t + \gamma)\}$, with **linear “chirps”** as a limit when $\text{Re}\{\beta\} = 0$ and $\text{Re}\{\alpha\} \rightarrow 0_-$



time-frequency alternatives

From stationarity...

Spectrum analysis “à la Wiener-Khintchine-Bochner” :
 $\Gamma_x(f) = \mathcal{F}\{\gamma_x\}(f)$, with $\gamma_x(\tau) := \langle x, \mathbf{T}_\tau x \rangle$ *correlation function*
independent of time

... to nonstationarities (Wigner, '32; Ville, '48)

$\gamma_x \rightarrow$ *time-frequency correlation* $\langle x, \mathbf{T}_{\tau,\xi} x \rangle$ + 2D Fourier transform \Rightarrow *Wigner-type transforms*

- *intrinsic definitions: no dependence on some measurement device (window, wavelet)*
- *perfect localization for linear chirps (with possible extensions to non linear cases)*

“distribution/correlation” duality

Definition

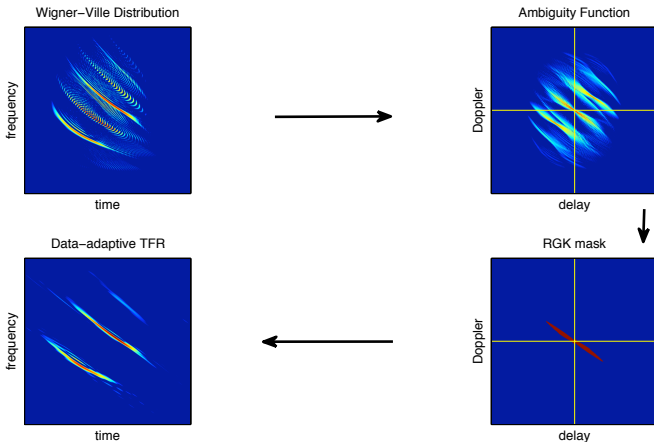
The 2D Fourier transform $A_x(\xi, \tau)$ of the Wigner distribution $W_x(t, f)$ is referred to as the (narrowband) **ambiguity function (AF)**

Interpretation

The TF-shift operator $(\mathbf{T}_{\xi, \tau} x)(t) := x(t - \tau) e^{-i2\pi\xi(t - \tau/2)}$ is such that $A_x(\xi, \tau) = \langle x, \mathbf{T}_{\xi, \tau} x \rangle \Rightarrow$ **AF = TF correlation**, with

- **“auto-terms”** neighbouring the origin of the plane
- **“cross-terms”** at a distance from the origin that equals the TF distance between components

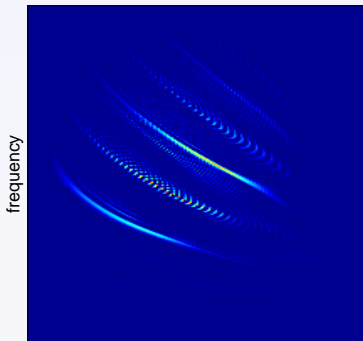
the other trade-off and its “classical” way out



