# Wavelets and scaling processes

#### Patrick Flandrin

CNRS & École Normale Supérieure de Lyon, France



thanks to Patrice Abry, Paulo Gonçalves, Darryl Veitch & Herwig Wendt

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# outline

- Examples of "real life" scaling processes
- Scaling concepts (self-similarity, fractality, long-range dependence, "1/f") and simple models (fractional Brownian motion, FARIMA, ON/OFF)
- ③ Wavelet basics
- A wavelet-based framework for scaling processes (analysis and synthesis)
- 5 Extension to multifractal processes
- 6 References

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### random walk

#### Definition

Brownian motion is the zero-mean Gaussian process  $\{X(t), t \ge 0; X(0) = 0\}$  with increments  $\{\Delta_X(t, s) := X(t) - X(s), t \ge s\}$  such that 1 var $\Delta_X(t, t - 1) = 2D$ , for any t (stationarity) 2  $\mathbb{E}\Delta_X(t_1, t_2)\Delta_X(t_3, t_4) = 0$ , if  $t_1 > t_2 \ge t_3 > t_4$ (independence)

Ordinary diffusion — We then have (Einstein, 1905):

$$\operatorname{var} X(t) = \mathbb{E}[\Delta_X(t, t-1) + \Delta_X(t-1, t-2) + \ldots + \Delta_X(1, 0)]^2$$
  
=  $t \times \operatorname{var} \Delta_X(1, 0)$   
=  $2Dt \Rightarrow \sigma_X(t) := (\operatorname{var} X(t))^{\frac{1}{2}} \propto \sqrt{t}$ 



credit: Leonardo da Vinci

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### disordered fluids

Prediction (Kolmogorov, '41)

Assuming homogeneity, isotropy and a constant energy transfer rate in turbulent flows leads to the following (K41) predictions for the velocity field v(x) and energy spectrum E(k):

$$\langle |v(x+r)-v(x)|^q 
angle \propto r^{q/3} \stackrel{(q=2)}{\Rightarrow} E(k) \propto k^{-5/3}$$

- Experiments Observed spectra are "in  $k^{-5/3}$ ", but scaling laws have anomalous exponents  $\zeta_q \neq q/3$
- Issues Reconsidering initial hypotheses, in terms of statistics (non-Gaussianity...) and events (coherent structures...)

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# velocity spectrum



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## complex systems



### network traffic

Observation (Leland & et al., '94)

Experimental evidence of scaling phenomena and long-range dependence in Ethernet traffic

- Extensions Similar behaviors observed in other types of traffic (Internet, VBR, WAN ...).
- Issues Traffic control (congestion, dimensioning, anomalies detection, prediction, protocols ...) ⇒ modeling and analysis.

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### **IP** Packets



credit: D. Veitch

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### Nature and finance!





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### scaling ubiquitous in many domains

- Solid-state physics
- Hydrology
- Astrophysics
- Heart-rate variability
- Brain activity (fMRI)
- Earthquakes
- Social networks
- etc.

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### scaling and power-laws

 Power-law spectra — Power-laws correspond to homogeneous functions:

$$\mathcal{S}(f) = \mathcal{C} |f|^{-\alpha} \Rightarrow \mathcal{S}(kf) = \mathcal{C} |kf|^{-\alpha} = k^{-\alpha} \mathcal{S}(f),$$

for any k > 0

Fourier transform — Frequency scaling carries over to the time domain. If we let s(t) := (𝒴<sup>-1</sup>𝔅)(f), we get:

$$\int \mathcal{S}(kf) e^{i2\pi ft} df = k^{-1} \int \mathcal{S}(f') e^{i2\pi f'(t/k)} df' = s(t/k)/k$$

It follows that  $s(t/k) = s(t)/k^{\alpha-1} \Rightarrow self-similarity$ 



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credit: S. Kim

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### no characteristic scale

#### Definition

A process  $\{X(t), t \in \mathbb{R}\}$  is said to be self-similar of index H (or "H-ss") if, for any k > 0,

$$\{X(kt), t \in \mathbb{R}\} \stackrel{d}{=} k^{H}\{X(t), t \in \mathbb{R}\}.$$

#### Interpretation

Any zoomed (in or out) version of an H-ss process looks (statistically) the same

#### Remark

If a process X is self-similar, it is necessarily nonstationary

#### zooming in on an *H*-ss process

credit: H. Wendt

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### stationary increments

Definition

A process  $\{X(t), t \in \mathbb{R}\}$  is said to have stationary increments if and only if, for any  $\theta \in \mathbb{R}$ , the increment process:

$$\left\{X^{( heta)}(t) := X(t+ heta) - X(t), t \in \mathbb{R}
ight\}$$

has a distributional law which does not depend upon t

#### Remark

The concept of stationary increments can be naturally extended to higher orders ("increments of increments")

