

L'ordinateur quantique à l'épreuve des erreurs

Christophe VUILLOT

Inria

SIESTE - 23 octobre 2024

About Me

Current Position (Jan. 2021 - now)

The logo for Inria, featuring the word "Inria" in a stylized, red, cursive script font.

Permanent, “Chargé de Recherche”, at Inria Nancy, *Mocqua* team

About Me

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Permanent, “Chargé de Recherche”, at Inria Nancy, *Mocqua* team

Previous Research Experience



Postdoc
2019 - 2020



PhD
2015 - 2019



internships
2015 2014 2012

Education



Upcoming Update

Future Position



Will be on (temporary) leave from Inria to work at Alice&Bob on trying to build a quantum computer

Computation

Computability

Church-Turing Thesis (20th): *Computability* = Turing machines

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- *Complexity* theory → accounts for time-scale and resources
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Physics of Computation Conference in 1981

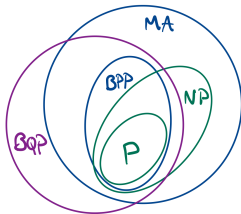


Quantum processes $\notin BPP$, (although $\in PSPACE$).

Quantum Computation

Some Striking Quantum Algorithms

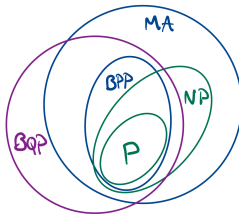
- Shor (1994) factors in polynomial time (break RSA).
- Grover (1996) searches an unstructured list in \sqrt{N} queries.



Quantum Computation

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Quantum Computer Model

Assume *idealized* perfect quantum computers!

Error Correction and Fault-Tolerance

Quantum Error Correction

- Encode quantum information redundantly
- Extract information about errors, process it and correct

Fault-Tolerant Quantum Computation

- Compute directly on encoded quantum information without weakening the protection
- Works when every quantum process is unreliable

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Threshold Theorems (informal)

Given an error model there exist so-called fault-tolerant protocols and a threshold such that any computation can be simulated to any desired precision if the noise strength is below the threshold.

The Quantum Stack

FTQC Aware Compilation

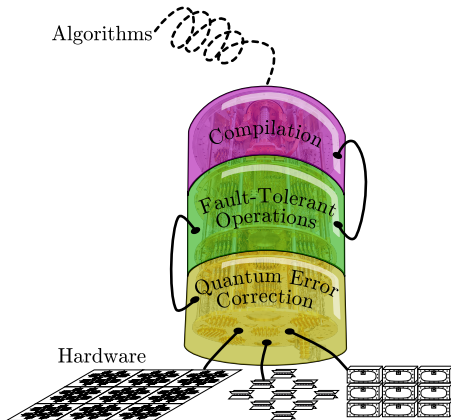
Specific gate sets

Fault-Tolerant Logic

Transversal gates, code deformation

Quantum Error Correction

Finite as well as infinite dimensional systems



Quantum Bit

A physical system with two perfectly distinguishable states:

Examples

- Electron around an atom/ion
- Photon in one of two paths
- Spin of a particle
- Collective degree of freedom
- ...

Notation

$|0\rangle$ $|1\rangle$

Superposition

Qubit Hilbert Space

Quantum states live in a Hilbert space, $|\Psi\rangle \in \mathcal{H}$:

$$|\Psi\rangle = \alpha |0\rangle + \beta |1\rangle, \quad |\alpha|^2 + |\beta|^2 = 1, \quad \alpha, \beta \in \mathbb{C}.$$

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$$\langle 0|0\rangle = \langle 1|1\rangle = 1, \quad \langle 0|1\rangle = 0$$

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Dual Basis

$$|+\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}, \quad |-\rangle = \frac{|0\rangle - |1\rangle}{\sqrt{2}},$$

defines a perfectly valid basis

$$\langle +|+\rangle = \langle -|-\rangle = 1, \quad \langle +|-\rangle = 0.$$

Standard Measurement

$$|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle \rightarrow \{|0\rangle, |1\rangle\}?$$