from Wigner-Ville to spectrogram, and back

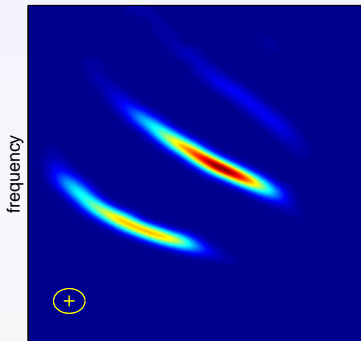
spectrogram = smoothed Wigner

Wigner-Ville



time

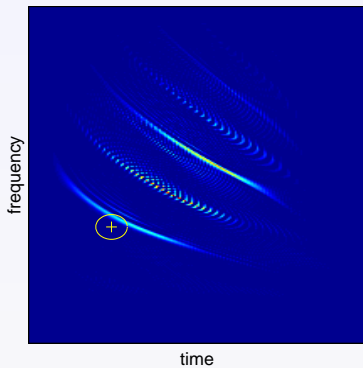
spectrogram



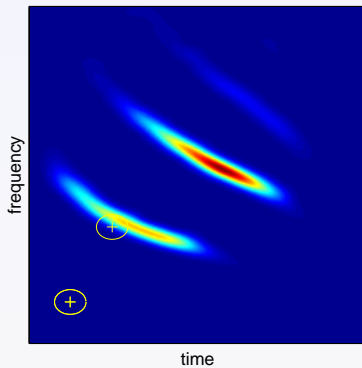
time

spreading of auto-terms

Wigner-Ville

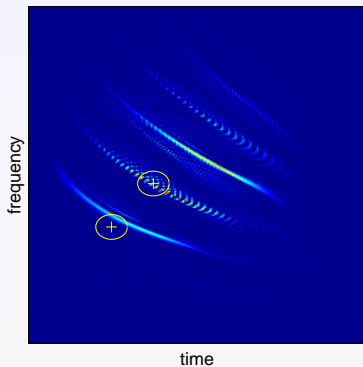


spectrogram

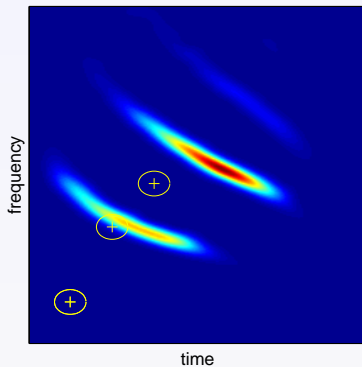


cancelling of cross-terms

Wigner-Ville

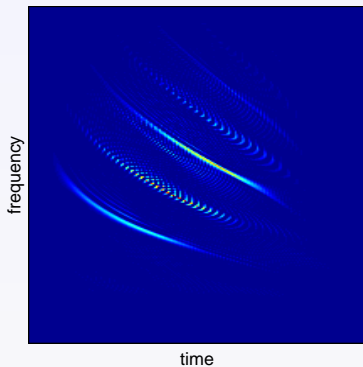


spectrogram

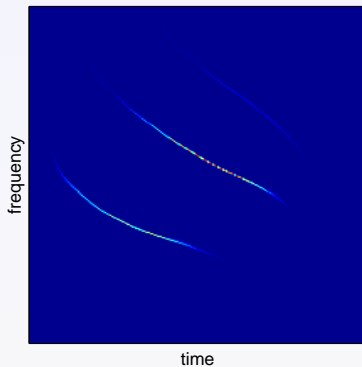


reassignment (Kodera *et al.*, '76, Auger & F., '95)

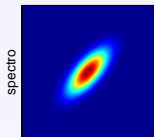
Wigner-Ville



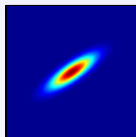
reassigned spectrogram



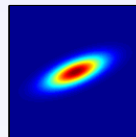
independence w.r.t. window size



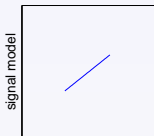
window = 21



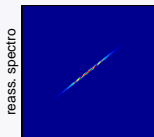
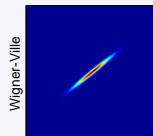
63



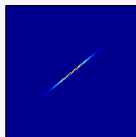
127 points



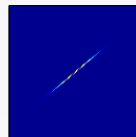
128 points



window = 21



63



127 points

a “compressed sensing” approach

Discrete time

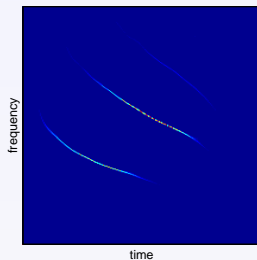
*signal of dimension $N \Rightarrow TF$
distribution of dimension $\approx N^2$*

Few components

*$K \ll N \Rightarrow$ at most $KN \ll N^2$ non zero
values in the TF plane*

Sparsity

*minimizing the ℓ_0 “norm” not feasible, but almost optimal
solution by minimizing the ℓ_1 norm*



a “compressed sensing” approach

Idea (F. & Borgnat, '08-10)

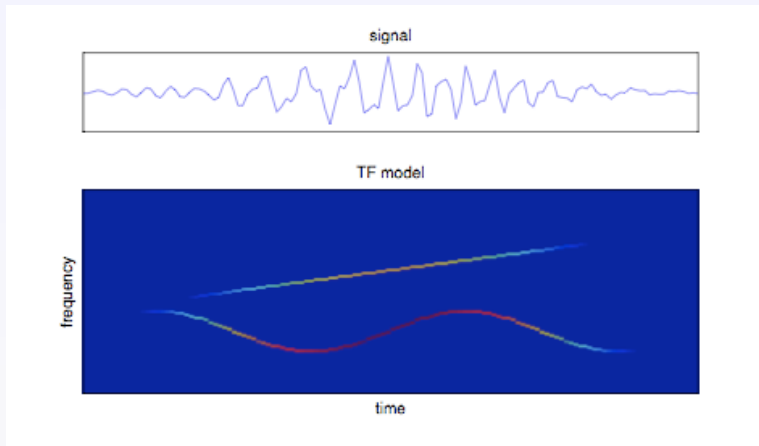
- ① *choose a domain Ω neighbouring the origin of the AF plane*
- ② *solve the program*

$$\min_{\rho} \|\rho\|_1 ; \mathcal{F}\{\rho\} - \mathbf{A}_x = \mathbf{0} |_{(\xi, \tau) \in \Omega}$$