Definition *H-ss processes with stationary increments are referred to as "H-sssi" processes* 

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#### covariance function of *H*-sssi processes

#### Theorem

The structure of the covariance function is the same for all *H*-sssi processes

*Proof* — Assuming that X(t) is *H*-sssi, with X(0) = 0 and  $X(1) \neq 0$ , we have necessarily:

$$\begin{split} \mathbb{E}X(t)X(s) &= \frac{1}{2} \left( \mathbb{E}X^2(t) + \mathbb{E}X^2(s) - \mathbb{E}\left(X(t) - X(s)\right)^2 \right) \\ &= \frac{1}{2} \left( \mathbb{E}X^2(t) + \mathbb{E}X^2(s) - \mathbb{E}\left(X(t-s) - X(0)\right)^2 \right) \\ &= \frac{\text{var}X(1)}{2} \left( |t|^{2H} + |s|^{2H} - |t-s|^{2H} \right) \end{split}$$

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# asymptotic self-similarity

#### Definition

A stationary process  $\{X(t), t \in \mathbb{R}\}$  is said to be asymptotically self-similar of index  $\beta \in (0, 1)$  if

$$(\operatorname{var} X(t))^{-1} \mathbb{E} X(t) X(t+\tau) \sim \tau^{-\beta}$$

when  $\tau \to \infty$ 

- *H*-sssi processes are asymptotically self-similar of index  $\beta = 2(1 H)$
- non-summability (and power-law decay) of the autocorrelation ⇒ (power-law) divergence of the PSD at f = 0
- asymptotic self-similarity ⇒ long-range dependence (LRD) (also referred to as long memory)

### fractional Brownian motion

#### Definition 1

A process  $B_H(t)$  is referred to as a fractional Brownian motion (fBm) of index 0 < H < 1, if and only if it is H-sssi and Gaussian

- fBm has been introduced in (Mandelbrot & van Ness, '68), as an extension of the ordinary Brownian motion  $B(t) \equiv B_H(t)|_{H=1/2}$  (anomalous diffusion)
- the index *H* is referred to as the *Hurst exponent*, and its limited range guarantees the *non-degeneracy* (*H* < 1) and the *mean-square continuity* (*H* > 0) of fBm

### fractional Brownian motion

Definition 2 *fBm admits the moving average representation:*  $B_{H}(t) - B_{H}(0) = \frac{1}{\Gamma(H + \frac{1}{2})} \left\{ \int_{-\infty}^{0} [(t - s)^{H - \frac{1}{2}} - (-s)^{H - \frac{1}{2}}] B(ds) + \int_{0}^{t} (t - s)^{H - \frac{1}{2}} B(ds) \right\}$ 

- fBm results from a *"fractional integration"* of white noise
- no specific role attached to time t = 0

### fractional Brownian motion

#### Definition 3

fBm admits the (harmonizable) spectral representation:

$$B_{H}(t) = C \int_{-\infty}^{+\infty} |f|^{-(H+\frac{1}{2})} (e^{i2\pi tf} - 1) W(df),$$

with W(df) the Wiener measure

- the "average spectrum" of fBm behaves as  $|f|^{-(2H+1)}$
- fBm is a widespread model for (nonstationary) Gaussian processes with a *power-law* (empirical) spectrum

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### fractional Gaussian noise

#### Definition

The (stationary) increment process  $B_H^{(\theta)}(t)$  of fBm  $B_H(t)$  is referred to as fractional Gaussian noise (fGn)

*Autocorrelation* — The (stationary) autocorrelation function of fGn,  $c_H(\tau) := \mathbb{E}B_H^{(\theta)}(t)B_H^{(\theta)}(t+\tau)$ , reads:

$$c_{\mathcal{H}}(\tau) = \frac{\sigma^2}{2} \left( |\tau + \theta|^{2\mathcal{H}} - 2|\tau|^{2\mathcal{H}} + |\tau - \theta|^{2\mathcal{H}} \right)$$

