

Viscosity as a function of Tg/T

- Tg is the glass transition temperature
- At Tg the viscosity is about 10^{12} Pa s
- For T>Tg the Young modulus falls down of several orders of magnitude

Mechanical measurements



Fig. 3. Evolution of $\tan(\Phi) = G'/G'$ with temperature for successive heating runs for deformed polycarbonate (applied deformation close to 50%, in compression at ambient temperature). (•) first scan up to 339 K; (O) second scan up to 368 K; (+) third scan up to 413 K; (×) last scan up to 448 K, similar to undeformed sample. Between two successive heating runs, the sample is cooled at 6 K/min down to 100 K.

Dielectric measurements



FIG. 4. Dielectric loss vs temperature at 1.2 Hz for pure PC and PC-EG systems during heating and cooling.



Vogel-Fulcher-Tamman law for T>Tg

$$\tau = \tau_o \, \exp\left(\frac{B}{T - To}\right)$$

where To<Tg

Type of glasses

- Structural glasses
- Magnetic glasses
- •Colloids

Frustation AF AF AF Spin is frustrated!

Spin Frustration on the Kagomé Lattice

Energy landscape Bouchaud trap model







(a) Experimental set-up for PMMA. PMMA is the dielectric of a capacitor whose vacuum capacitance is $C_0 = 230 pF$.

 $\epsilon = \epsilon' + i \epsilon''$ is the PMMA dielectric constant.

(b-d) Typical thermal cycles applied to the sample



Dependence on t of ϵ after a quench.

(a) Aging measured at f = 1Hz after a quench at various T_{stop} .

(b) Aging measured after a quench at $T_{stop} = 365K$ at various f.

Memory effect in PMMA

Evolution of \mathbf{E} at f=0.1Hz as a function of T



 ϵ_r = dielectric constant measured with continuous ramp ϵ_m = dielectric constant measured with a cooling stop



Evolution of $\mathbf{\mathcal{E}}$ at f=0.1Hz as a function of T

 ϵ_r = dielectric constant measured without a cooling stop ϵ_m = dielectric constant measured with a cooling stop



From:

V. Dupuis, E. Vincent, J.P. Bouchaud, J. Hammann, A.Ito, H. Aruga Katori,

Aging, rejuvenation and memory effects in Ising and Heisenberg spin glasses, Phys. Rev B 64 (17),174204,(2001).

Also in cond-mat/0104399



- The sample reminds its thermal history

- The response of the system depends on the quench speed

Memory effects and trap model



Fig. 6. Schematic picture of the hierarchical structure of the metastable states as a function of temperature.

Aging in glassy materials

Aging has been often characterized by studing the response functions of the systems

Smart experimental procedures, based either

on multiple cycles of cooling, heating and waiting times

or

on the modulation of the applied external fields

have shown the existence of spectacular effects of aging in glassy materials, such as

rejuvenation and memory.

These studies have been extremely useful to fix several important constraints for the phenomenological models of aging.

Question: is the analysis of fluctuations useful ?

Outline

- 1) Thermal fluctuations and the Fluctuation Dissipation Relations during aging
- 2) The electrical thermal noise of two materials:
 - a) a polymer after a quench
 - b) a colloidal glass during the sol-gel transition.
- 3) Comparisons of the experimental results with those of other experiments and of models of aging.
- 4) The mechanical noise.
- 5) Conclusions

FLUCTUACTION DISSIPATION THEOREM

in thermodynamic equilibrium

V and q are two conjugate variables

 $R(\omega) = \frac{\delta V(\omega)}{\delta q(\omega)}$ is the response function

The thermal fluctuation spectrum $S(\omega) = \langle |V(\omega)|^2 \rangle$ is

$$S(\omega) = \frac{4 \text{ K T}}{\omega} \quad Im\{ \text{ R}(\omega) \}$$

Typical examples are :



Fluctuation Dissipation Relation (FDR) in a weakly out of equilibrium system (Cugliandolo,Kurchan 1992.)