$|0\rangle$ with proba $|\alpha|^2$

$|1\rangle$ with proba $|\beta|^2$

Dual Measurement

$$|\Psi\rangle = \alpha |0\rangle + \beta |1\rangle \rightarrow \{|+\rangle, |-\rangle\}?$$

$|+\rangle$ with proba $\frac{|\alpha+\beta|^2}{2}$

$|-\rangle$ with proba $\frac{|\alpha-\beta|^2}{2}$

Quantum Gates

Unitary Transformations

Quantum transformations are linear and *unitary* (consistent with probabilistic interpretation of measurements)

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Pauli Operators form a basis

$$\forall M, \quad M = \alpha_{\mathbb{1}}\mathbb{1} + \alpha_X X + \alpha_Z Z + \alpha_{XZ} XZ$$

Several Quantum Systems

Tensor Product of Hilbert Spaces

The joint Hilbert space of two quantum systems described by $|\Psi\rangle \in \mathcal{H}_1$ and $|\Phi\rangle \in \mathcal{H}_2$ is given by the tensor product

$$\mathcal{H}_1 \otimes \mathcal{H}_2.$$

A basis for $\mathcal{H}_1 \otimes \mathcal{H}_2$ can be obtained by the cartesian product of a basis of \mathcal{H}_1 and a basis of \mathcal{H}_2 .

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Entangled States

Entangled states are states that cannot be factorized as a single product state

$$\alpha |\Psi_1\rangle \otimes |\Phi_1\rangle + \beta |\Psi_2\rangle \otimes |\Phi_2\rangle \neq |\Psi\rangle \otimes |\Phi\rangle.$$

Many Qubits

We consider having $n \in \mathbb{N}^*$ qubits

Hilbert Space

This is a 2^n dimensional space with standard basis

$$\{|\mathbf{b}\rangle \mid \forall \mathbf{b} \in \mathbb{F}_2^n\}, \quad \mathbb{F}_2 = \{0, 1\}.$$

Quantum States

Quantum states are given by

$$|\Psi\rangle = \sum_{\mathbf{b} \in \mathbb{F}_2^n} \alpha_{\mathbf{b}} |\mathbf{b}\rangle, \quad \sum_{\mathbf{b} \in \mathbb{F}_2^n} |\alpha_{\mathbf{b}}|^2 = 1.$$

Multi-Qubit Pauli Operators

Single Qubit X and Z

$$X_j |\mathbf{b}\rangle = |\mathbf{b} \oplus \mathbf{e}_j\rangle, \quad \mathbf{e}_j = (0 \ \cdots \ 0 \ 1 \ 0 \ \cdots \ 0)$$

$$Z_j |\mathbf{b}\rangle = (-1)^{b_j} |\mathbf{b}\rangle = (-1)^{\mathbf{b} \cdot \mathbf{e}_j} |\mathbf{b}\rangle.$$

Multi-Qubit X and Z

$$X(\mathbf{x}) = \prod_{j=1}^n X_j^{x_j}, \quad \mathbf{x} \in \mathbb{F}_2^n \quad Z(\mathbf{z}) = \prod_{j=1}^n Z_j^{z_j}, \quad \mathbf{z} \in \mathbb{F}_2^n.$$

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Pauli Measurements

Z Parity Measurement

Given $\mathbf{z} \in \mathbb{F}_2^n$, it splits in half $\mathbf{b} \in \mathbb{F}_2^n$ between:

- $(-1)^{\mathbf{b} \cdot \mathbf{z}} = 1$
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There is a quantum measurement that measures the parity of $\mathbf{z} \cdot \mathbf{b}$

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$$\begin{aligned}
 |\Psi\rangle = \sum_{\mathbf{b} \in \mathbb{F}_2^n} \alpha_{\mathbf{b}} |\mathbf{b}\rangle &\xrightarrow{M_{\mathbf{z}=m \in \{0,1\}}} |\psi'_m\rangle \propto \sum_{\mathbf{b}, \mathbf{z} \cdot \mathbf{b} = m} \alpha_{\mathbf{b}} |\mathbf{b}\rangle \\
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 \end{aligned}$$