- if  $\theta = 1$  and  $H = \frac{1}{2}$ , we have  $c_H(k) = \sigma^2 \, \delta(k), k \in \mathbb{Z}$ (discrete-time white noise)
- for large lags  $\tau$ , one has  $c_H(\tau) \sim \sigma^2 \theta^2 H(2H-1)\tau^{2(H-1)}$ (subexponential, *power-law* decay)

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### fGn autocorrelation function



### fGn spectrum

#### Result

If  $\theta = 1$ , the power spectrum density of discrete-time fGn is given by:

$$S(f) = C \sigma^2 |e^{i2\pi f} - 1|^2 \sum_{k=-\infty}^{\infty} \frac{1}{|f+k|^{2H+1}},$$

with  $-\frac{1}{2} \le f \le +\frac{1}{2}$ 

if H ≠ 1/2, we have S(f) ~ C σ<sup>2</sup> |f|<sup>1−2H</sup> when f → 0
 0 < H < 1/2 ⇒ S(0) = 0</li>
 1/2 < H < 1 ⇒ S(0) = ∞ (spectral divergence)</li>

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#### fGn power spectrum density



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### importance of fGn

Aggregation — Renormalized by T<sup>H-1</sup>, fGn is invariant under the recomposition rule

$$X(n) \mapsto X^{T}(n) := T^{-1} \sum_{k=(n-1)T+1}^{nT} X(k)$$

- ② Attraction As T → ∞, aggregating any asymptotically H-ss process ends up with a process whose covariance structure is that of fGn
- 3 Long-range dependence fGn is LRD when  $\frac{1}{2} < H < 1$

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### Bm sample paths



# fBm sample paths 1

#### Definition

The local regularity of a signal x(t) at a given point  $t_0$  is measured by the Hölder exponent  $h(t_0)$  defined as the supremum of  $\alpha$ 's such that  $|x(t) - x(t_0)| < C |t - t_0|^{\alpha}$  when  $|t - t_0| \rightarrow 0$  (the larger the exponent, the smoother the signal)

#### Result

For any (small enough)  $\epsilon > 0$  and any  $t \in \mathbb{R}$ , we have  $|B_{H}^{(\epsilon)}(t)| \leq C |\epsilon|^{H}$ , with probability 1

- fBm is everywhere continuous, but nowhere differentiable
- sample paths have a *uniform* Hölder regularity h = H
- sample paths have a uniform (Haussdorf and box) *fractal dimension* dim<sub>B</sub> graph  $B_H = 2 - H$

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#### examples of fBm sample paths



### correlation between increments

# Interpretation It follows from the covariance structure of fBm that, for any $t \in \mathbb{R}$ , $\pi P(=\theta)(x) P(\theta)(x)$

$$C_{H}(\theta) := -rac{\mathbb{E}B_{H}^{(-\theta)}(t)B_{H}^{(\theta)}(t)}{\operatorname{var}B_{H}^{(\pm\theta)}(t)} = 2^{2H-1} - 1$$

- $H = \frac{1}{2}$ : no correlation (Brownian motion, D = 1.5)
- $H < \frac{1}{2}$ : negative correlation (more erratic,  $\lim_{H\to 0} D = 2$ )
- $H > \frac{1}{2}$ : *positive* correlation (less erratic,  $\lim_{H \to 1} D = 1$ )

Interpretation

H is a roughness measure of sample paths

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### "1/*f*" processes

#### Definition

A process is said to be of "1/f"-type if its empirical PSD behaves as  $f^{-\alpha}$  ( $\alpha > 0$ ) over some frequency range [A, B]

Depending on A and B, one can end up with:

- *LRD*, if  $A \rightarrow 0$  and  $B < \infty$
- scaling in some "inertial range", if  $0 < A < B < \infty$
- small-scale *fractality*, if  $A < \infty$  and  $B \rightarrow \infty$

#### Remark

In the fBm case, the only Hurst exponent H controls all 3 situations

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# **ON/OFF Models 1**

**ON/OFF sources** — Let  $S_N(t)$  result from the superposition of *N* independent *ON/OFF* sources:  $S_N(t) := \sum_{i=1}^{N} X_i(t)$ 



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### ON/OFF Models 2

Observation Given N and T, the statistical properties of the cumulative process:

$$Y_N(Tt) := \int_0^{Tt} S_N(s) \, ds$$

depend on the pdf's of the ON and OFF periods au

Theorem (Willinger *et al.*, '95)