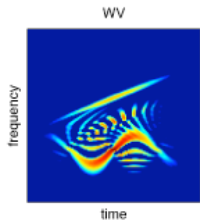
- ③ *the exact equality over Ω can be relaxed to*

$$\min_{\rho} \|\rho\|_1 ; \|\mathcal{F}\{\rho\} - \mathbf{A}_x\|_2 \leq \epsilon |_{(\xi, \tau) \in \Omega}$$

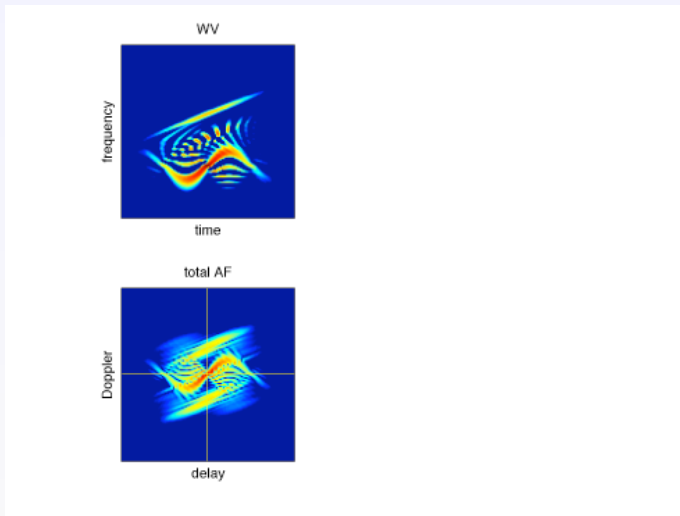
a toy example



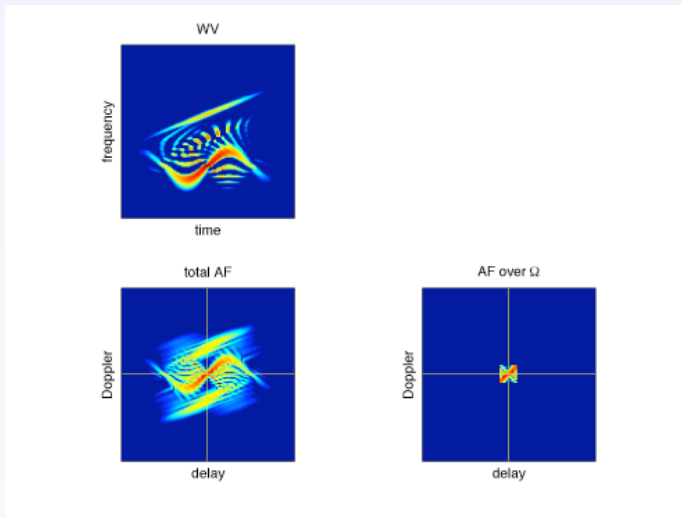
Wigner



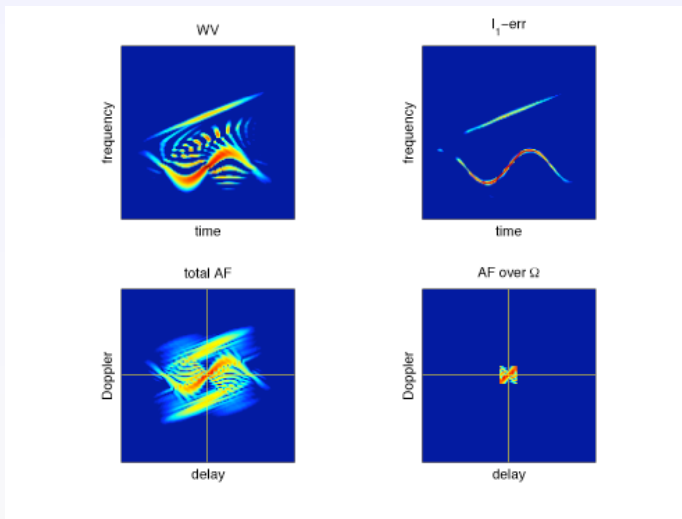
ambiguity



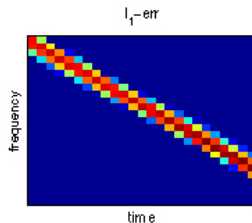
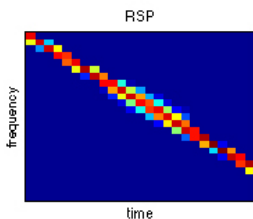
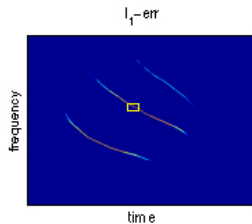
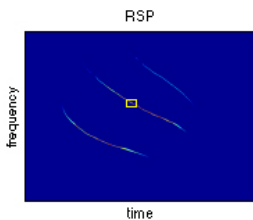
selection



sparse solution



comparison sparsity vs. reassignment