X Parity Measurement

Same but in the dual basis $\{|+\rangle, |-\rangle\}$ and after measurement:

$$X(\mathbf{x}) |\Psi'_m\rangle = (-1)^m |\Psi'_m\rangle.$$

Classical Error Correction

3-bit Repetition Code

Define the code $\mathcal{C} = \{\mathbf{b}_0, \mathbf{b}_1\}$

$$0 \rightarrow \mathbf{b}_0 = (0 \ 0 \ 0),$$

$$1 \rightarrow \mathbf{b}_1 = (1 \ 1 \ 1).$$

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Parity-check matrix

$$H = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 1 \end{pmatrix}, \quad H\mathbf{b}_0^T = H\mathbf{b}_1^T = \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$

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Syndrome of bit-flip error \mathbf{e}

The parity checks give information about errors only

$$H(\mathbf{b}_0 + \mathbf{e})^T = H(\mathbf{b}_1 + \mathbf{e})^T = H\mathbf{e}^T$$

Linear Error Correcting Code

Definition

A linear code $\mathcal{C} \subset \mathbb{F}_2^n$ can be defined by a parity check matrix $H \in \mathbb{F}_2^{r \times n}$ with

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Syndrome and Decoding

$$\forall \mathbf{c} \in \mathcal{C}, H(\mathbf{c} + \mathbf{e})^T = H\mathbf{e}^T$$

Given \mathbf{s} find the smallest \mathbf{e} such that $H\mathbf{e}^T = \mathbf{s}$.

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Parameters $[n, k, d]$

The length of the codewords, n , the dimension of \mathcal{C} , k , the minimum distance, d ,

$$k = n - \text{rank}H, \quad d = \min_{\mathbf{c} \in \mathcal{C}, \mathbf{c} \neq \mathbf{0}} |\mathbf{c}|.$$

Quantum Error Correcting Code

General Idea

- Correct for bit-flips and phase-flips separately with X and Z parity measurements
- Problem: not all X and Z parity measurements are compatible because $XZ = -ZX$

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Quantum Commutation Constraint

$$\begin{aligned} Z(\mathbf{z})|\Psi\rangle = |\Psi\rangle \quad \wedge \quad X(\mathbf{z})|\Psi\rangle = |\Psi\rangle \\ \Rightarrow X(\mathbf{x})Z(\mathbf{z}) = Z(\mathbf{z})X(\mathbf{x}) \\ \Rightarrow \mathbf{x} \cdot \mathbf{z} = 0 \end{aligned}$$

Quantum Error Correcting Code

Definition

Given two parity check matrices $H_Z \in \mathbb{F}_2^{r_Z \times n}$ and $H_X \in \mathbb{F}_2^{r_X \times n}$ such that

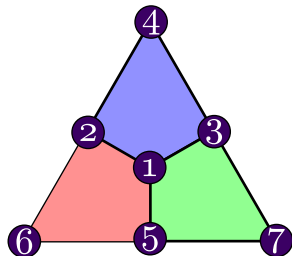
$$H_Z H_X^T = 0,$$

We define the quantum error correcting code $\mathcal{C}(H_Z, H_X)$ on n qubits as the quantum states satisfying the Z and X parities defined by H_Z and H_X :

$$\mathcal{C} = \{|\Psi\rangle \mid \forall \mathbf{s}_Z \in \mathbb{F}_2^{r_Z}, \forall \mathbf{s}_X \in \mathbb{F}_2^{r_X}, Z(\mathbf{s}_Z H_Z) |\Psi\rangle = X(\mathbf{s}_X H_X) |\Psi\rangle = |\Psi\rangle\}$$