Assuming that the ON and OFF periods are both heavy-tailed, i.e., that  $Prob\{\tau > x\} \sim c x^{-\alpha}, \tau \to \infty$ , with  $1 < \alpha < 2$ , then (up to a linear trend)  $Y_N(Tt)$  behaves asymptotically as fBm of Hurst exponent  $H = (3 - \alpha)/2$  when N and  $T \to \infty$ 

### FARIMA

Definition (Hosking, '81) In the discrete-time case where  $k \in \mathbb{Z}$ , Fractional Auto-Regressive Integrated Moving Average (FARIMA) processes X(k) are defined by

 $\Phi(z)(1-z^{-1})^d X(k) = \Theta(z) W(k),$ 

where z is the unit delay operator,  $\Phi$  and  $\Theta$  are polynomial of order p and q, respectively,  $d \in ]-1/2, +1/2[$  is the order of fractional differencing and W(k) is white Gaussian noise

Interpretation

FARIMA(p, d, q) processes generalize the classical random walk, which is a FARIMA(0, 1, 0)

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## importance of FARIMA

**1** FARIMA(0, d, 0) — In the simplest case of a *fractionally integrated white noise*, the spectrum reads  $S(f) \propto |\sin \pi f|^{-2d}$  and

$$\mathcal{S}(f) \sim |f|^{-2d}, \ f 
ightarrow 0$$

- 2 *fGn-like* At low frequencies, the spectrum is identical to that of fGn, with the identification  $d \equiv H 1/2$
- 3 Long-range dependence FARIMA(0, d, 0) processes are suitable models for *asymptotic self-similarity*, and *long-range dependence* when d > 0

## evidencing scaling in data ?

#### Observation

- Different and complementary signatures of scaling can be observed with respect to time (correlation, fractality,...) or frequency/scale (spectrum, zooming, ...)
- Iterating aggregation reveals scale invariance

#### Idea

- Use explicitly an approach which combines time and frequency/scale
- Use explicitly a multiresolution approach

#### $\Rightarrow$ Wavelets !

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#### $\Rightarrow$ Wavelets !

## rationale

#### Idea "signal = (low-pass) approximation + (high-pass) detail" + iteration

- successive approximations (at coarser and coarser resolutions)  $\sim$  aggregated data
- details (information differences between successive resolutions)  $\sim$  *increments*

Multiresolution is a natural language for scaling processes

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## formalization

#### Definition (Mallat & Meyer, '86)

A MultiResolution Analysis (MRA) of  $L^2(\mathbb{R})$  is given by:

- ④ a hierarchical sequence of embedded approximation spaces ...,  $V_1 ⊂ V_0 ⊂ V_{-1}$ ..., whose intersection is empty and whose closure is dense in  $L^2(\mathbb{R})$
- a dyadic two-scale relation between successive approximations

$$X(t) \in V_j \Leftrightarrow X(2t) \in V_{j-1}$$

3 a scaling function  $\varphi(t)$  such that all of its integer translates  $\{\varphi(t-n), n \in \mathbb{Z}\}$  form a basis of  $V_0$ 

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## multiresolution expansion

#### Definition

The wavelet  $\psi(.)$  is constructed in such a way that all of its integer translates form a basis of  $W_0$ , defined as the complement of  $V_0$  in  $V_{-1}$ 

For a given resolution depth *J*, any signal  $X(t) \in V_0$  can be expanded as :



with  $\{\xi_{j,k}(t) := 2^{-j/2} \zeta(2^{-j}t - k), j \text{ and } k \in \mathbb{Z}\}$ , for  $\xi = \varphi$  and  $\psi$ 

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## a simple construction

#### Example (Haar, 1911)

The simplest choice for a MRA is given by the Haar basis, attached to the scaling function  $\varphi(t) = \chi_{[0,1]}(t)$  and the wavelet  $\psi(t) = \chi_{[0,1/2]}(t) - \chi_{[1/2,1]}(t)$ 

#### Remark

When aggregated over dyadic intervals, data samples identify to Haar approximants

Interpretation

Wavelet analysis offers a refined way of both aggregating data and computing increments