- 1 **Biology** — Bats echolocation calls
- 2 **Physics** — Gravitational waves
- 3 **Mathematics** — Riemann and Weierstrass functions

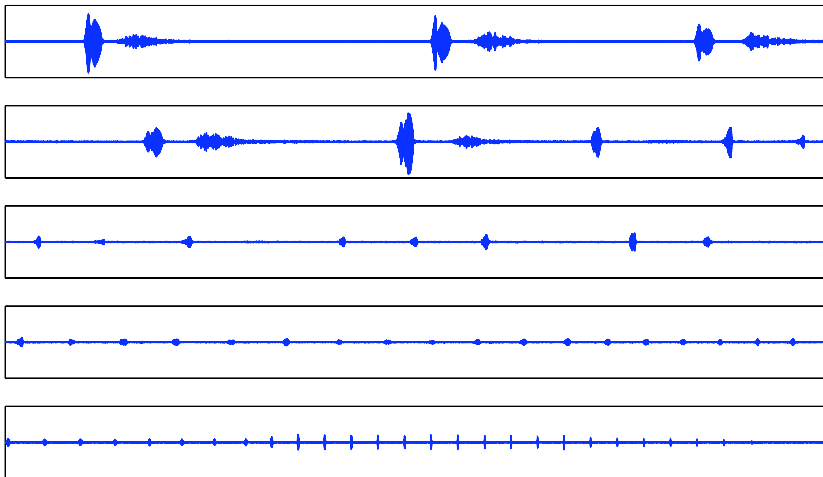


bat sonar system

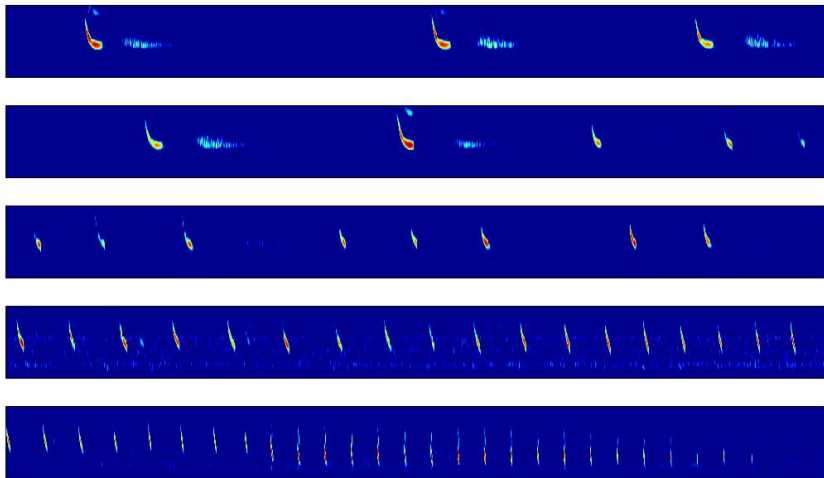
Observation

- **Echolocation** — *Active system for navigation, natural airborne sonar*
- **Signals** — *Ultrasonic acoustic waves, transient (some ms) and wideband (some tens of kHz between 40 and 100kHz) chirp signals*
- **Performance** — *Close to optimality, with adaptation of the waveforms to multiple tasks (detection, estimation, recognition, interferences rejection, . . .)*