In a glass at $T < T_G$ the physical properties of the material depend on the aging time t_w after the temperature quench. Thus FDR takes the following form:

$$S(\omega, t_w) = \frac{4 K_B T_{eff}(\omega, t_w)}{\omega} Im\{R_{Vq}(\omega, t_w)\}$$

FDR can be used to define an effective temperature of the system

$$T_{eff}(\omega, t_w) = \frac{S(\omega, t_w) \ \omega}{4 \ K_B \ Im\{R_{Vq}(\omega, t_w)\}}$$

At equilibrium $T_{eff}(\omega, t_w) = T$

In terms of correlation function FDR takes the form

$$-C(t,t_w) + C(t_w,t_w) = K_B T_{eff}(t,t_w) R(t,t_w)$$

where $C(t, t_w)$ is the correlation function and $R(t, t_w)$ the integrated response

KOB, **BARRAT**, Fluctuation dissipation ratio in an aging Lennard-Jones glass, Europhys. Lett. 46, 637 (1999)



 $-C(t_w + \tau, t_w) + C(t_w, t_w) = K_B T_{eff}(t_w + \tau, t_w) M(t_w + \tau, t_w)$



FDR in out of equilibrium system

$$S(\omega, t_w) = \frac{4 K_B T_{eff}(\omega, t_w)}{\omega} Im\{R_{Vq}(\omega, t_w)\}$$

 $-C(t,t_w)+C(t_w,t_w)=K_B T_{eff}(t,t_w) R_{Vq}(t,t_w)$

Theoretical Background

 This definition of temperature seems to be appropriate for several systems.
Cugliandolo, Kurchan, Peliti (1997), Kob, Barrat (1999)

Berthier, Barrat (2002), Liu, Nagel (2002),

Sciortino(2002).....

2) The robustness of this definition of temperature has been questioned .

S. Fielding, P. Sollich, (2002), Perez-Madrid, Reguera, Rubi (2002).

Experiments

- 1970 x-ray scattering on PMMA (Weandorf and Fisher)
- 1999 Grigera, Israeloff, super-cooled liquid
- 2001 Bellon, Ciliberto, sol-gel transition
- 2002 Herisson and Ocio, spin-glass
- 2002 et 2005 Buisson, Ciliberto, polymer

X-ray experiments

Intensity I(0) of scattered x-rays at small angles is related to the density fluctuations $\delta \rho$:

$$\frac{<\delta\rho^2>}{\rho^2}\propto I(o)$$

From FDT

$$<\delta\rho^2> = \frac{K_B T \rho^2 \chi_T}{V}$$

where χ_T is the isothermal compressibility.

Weandorf and Fisher found a violation between 2.5 and 5 of this expression for various polymers.

Comparison with theory:

$$-C(t_w + \tau, t_w) + C(t_w, t_w) = K_B T R(t_w + \tau, t_w)$$

only two points on the plane (C, R) are available



Herisson and Ocio

Insulating spin glass $CdCr_{1.7}In_{0.3}S_4$

Tg=16K

Quench at 0.8Tg.



Fig. 4. A typical thermal history of the sample for a 4,500 srecord at 10 K. In inset, detail on the crucial part, the last 3 K's cooling.



FIG. 1: a) Schematic of the detection circuit. The pick-up coil (right side), containing the cylindrical sample, is a third order gradiometer made of +3 -6 +6 -3 turns. b) Calibration is obtained by measuring relaxation versus correlation in a high conductivity copper sample at equilibrium at 4.2K.



FIG. 2: Aging and scaling of (a) correlation (b) relaxation at $T = 0.8T_g$. Both are measured for waiting times t' = 100(\circ), 200 (\triangle), 500(\bigtriangledown), 1000(\diamond), 2000 (+), 5000(\times), 10000(*) seconds from bottom to top. Reported error-bars on correlation have a length of two standard-deviation, corresponding to averages over records. In insets, scaling of the aging parts versus $\zeta = (t^{1-\mu} - t'^{1-\mu})/(1-\mu)$, using $\mu = 0.87$. The stationary parts are found to obey a power-law decrease with an exponent $\alpha = 0.05$.



FIG. 3: FD-plot. Relaxation measurements are plotted versus correlation functions for each t'. The dot-dashed line (FDT line) is calculated for T = 0.8Tg = 13.3K, from the calibration obtained with the copper sample. The dashed line represents the scaling extrapolation for $t' \to \infty$. The branching point with the FDT line, corresponds to $\tilde{C} = q_{EA}$ (square symbol, with size giving the error range). In Inset, the same data in the whole range.