## two key properties

(1) Admissibility — By construction, a scaling function (resp., a wavelet) is a low-pass (resp., high-pass) function  $\Rightarrow$  an admissible wavelet  $\psi(t)$  is necessarily zero-mean:

$$\Psi(\mathbf{0}) := \int_{-\infty}^{+\infty} \psi(t) \, dt = \mathbf{0}$$

② Cancellation — A further key property for a wavelet is the number of its vanishing moments, i.e., the integer N ≥ 1 such that

$$\int_{-\infty}^{+\infty} t^k \,\psi(t) \,dt = 0, \text{ for } k = 0, 1, \dots N-1$$

## the example of Daubechies wavelets



## effective computation

 Theory — The wavelet coefficients d<sub>X</sub>(j, k) are given by the inner products:

$$d_X(j,k) := \langle X, \psi_{j,k} \rangle$$

- Practice They can rather be computed in a recursive fashion, via efficient pyramidal algorithms (faster than FFT)
  - no need for knowing explicitly  $\psi(t)$
  - enough to characterize a wavelet by its *filter coefficients*  $\{g(n) := (-1)^n h(1 n), n \in \mathbb{Z}\}$ , with

$$h(n) := \sqrt{2} \int_{-\infty}^{+\infty} \varphi(t) \, \varphi(2t-n) \, dt$$

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# Mallat's algorithm



## wavelets as filters

- Input-output Given the statistics of the analyzed signal, statistics of its wavelet coefficients can be derived from imput-ouput relationships of *linear filters*
- Stationary processes In the case of stationary processes with autocorrelation  $\gamma_X(\tau) := \mathbb{E}X(t)X(t + \tau)$ , stationarity carries over to wavelet sequences:

$$\mathcal{C}_X(j,n) := \mathbb{E} d_X(j,k) d_X(j,k+n) = \int_{-\infty}^{+\infty} \gamma_X(\tau) \, \gamma_\psi(2^{-j} au + n) \, d au$$

$$\sum_{n=-\infty}^{\infty} C_X(j,n) e^{-i2\pi fn} = \Gamma_X(2^{-j}f) \times \sum_{n=-\infty}^{\infty} \gamma_{\psi}(n) e^{-i2\pi fn}$$

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### wavelets as stationarizers 1

Theorem (F., '89 & '92)

Wavelet admissibility ( $N \ge 1$ ) guarantees that, if X(t) has stationary increments, then  $d_X(j, k)$  is stationary in k, for any given scale  $2^j$ 

*Proof* — Assuming that X(t) is a s.i. process with X(0) = 0 and  $varX(t) := \rho(t)$ , we have

$$\begin{split} \mathbb{E}X(t)X(s) &= \frac{1}{2} \left( \mathbb{E}X^2(t) + \mathbb{E}X^2(s) - \mathbb{E}\left(X(t) - X(s)\right)^2 \right) \\ &= \frac{1}{2} \left( \mathbb{E}X^2(t) + \mathbb{E}X^2(s) - \mathbb{E}\left(X(t-s) - X(0)\right)^2 \right) \\ &= \frac{1}{2} \left( \rho(t) + \rho(s) - \rho(t-s) \right) \end{split}$$

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### wavelets as stationarizers 2

and it follows that

$$\mathbb{E}d_{X}(j, n)d_{X}(j, m) = \iint \mathbb{E}X(t)X(s)\psi_{jn}(t)\psi_{jm}(s) dt ds$$

$$= \frac{1}{2}\int \rho(t)\psi_{jn}(t)\underbrace{\left(\int \psi_{jm}(s) ds\right)}_{=0} dt$$

$$+\frac{1}{2}\int \rho(s)\psi_{jm}(s)\underbrace{\left(\int \psi_{jn}(t) dt\right)}_{=0} ds$$

$$-\frac{1}{2}\iint \rho(t-s)\psi_{jn}(t)\psi_{jm}(s) dt ds$$

$$= -\frac{1}{2}\int \rho(\tau)\gamma_{\psi}(2^{-j}\tau - (n-m)) d\tau$$

### wavelets as stationarizers 3

- Extension Stationarization can be extended to processes with stationary increments of order p > 1, under the vanishing moments condition N ≥ p
- Application Stationarization applies to *H*-sssi processes (e.g., fBm), with  $\rho(t) = |t|^{2H}$