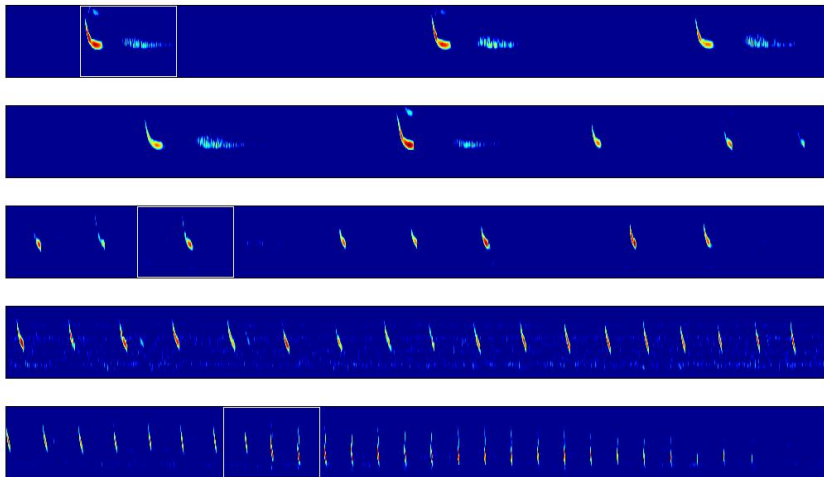
a typical pursuit sequence (*Myotis mystacinus*)



a typical pursuit sequence (*Myotis mystacinus*)

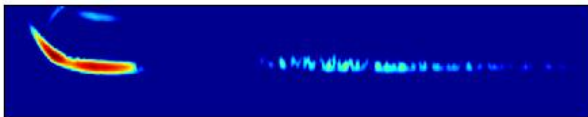


a typical pursuit sequence (*Myotis mystacinus*)

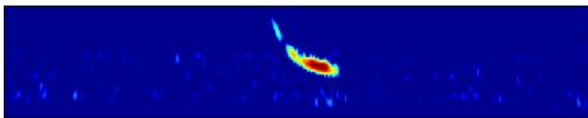


close-up (spectrogram)

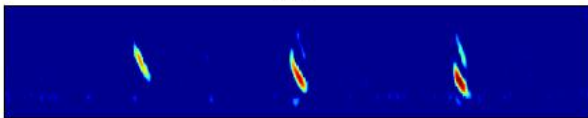
cruise



pursuit



catch

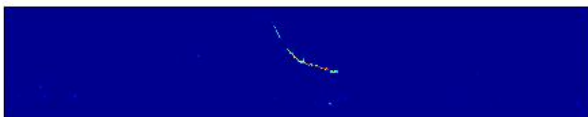


close-up (reassigned spectrogram)

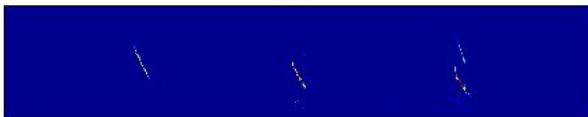
cruise



pursuit



catch



bats and signal processing

Aim

Understand the *signal design* of bat echolocation calls:

- 1 evolution within a sequence?
- 2 adaptation to environment?
- 3 optimality if any?
- 4 ...

Result

For an emitted signal of duration T and bandwidth B , the accuracy in estimating the Doppler shift and the delay of the returning echo is roughly given by $\delta f \sim T^{-1}$ and $\delta t \sim (B \sqrt{\text{SNR}})^{-1}$ (Woodward's formula)

bats and signal processing

Observation

- ① **Cruise** — Importance of estimating both distance (delay) and speed (Doppler)
⇒ *broadband* chirp + *almost Constant Frequency* part
- ② **Pursuit** — Importance of estimating distance whatever the Doppler
⇒ *adapted* chirp + *progressive suppression* of the almost Constant Frequency part
- ③ **Catch** — Importance of precise localization with shorter pulses
⇒ increasing the effective bandwidth B by *lowering* the fundamental and increasing distortion (*harmonics*)

chirp detection/estimation

Interpretation (time-frequency)

- ① **Matched filtering** — Emitted signal $s(t)$ as a template + echo $e(t)$ as a delayed version of $s(t)$ embedded in wGn
 \Rightarrow optimal estimation of delay by $\hat{\tau} = \arg \max_{\tau} |\langle e, \mathbf{T}_{\tau} s \rangle|^2$
- ② **Unitarity** — Inner product equivalence
 \Rightarrow (Moyal's formula) $|\langle x, y \rangle|^2 = \iint \rho_x(t, f) \rho_y(t, f) dt df$
- ③ **Localization** — Energy along instantaneous frequency $f_s(t)$
 $\Rightarrow \hat{\tau} = \arg \max_{\tau} \int a_s^2(t - \tau) \rho_e(t, f_s(t - \tau)) dt$
path integration in the time-frequency plane

optimal integration (Wigner-Ville)

approximation (reassigned spectrogram)

robustness

Doppler tolerance

Problem

Coupled errors in the joint estimation of delay and Doppler

Way out (Altes & Titlebaum, '70)