Steane's Code Example

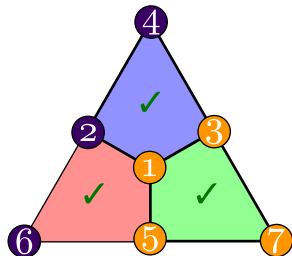
$$\left. \begin{array}{l} H_X \\ \parallel \\ H_Z \end{array} \right\} = \begin{array}{c} c_1 \\ c_2 \\ c_3 \end{array} \begin{pmatrix} \textcircled{1} & \textcircled{2} & \textcircled{3} & \textcircled{4} & \textcircled{5} & \textcircled{6} & \textcircled{7} \\ 1 & 1 & 1 & 1 & 0 & 0 & 0 \\ 1 & 1 & 0 & 0 & 1 & 1 & 0 \\ 1 & 0 & 1 & 0 & 1 & 0 & 1 \end{pmatrix}$$



Stabilizer checks

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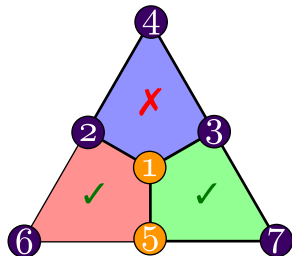
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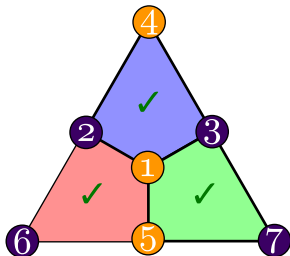
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Error

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Logical operator

Properties of Quantum Codes

Logical operators

Pauli operators which are not detectable but not stabilizers act non-trivial on the codespace

$$\mathcal{L}_X = \ker H_Z / \text{im} H_X, \quad \mathcal{L}_Z = \ker H_X / \text{im} H_Z$$

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Parameters $[[n, k, d]]$

The number of qubits, n , the number of logical qubits in \mathcal{C} , k , the minimum distance, d ,

$$k = n - \text{rank } H_Z - \text{rank } H_X, \quad d = \min_{c \in \mathcal{L}_X \cup \mathcal{L}_Z, c \neq 0} |c|.$$

Desirable Properties of Quantum Codes

Given a quantum code \mathcal{C} with parameters $[[n, k, d]]$ one might want it to have

- Large encoding rate: k/n
- Large distance: d
- Good and efficient syndrome decoding algorithm
- (Large) Threshold
- Easy implementation \rightarrow sparse/local matrices H_Z and H_X
- Fault-tolerant quantum gates...

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Good quantum LDPC codes exists

The existence of good quantum LDPC codes with efficient decoders has been established first in 2022, i.e $[[n, k = \Theta(n), d = \Theta(n)]]$.

Fruitful Connection to Homology

Chain Complex

$$\mathcal{C} : \quad C_2 \quad \xrightarrow{\partial} \quad C_1 \quad \xrightarrow{\sigma} \quad C_0 \quad \text{with } \sigma \circ \partial = 0$$

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H_X H_Z^T

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Homology Group = X Logical Operators

$$H_1(\mathcal{C}, \mathbb{Z}) = \ker \sigma / \text{im} \partial = \ker (H_Z) / \text{im} (H_X)$$

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 &= \mathcal{L}_X
 \end{aligned}$$

Cohomology = Logical Z

Cochain Complex

$$\begin{array}{ccccc}
 \mathcal{C}^* : & C_2^* & \xleftarrow{\partial^*} & C_1^* & \xleftarrow{\sigma^*} & C_0^* \\
 & \parallel & & \parallel & & \parallel \\
 & \mathbb{F}_2^{r_x} & & \mathbb{F}_2^n & & \mathbb{F}_2^{r_z} \\
 & \parallel & & \parallel & & \parallel \\
 & \text{syndrome} & & \text{operators} & & \text{stabilizers}
 \end{array}$$

H_X^T H_Z

Cohomology = Logical Z

Cochain Complex

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Cohomology Group = Z Logical Operators

$$H^1(\mathcal{C}, \mathbb{T}) = \ker \partial^* / \text{im} \sigma^* = \ker (H_X) / \text{im} (H_Z)$$

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Cohomology Group = Z Logical Operators

$$\begin{aligned}
 H^1(\mathcal{C}, \mathbb{T}) &= \ker \partial^* / \text{im} \sigma^* = \ker (H_X) / \text{im} (H_Z) \\
 &= \mathcal{L}_Z
 \end{aligned}$$