#### Remark

Nonstationarity is contained in the approximation sequence

### wavelets and scale invariance

 Self-similarity — The multiresolution nature of wavelet analysis guarantees that, if X(t) is H-ss, then

$$\{d_X(j,k), k \in \mathbb{Z}\} \stackrel{d}{=} 2^{j(H+1/2)} \{d_X(0,k), k \in \mathbb{Z}\}$$

for any  $j \in \mathbb{Z}$ 

 Spectral interpretation — Given a "1/f" process, the wavelet tuning condition N > (α – 1)/2 guarantees that

$$\mathcal{S}_X(f) \propto |f|^{-lpha} \Rightarrow \mathbb{E} d_X^2(j,k) \propto 2^{jlpha}$$

### wavelets as decorrelators 1

Theorem (F., '92; Tewfik & Kim, '92)

In the case where X(t) is H-sssi, the condition N > H + 1/2guarantees that

$$\mathbb{E} d_X(j,k) d_X(j,k+n) \sim n^{2(H-N)}, \,\, n 
ightarrow \infty$$

Interpretation

*Competition,* at f = 0, between the (divergent) spectrum of the process and the (vanishing) transfer function of the wavelet:

$$\mathbb{E} d_X(j,k) d_X(j,k+n) \propto \int_{-\infty}^{+\infty} rac{|\Psi(2^j f)|^2}{|f|^{2H+1}} \, e^{i 2\pi n f} \, df$$

## LRD and vanishing moments



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### wavelets as decorrelators 2

Consequence — Long-range dependence (LRD) of a process X can be transformed into short-range dependence (SRD) in the space of its wavelet coefficients d<sub>X</sub>(j,.), provided that the number N of the vanishing moments is high enough

#### Remark

Residual LRD in the approximation sequence

• The case of H-sssi processes — LRD when  $H > 1/2 \Rightarrow$  wavelet SRD needs  $N > 1 \Rightarrow$  Haar not suitable

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## wavelet correlation of fBm in the Haar case



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## wavelet correlation and vanishing moments



## multiresolution building blocks

• Brownian motion — The representation

$$B(t)=\int_0^t W(ds),$$

where W(.) is white Gaussian noise, can be equivalently expressed as

$$B(t) = \sum_{j=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} d_W(j,k) \left( 2^{-j/2} \int_{2^j k}^{2^j (k+t)} \psi(s) \, ds \right)$$

with uncorrelated details (Karhunen-Loève expansion).

 Haar — In the Haar case, this is Lévy's construction (1954) from the Schauder basis

#### approximate construction

 Almost K-L expansion — The quasi-decorrelation property of wavelets (with enough vanishing moments) suggests to approximate fBm by (Wornell, '90):

$$\tilde{B}_{H}(t) = \sum_{j=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} \epsilon(j,k) \psi_{jk}(t),$$

with *uncorrelated* weights such that  $var\epsilon(j, k) = \sigma^2 2^{j(2H+1)}$ 

Spectrum — One can, however, only guarantee that

$$C_1 |f|^{-(2H+1)} \le \mathcal{S}_{ ilde{B}_H}(f) \le C_2 |f|^{-(2H+1)}$$

### exact construction

 Exact expansion — Taking into account low-pass contributions leads to (Sellan et al., '00):

$$B_{H}(t) = \sum_{k=-\infty}^{\infty} S(k)\phi^{(H)}(t-k) + \sum_{j=-\infty}^{0} \sum_{k=-\infty}^{\infty} \epsilon(j,k)\psi_{jk}^{(H)}(t) - b_{0},$$

with S(k) a sum of k FARIMA(0, H - 1/2, 0) processes,  $\epsilon(j, k)$  uncorrelated Gaussian variables with var $\epsilon(j, k) = \sigma^2 2^{j(2H+1)}$ , and  $b_0$  a correcting term ensuring that  $B_H(0) = 0$ 

 Basis functions — Both φ<sup>(H)</sup>(t) and ψ<sup>(H)</sup>(t) have to be specifically designed for a given H

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## assessing scaling via "log-scale" diagrams 1

#### Idea

Given the variance  $v_X(j) := \mathbb{E}d_X^2(j, k)$ , scale invariance is revealed by the linear relation :

$$\log_2 v_X(j) = \alpha j + \text{Const.}$$

*From theory to practice* — The further properties of 1) *stationarization* and 2) *quasi-decorrelation* suggest to use as estimator of  $v_X(j)$  the *empirical variance* 

$$\hat{v}_X(j) := rac{1}{N_j}\sum_{k=1}^{N_j} d_X^2(j,k),$$

where  $N_0$  stands for the data size, and  $N_j := 2^{-j} N_0$ 

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## assessing scaling via "log-scale" diagrams 2