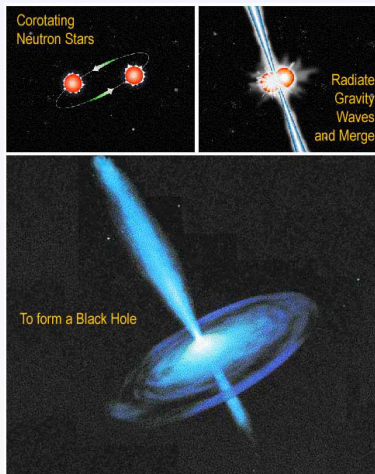
- A waveform is said to be *Doppler-tolerant* if it permits an unbiased estimation of delay whatever the (unknown) Doppler
- For broadband signals, Doppler has to be considered as a *dilation* (shift = approximation for narrowband signals)
- Analytic solution = *hyperbolic chirps*

Doppler tolerance

Graphical solution



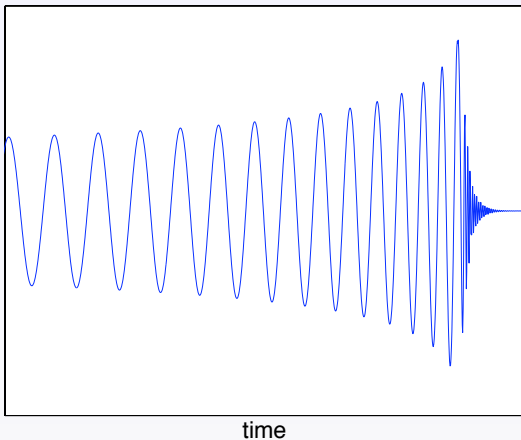
coalescing binaries



coalescing binaries

expected waveform

gravitational wave



the quest

Challenge

First direct proof on Earth of the existence of gravitational waves (GWs), as predicted by general relativity

- **Projects** — VIRGO (France-Italy) + LIGO (USA): giant Michelson interferometers (~ 3 km long arms)
- **Measurements** — GWs impinging the interferometer modify locally the space-time geometry and result in a differential variation of the arms length \Rightarrow interference fringes
- **Difficulties** — Signals are very weak and can be efficiently observed only in a very short time window (some seconds) corresponding to the frequency window above ~ 10 Hz (seismic noise) and below ~ 1 kHz (photon noise)

GW detection/estimation

Model

The inspiral part of the GW radiated by a coalescing binary made of two objects of respective masses m_1 and m_2 can be modelled as a **power-law chirp**

$$C_{\alpha,\beta}(t) = a(t_c - t)^\alpha \exp\{i(b(t_c - t)^\beta + c)\} U(t_c - t),$$

with $(\alpha, \beta) = (-1/4, 5/8)$

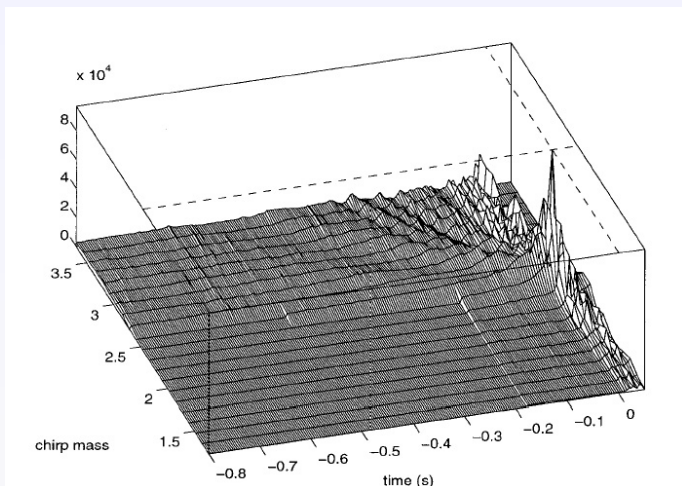
Parameters

① the **coalescence time** t_c

② the **chirp mass** defined as

$\mathcal{M} = (m_1 + m_2)^{2/5} (m_1^{-1} + m_2^{-1})^{-3/5}$, and related to the “chirp rate” b according to $b \approx 38.6 (\mathcal{M}/M_\odot)^{-5/8}$, where M_\odot stands for the solar mass

reassigned spectrogram + path integration



(Chassande-Mottin & F., *ACHA* '98)