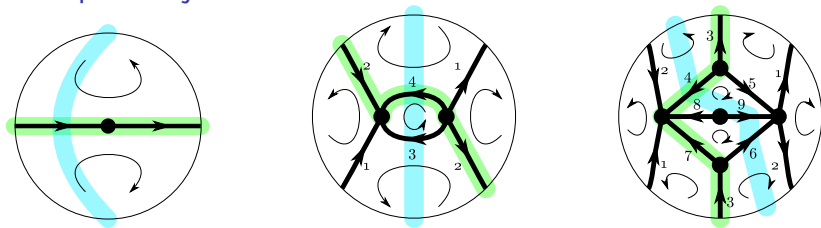
Codes from Cellular Homology in 2D

$$\begin{array}{ccccc}
 \mathcal{C} : & C_2 & \xrightarrow{\partial} & C_1 & \xrightarrow{\sigma} & C_0 & \text{with } \sigma \circ \partial = 0 \\
 & \parallel & & \parallel & & \parallel & \\
 & \mathbb{F}_2^F & & \mathbb{F}_2^E & & \mathbb{F}_2^V & \\
 & \parallel & & \parallel & & \parallel & \\
 & \text{faces} & & \text{edges} & & \text{vertices} &
 \end{array}$$

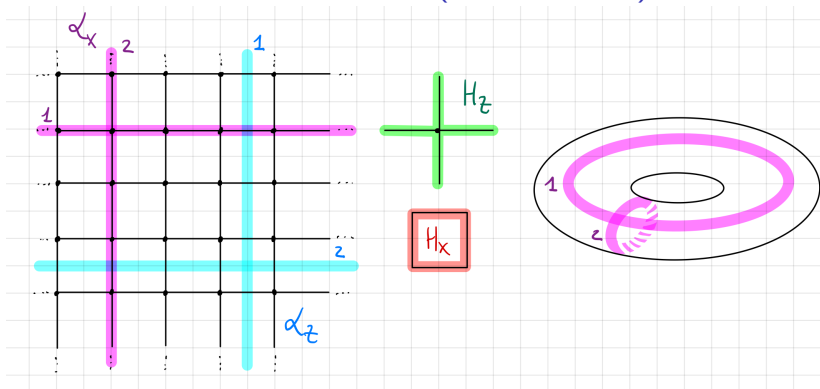
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 \end{array}$$

Example: Projective Plane



The Toric Code (Kitaev 1997)



For a genus g surface, distance correspond to shortest non-trivial cycles and

$$H_1(\mathcal{C}, \mathbb{F}_2) = H^1(\mathcal{C}, \mathbb{F}_2) = \mathbb{F}_2^{2g}$$

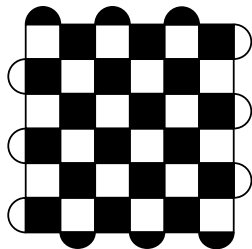
The Surface Code: The Volkswagen of Quantum Codes

Many Advantages

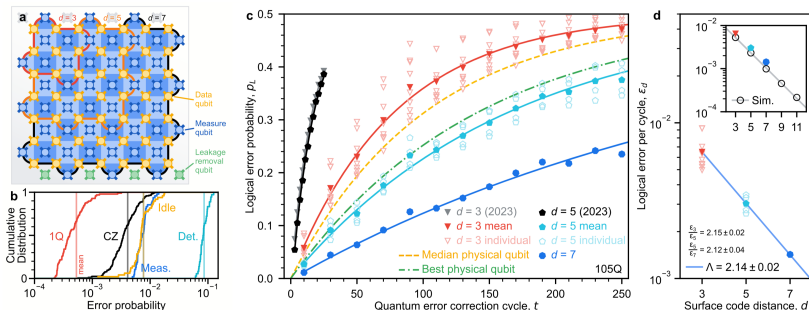
- 2D local connectivity with degree 4
- Good decoder with high threshold
- Fault-tolerant gates understood
- Experimentally demonstrated (2024)

One Weakness

Vanishing encoding rate!