• *Bias correction* — Given that "log  $\mathbb{E} \neq \mathbb{E}$  log", the effective estimator is  $y_X(j) := \log_2 \hat{v}_X(j) - g(j)$ , with

$$g(j) = \psi(N_j/2)/\log 2 - \log_2(N_j/2)$$

and  $\psi(.)$  the derivative of the Gamma function, so that  $\mathbb{E}y_X(j) = \alpha j + \text{Const.}$  in the *uncorrelated* case

 Variance — Assuming stationarization and quasi-decorrelation guarantees furthermore that

$$\sigma_j^2 := \operatorname{var} y_X(j) = \zeta(2, N_j/2) / \log^2 2,$$

where  $\zeta(z, \nu)$  is a generalized Riemann function

### scaling exponent estimation

$$\widehat{\alpha} = \sum_{j=j_{\min}}^{j_{\max}} \frac{S_0 j - S_1}{S_0 S_2 - S_1^2} \frac{1}{\sigma_j^2} y_X(j),$$

with 
$$S_k := \sum_j k / \sigma_j^2, \ k = 0, 1, 2$$

 Bias and variance — We have E α̂ ≡ α, by construction. Assuming Gaussianity, the estimator is moreover asymptotically efficient in the limit N<sub>j</sub> → ∞ (for any j), with

$$\operatorname{var}\widehat{\alpha} \sim 1/N_0$$

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## examples of log-scale diagrams



#### robustness

• Detrending — The vanishing moments condition

$$\int_{-\infty}^{+\infty} t^k \, \psi(t) \, dt = 0, \text{ for } k = 0, 1, \dots N - 1,$$

guarantees that  $d_T(j, n) \equiv 0$  for any T(t) of the form

$$T(t) = \sum_{k=0}^{N-1} a_k t^k$$

In other words, a wavelet with enough vanishing moments makes the transform of Z(t) := X(t) + T(t) blind to a superimposed polynomial trend

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# robustness to polynomial trends



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## robustness to "jumps"



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### related techniques aimed at scaling processes 1

- Aggregation Wavelets offer a natural generalization to aggregation: Haar approximants → Haar details → wavelet details with higher N
- Variogram Wavelets generalize as well variogram techniques (Matheron, 1967), which are based on the increment property  $\mathbb{E}(X(t+\tau) X(t))^2 = \sigma^2 |\tau|^{2H}$ , since increments can be viewed as constructed on the "poorman's wavelet":

$$\psi(t) := \delta(t+\tau) - \delta(t)$$

 Structure functions — Same relation to the generalization of the variogram beyond second order examples concepts wavelets framework multifractals

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## related techniques aimed at scaling processes 2

 Allan variance — A refined notion of variance introduced in the study of atomic clocks stability — is the so-called Allan variance (Allan, '66), defined by

$$\operatorname{var}_{X}^{(\operatorname{Allan})}(T) := \frac{1}{2T^{2}} \mathbb{E}\left[\int_{t-T}^{t} X(s) \, ds - \int_{t}^{t+T} X(s) \, ds\right]^{2}$$

- in the case of *H*-ss processes, Allan variance is such that  $\operatorname{var}_{X}^{(\operatorname{Allan})}(T) \sim T^{2H}$  when  $T \to \infty$
- when evaluated over dyadic intervals, Allan variance identifies to the variance of Haar details:

$$\operatorname{var}_{X}^{(\operatorname{Allan})}(2^{j}) = \operatorname{var}d_{X}^{(\operatorname{Haar})}(j,k)$$

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## related techniques aimed at scaling processes 3

 Fano factor — In the case of a Poisson process P(t) of counting process N(.), one can define the Fano factor as:

$$F(T) := \operatorname{var} N(T) / \mathbb{E} N(T)$$

- for a uniform density  $\lambda$ , we have F(T) = 1 for any Twhereas, for a "fractal" density  $\lambda(t) = \lambda + B_{H}^{(\theta)}(t)$ , we have  $F(T) \sim T^{2H-1}$  when  $T \to \infty$
- interpretation as *fluctuations/average* suggests the *wavelet generalization* given by:

$$F(T) \mapsto F_W(j) := 2^{j/2} \operatorname{var} d_P(j,k) / \mathbb{E} a_P(j,k) \sim 2^{j(2H-1)}$$

when  $j \to \infty$ , and  $F_W^{(\text{Haar})}(j) \equiv F^{(\text{Allan})}(2^j)$ 