Bernhard Riemann (1826-1846)

a very special function

Definition

$$\sigma(t) := \sum_{n=1}^{\infty} n^{-2} \sin \pi n^2 t$$

Result

$\sigma(t)$ non differentiable if $t \neq t_0 = (2p + 1)/(2q + 1)$, $p, q \in \mathbb{N}$
(Hardy, '16) but differentiable in $t = t_0$ (Gerver, '70)

Theorem (Meyer, '96)

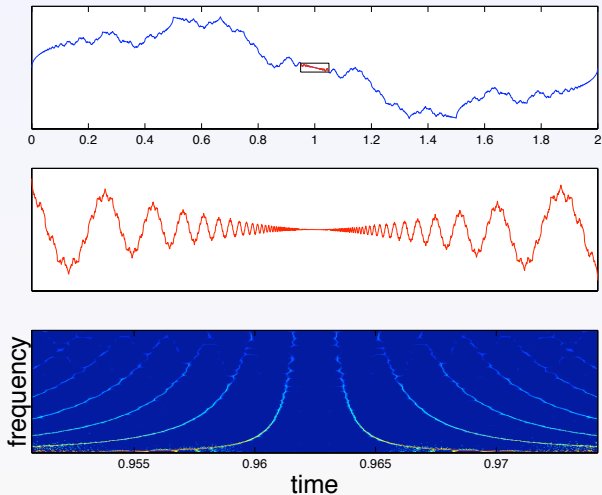
In the vicinity of $z = 1$, the holomorphic version of Riemann's function can be expressed as a combination of *local chirps*:

$$\sigma(1 + z) = \sigma(1) - \pi z/2 + \sum_{n=1}^{\infty} K_n(z) C_{3/2, -1}(z),$$

leading to $\sigma(1 + t) = \sigma(1) - \pi t/2 + O(|t|^{3/2})$ when $t \rightarrow 0$

power-law chirps

Riemann function





from Fourier to Mellin

Definition

The **Mellin Transform** (MT) of a signal $x(t) \in L^2(\mathbb{R}^+, t^{-2\alpha+1} dt)$ can be defined as the projection:

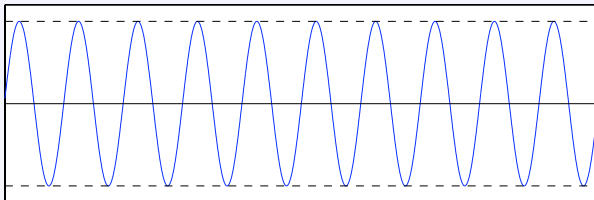
$$(\mathcal{M}x)(s) := \int_0^{+\infty} x(t) t^{-i2\pi s - \alpha} dt =: \langle x, c \rangle$$

Interpretation

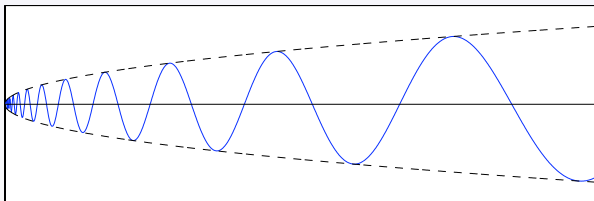
- ① Analysis over **hyperbolic** chirps $c(t) := t^{-\alpha} \exp\{i2\pi s \log t\}$
- ② $\dot{\varphi}_c(t)/2\pi = s/t \Rightarrow$ the Mellin parameter s can be interpreted as a **hyperbolic modulation rate**
- ③ The MT can also be viewed as a **warped FT**, since $\tilde{x}(t) := e^{(1-\alpha)t} x(e^t) \Rightarrow (\mathcal{M}x)(s) = (\mathcal{F}\tilde{x})(s)$

Mellin as warped Fourier

tone



chirp





Karl Weierstrass (1815-1897)

another very special function

Definition (Weierstrass, 1872)

Geometrically spaced Fourier modes

$$W(t) := \sum_{n=0}^{\infty} \lambda^{-nH} \cos \lambda^n t, \lambda > 1$$

Definition (Mandelbrot, 1977)

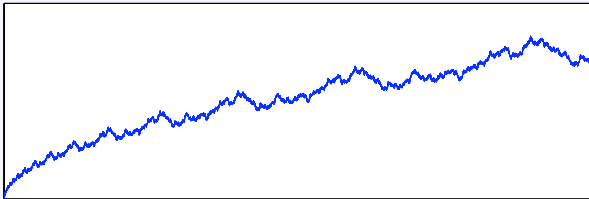
$$W_g(t) := \sum_{n=0}^{\infty} \lambda^{-nH} (g(0) - g(\lambda^n t)) e^{i\varphi_n}, \lambda > 1,$$

with $g(\cdot)$ 2π -periodic and $\varphi_n \in \mathcal{U}(0, 2\pi)$

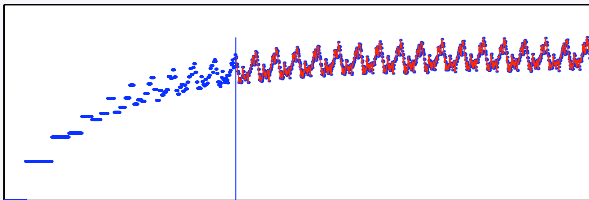
see also (Berry & Lewis, *Proc. Roy. Soc. London A*, 1980)