Google's 2024 Surface Code Experiment



Bravyi-Poulin-Terhal Bound

The surface code is optimal for 2D, for any $[[n, k, d]]$ quantum code:

Constraints from 2D Locality

$$kd^2 = O(n)$$

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Constraints from 2D Locality

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How to Improve

- Have to soften connectivity constraints

In the Lab

Physical Systems

Physical systems are often richer than qubits, for instance bosonic modes are infinite dimensional quantum systems

- photons in a cavity
- motion of an ion

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Physical Systems

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Error Correction

Can be formulated more generally as $\mathcal{H}_{\text{code}} \subset \mathcal{H}$. Often just replacing \mathbb{F}_2 with the relevant G is enough

- $G = \mathbb{Z}/\mathbb{Z}_d \rightarrow$ Qudit error correction
- $G = \mathbb{R} \rightarrow$ Grid-state bosonic codes (GKP)
- $G = \mathbb{Z} \rightarrow$ Quantum Rotor codes

Cat Qubits

Bosonic Codes

- Uses bosonic modes, i.e. $\mathcal{H} = \ell_2(\mathbb{N})$
- Can encode a qubit
- Can feature extremely biased noise

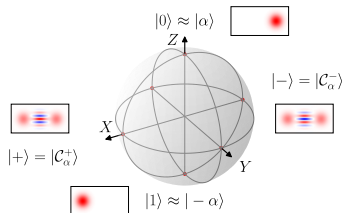
Cat Qubits 🐱

Bosonic Codes

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- Can encode a qubit
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Cat Qubit

- $|\bar{0}\rangle \simeq |\alpha\rangle$, $|\bar{1}\rangle \simeq |-\alpha\rangle$
- X -errors $\sim e^{-2|\alpha|^2}$
- Z -errors $\sim c|\alpha|^2$



Repetition Cat Qubit

Assuming we can get negligible X errors \rightarrow Classical code to correct remaining Z errors

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- 1D layout
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$$k\sqrt{d} = O(n)$$

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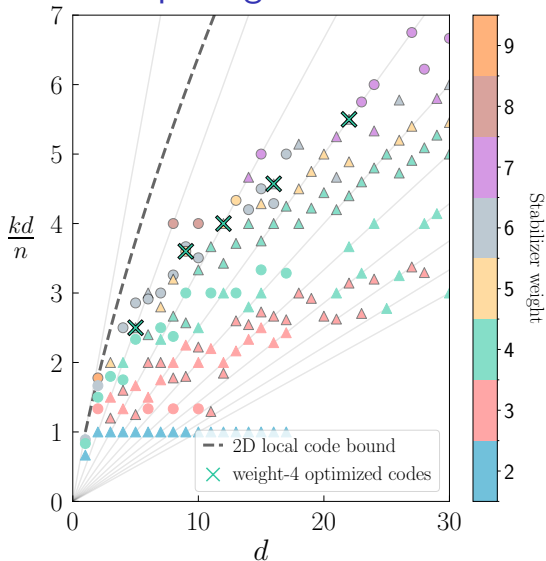
Classical Bravyi-Poulin-Terhal bound in 2D

$$k\sqrt{d} = O(n)$$

Repetition Code not Optimal

$$k\sqrt{d} = 1 \times \sqrt{n}$$

Exploring All Codes

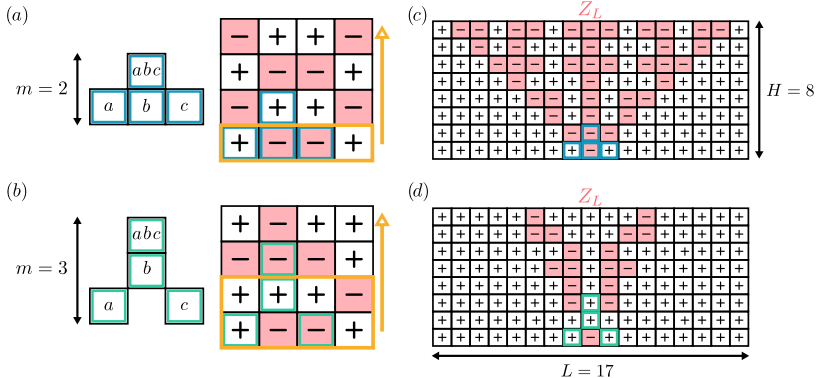


Cellular Automaton Codes

Characteristics

- $L \times H$ Cylinder
- Translation invariant check with “pointed” shape
- k linear in L
- d increases with H
- fixed $H \rightarrow$ fixed kd/n