## beyond second order

#### Idea

Given the renormalized definition  $T_X(a) := 2^{-j/2} d_X(j,n)|_{j=\log_2 a}$ , one can consider scaling laws which generalize second order behaviors:

$$\mathbb{E}|T_X(a)|^q \propto a^{Hq} = \exp\{Hq \log a\} \quad (\text{``monoscaling''}) \\ \downarrow \\ \exp\{H(q) \log a\} \quad (\text{``multiscaling''}) \\ \downarrow \\ \exp\{H(q) n(a)\} \quad (\text{``cascade''}) \end{cases}$$

### beyond second order

#### Idea

Given the renormalized definition  $T_X(a) := 2^{-j/2} d_X(j,n)|_{j=\log_2 a}$ , one can consider scaling laws which generalize second order behaviors:

$$\mathbb{E}|T_X(a)|^q \propto a^{H(q)} =$$

 $\exp \{H(q) \log a\}$  ("multiscaling")
## back to wavelets and regularity

### Definition

The local regularity of a signal x(t) at a given point  $t_0$  is measured by the Hölder exponent  $h(t_0)$  defined as the supremum of  $\alpha$ 's such that  $|x(t) - x(t_0)| < C |t - t_0|^{\alpha}$  when  $|t - t_0| \rightarrow 0$  (the larger the exponent, the smoother the signal)

$$W_x(a,t) := rac{1}{\sqrt{a}} \int x(s) \,\psi\left(rac{s-t}{a}
ight) \,ds$$

Theorem (Jaffard, '88)

$$h(t_0) = \liminf_{a \to 0, t \to t_0} \frac{\log |W_x(a, t)|}{\log(a + |t - t_0|)}$$

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## multifractality

- some processes (e.g., fBm) are such that h(t) = h for any t, but other ones exhibit *fluctuating* h(t)'s
- global characterization for *h*(*t*) based on *E*(*h*), the set of points *t<sub>h</sub>* with the same Hölder exponent (i.e., such that *h*(*t<sub>h</sub>*) = *h*)



The multifractal spectrum is given by  $D(h) = \dim_{H} E(h)$ , where  $\dim_{H}$  stands for the Hausdorff dimension



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### multiplicative cascades as an example



## monofractal vs. multifractal



credit: P. Abry & H. Wendt 《ロ》《母》《茎》《茎》 薹 少くぐ

## multifractal formalism (Parisi & Frisch, '85)

Given an observation  $x(t) \in \mathbb{R}^d$ :

- **1** start from some *multiresolution* quantity  $M_x(a, t)$  (e.g., increments, aggregation, wavelet details, ...)
- 2 compute structure functions

$$S_x(a,q) = \langle |M_x(a,t)|^q \rangle_t \sim a^{\zeta(q)}, a \to 0_+$$

estimate the scaling exponents as

$$\zeta(q) = \liminf_{a \to 0_+} \frac{\log S_x(a, q)}{\log a}$$

④ deduce the spectrum from a Legendre transform

$$D(h) = \min_{q \neq 0} \{d + qh - \zeta(q)\}$$

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## multifractal formalism and wavelets

### Problem

- (1) instability for  $q \leq 0$
- 2 Legendre transform not valid in general

### Way out (Jaffard, '04)

make use of hierarchical multiresolution quantities  $\Rightarrow$  wavelet leaders



credit: P. Abry

### multifractal formalism and wavelet leaders





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### bootstrap



credit: P. Abry & H. Wendt

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### monofractal vs. multifractal







*credit: P. Abry & H. Wendt* ∢ □ ▷ ∢ ⓓ ▷ ∢ 힅 ▷ ∢ 힅 ▷ ■ うへで

## fully developed turbulence



credit: P. Abry & H. Wendt

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# VanGogh: original vs. copy



credit: P. Abry & H. Wendt

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examples concepts wavelets framework multifractals

# contact, (p)reprints & Matlab codes

- patrick.flandrin@ens-lyon.fr
- http://perso.ens-lyon.fr/patrick.flandrin/
- http://perso.ens-lyon.fr/patrice.abry/

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