“de-warping” (Lamperti, '62)

Weierstrass function ($H = 0.5$)



"Delampertized" Weierstrass function



Weierstrass meets Mellin

Result (Borgnat & F., '03)

*The Weierstrass-Mandelbrot admits the equivalent **Mellin** decomposition:*

$$W_g(t) = \sum_{m=-\infty}^{\infty} \frac{(\mathcal{M}_H G)(m/\log \lambda)}{\log \lambda} m_{H, m/\log \lambda}(t),$$

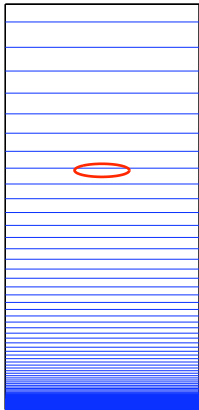
with $G(t) := g(0) - g(t)$

Interpretation

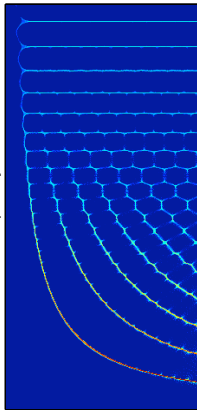
*Natural co-existence of **two readings** (Fourier and Mellin) in the time-frequency plane*

Fourier vs. Mellin

Fourier model

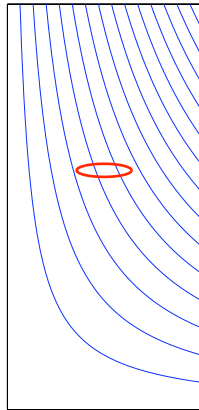


frequency

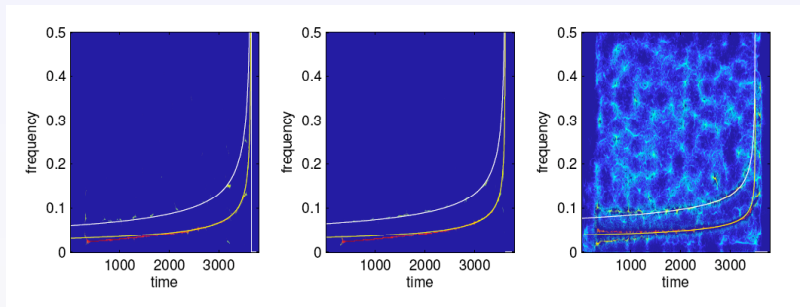


time

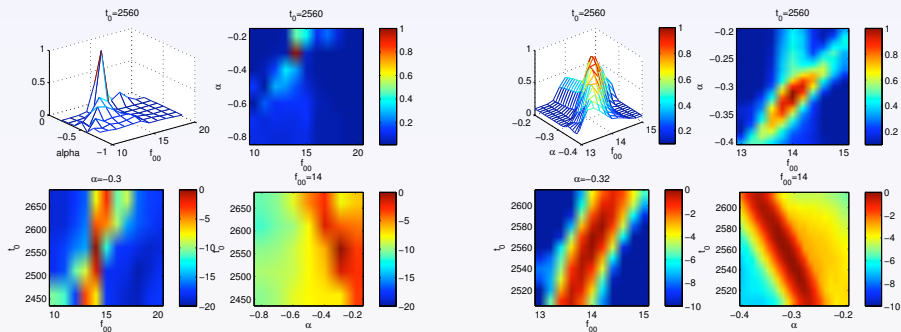
Mellin model



back to Euler



Hough transform



chirps “everywhere”

- bird songs
- ocean waves
- “whistling atmospherics”
- wideband impulses propagating in a dispersive medium
- vibroseismics
- EEG (epileptic seizure)
- uterine EMG
- coherent structures in turbulence
- precursors accumulation in earthquakes
- “speculative bubbles” prior a financial crash
- ...

(p)reprints, Matlab codes & contact

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- `Patrick.Flandrin@ens-lyon.fr`