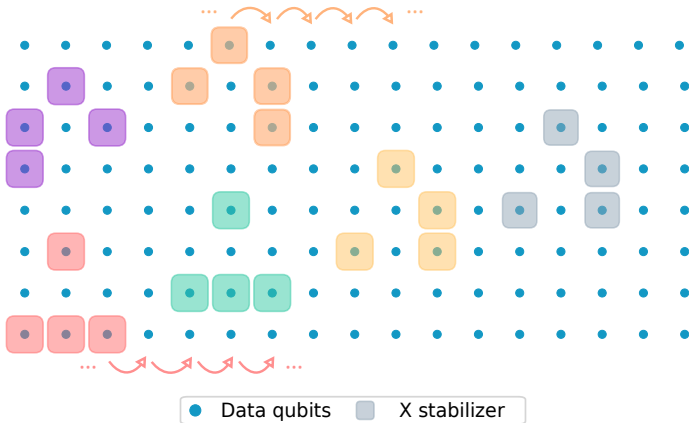
Cellular Automaton Codes Examples



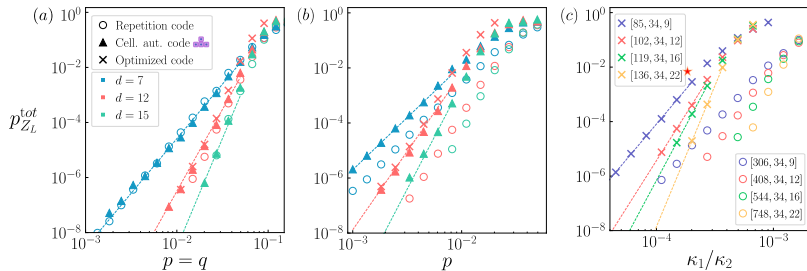
Relaxing Vertical Translation Invariance

$[n, k, d]$	kd/n	$(H, L = L^* + \ell)$	Stabilizer shapes (bottom to top)
$[20 + 4\ell, 10 + 2\ell, 5]$	2.5	$(4, 5 + \ell)$	
$[55 + 5\ell, 22 + 2\ell, 9]$	3.6	$(5, 11 + \ell)$	
$[78 + 6\ell, 26 + 2\ell, 12]$	4	$(6, 13 + \ell)$	
$[119 + 7\ell, 34 + 2\ell, 16]$	4.6	$(7, 17 + \ell)$	
$[136 + 8\ell, 34 + 2\ell, 22]$	5.5	$(8, 17 + \ell)$	

The Chosen One



Numerical Simulations



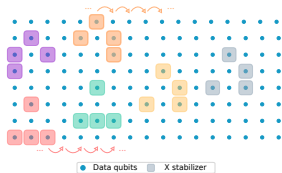
LDPC-Cat Code

LDPC-Cat Code

- $[136 + 8l, 34 + 2l, 22]$ code
- 2D local
- Degree 4 tanner graph

Improvement over Surface code

In an intermediate regime (logical error rate $\sim 10^{-8}$) can get 44-fold reduction in resource cost.



Summary

- Concepts works in theory
- First experimental proofs are emerging
- Demonstrating that it scales in practice is still daunting
- Active area of research with lot to understand

Quantum in France

Quantum Programming
Languages

Inria Nancy, Inria Saclay

Quantum Algorithms

Irif Paris

Quantum Info and Error
Correction

Inria Paris, Lyon, Bordeaux,
CEA Grenoble

